

Workshop on Fission Dynamics
Chongqing, China
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Microscopic Theory of Angular Momentum Distributions Across the Full Range of Fission Fragments

Petar Marević

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Outline

1. Introduction
2. Theoretical Framework
3. n-Induced Fission of ^{235}U and ^{239}Pu
4. Conclusion

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1. Introduction

2. Theoretical Framework

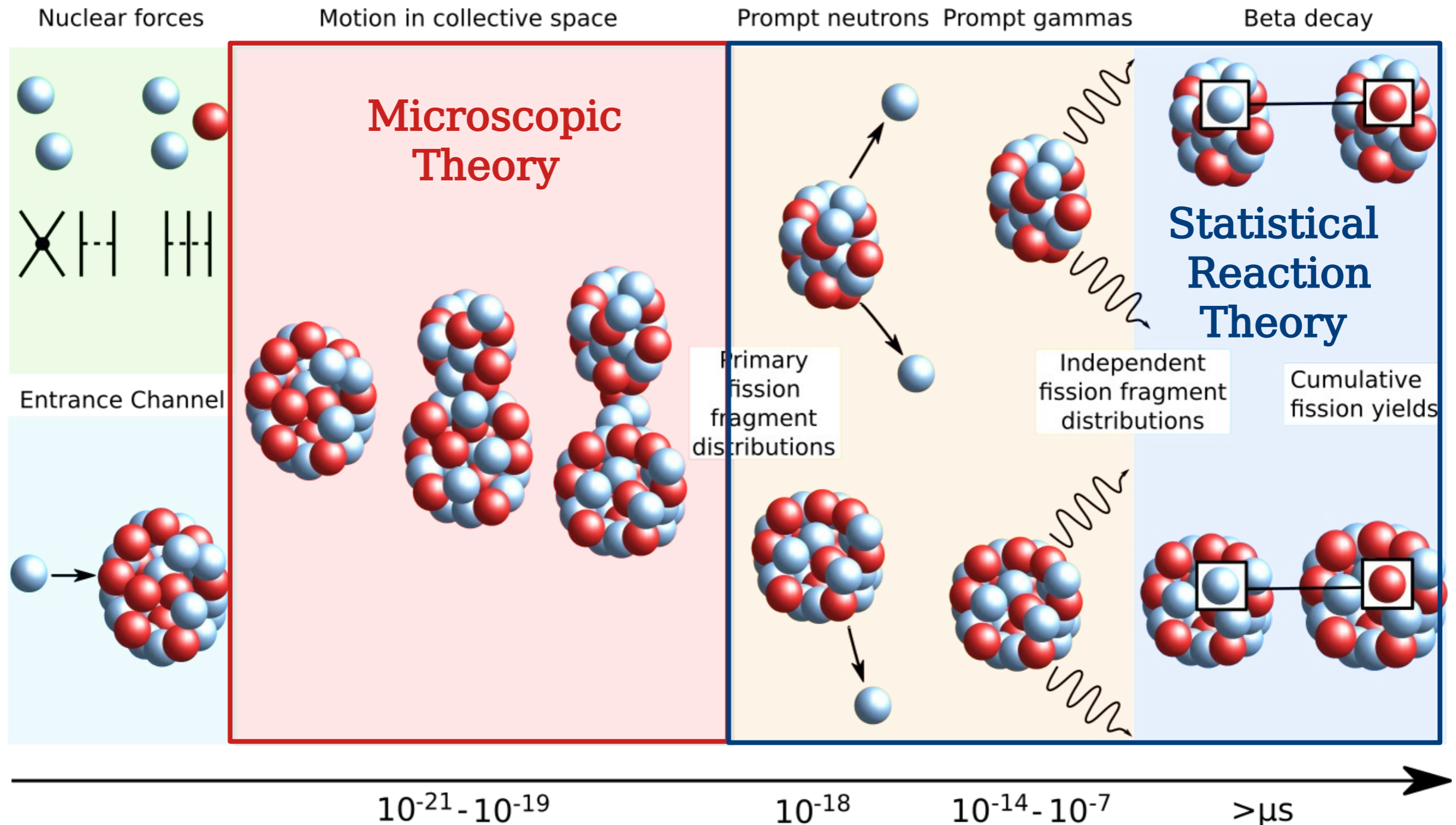
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Introduction

Microscopic Approach to Nuclear Fission

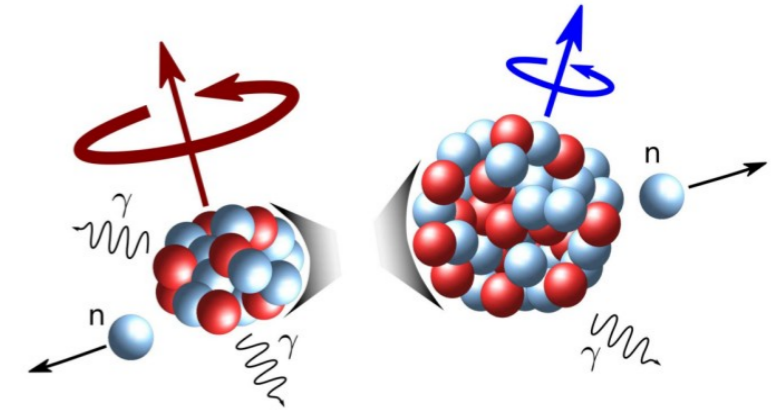
- Microscopic theory = using our best knowledge of internucleon forces and many-body techniques (DFT)



Introduction

Angular Momentum of Fission Fragments (FFs)

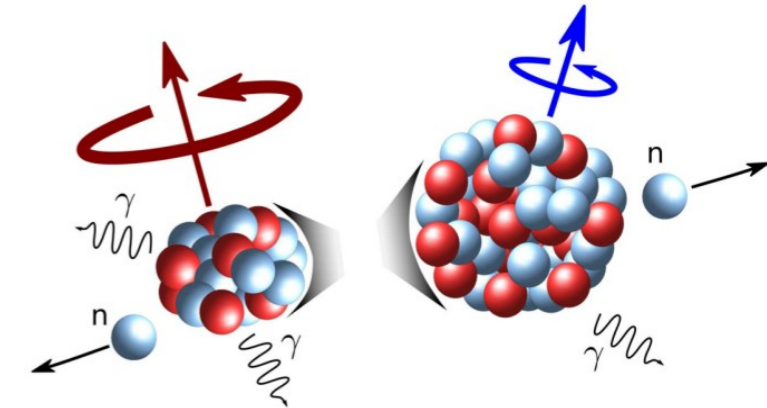
- FFs at scission are hot, deformed, and rotating
 - High excitation energy
 - Large deformation, typically different than the g.s.
 - Distribution of angular momentum (AM)



Introduction

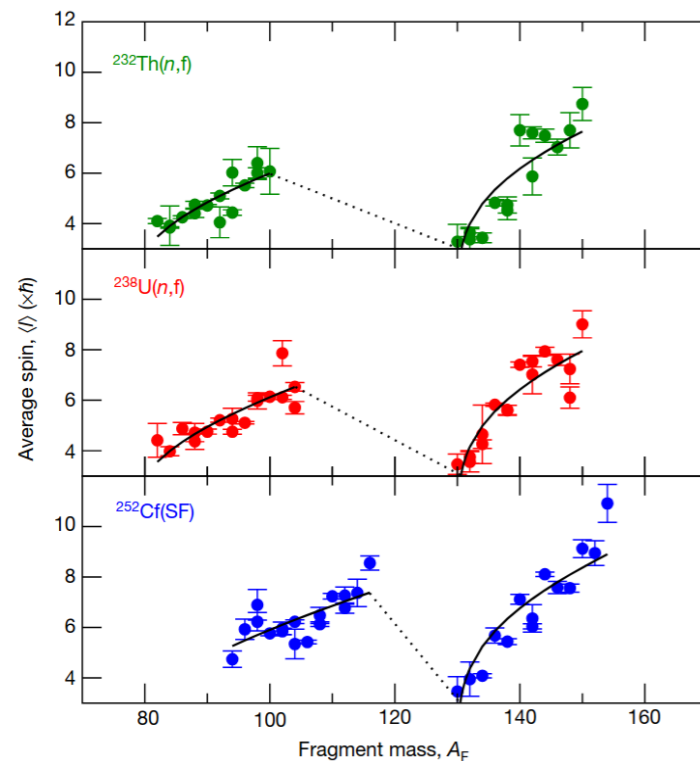
Angular Momentum of Fission Fragments (FFs)

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- Recently, the study of AM of FFs experienced a *renaissance*

Wilson *et al.*, Nature **590**, 566 (2021).



- A. Bulgac *et al.*, PRL **126**, 142502 (2021).
- P. Marević *et al.*, PRC **104**, L021601 (2021).
- J. Randrup and R. Vogt, PRL **127**, 062502 (2021).
- I. Stetcu *et al.*, PRL **127**, 222502 (2021).
- A. Bulgac *et al.*, PRL **128**, 022501 (2022).
- J. Randrup, PRC **106**, L051601 (2022).
- G. Scamps, PRC **106**, 054614 (2022).
- G. Scamps *et al.*, PRC **108**, L061602 (2023).
- T. Dossing *et al.*, PRC **109**, 034615 (2024).
- G. Scamps *et al.*, arXiv:2512.02207v1 (2025).
- ... and others

Most studies still focus on the **most probable fragmentation**.

Introduction

Statistical Decay of Fission Fragments

- FF initial conditions are essential inputs to decay models

- (N_F, Z_F) , Binding Energies
- Primary Yields $Y(N_F, Z_F)$
- Excitation E^* of FFs
- $|c(J^\pi)|^2$ Distributions
- Level Densities $\rho(E^*, J^\pi)$
- Transmission Coeff. $T_a(\varepsilon)$
- γ Strength Function
- β^- decay half-lives

Inputs from Data/Models



CGMF
FREYA
FIFRELIN
YAHFC
GEF
...

Statistical Modeling



- n and γ Multiplicities
- n and γ Spectra
- n and γ Angular Corr.
- Independent Fission Yields
- Cumulative Product Fission Yields
- ...

Correlated Observables

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Inputs from Data/Models

Statistical Modeling

Correlated Observables

- Traditionally, inputs are taken from experiments and empirical models
- We've entered an era where **microscopic theory is becoming useful**
 - Binding energy tables (~2000)
 - Level Densities (~2001) and γ strength function (~2002)
 - Primary FF yields for even (~2005) and odd (~2022) actinides
 - Excitation energies (~2016, but for most likely FFs)
 - **Angular momentum distributions (~2021)**

Introduction

What Are We Trying to Achieve?

Objectives:

1) Establish a microscopic framework for predicting AM distributions in the **full range** of FF masses and charges.

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- 4) Use the results to **inform phenomenological models** (deformation effects, isobaric dependence, ...).
- 5) Assess how microscopic AM distributions **impact the FF decay process** (neutron and photon spectra).

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Theoretical Framework

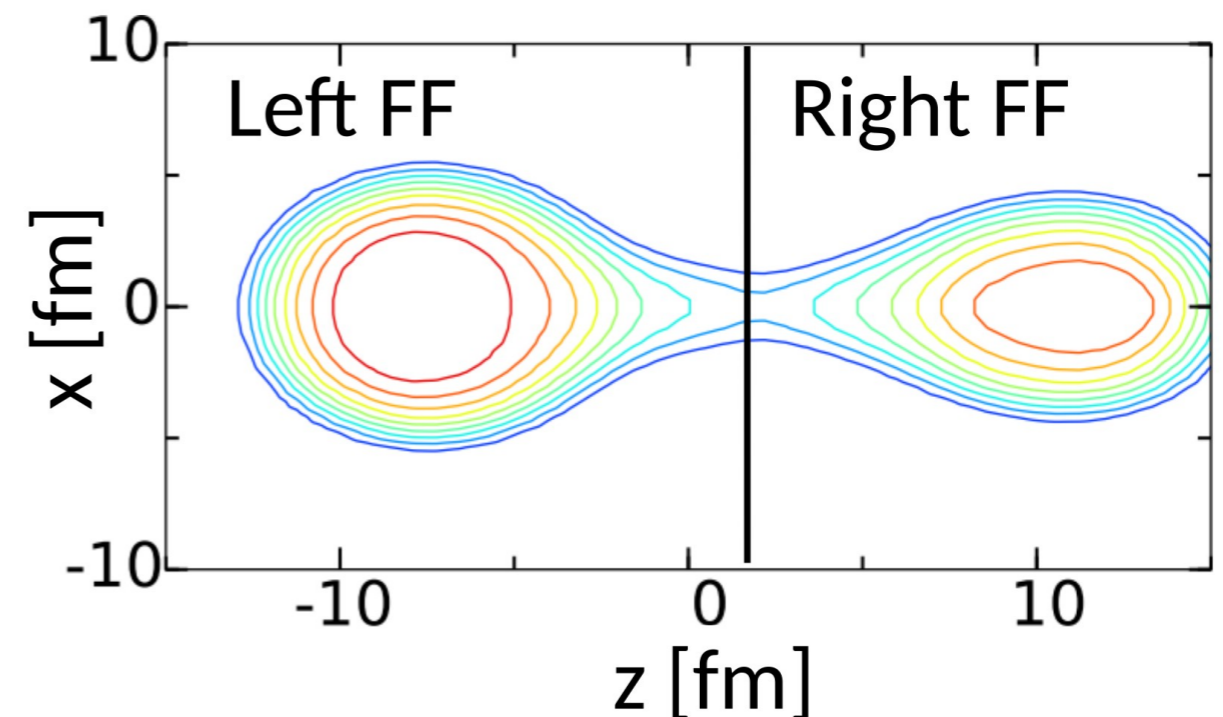
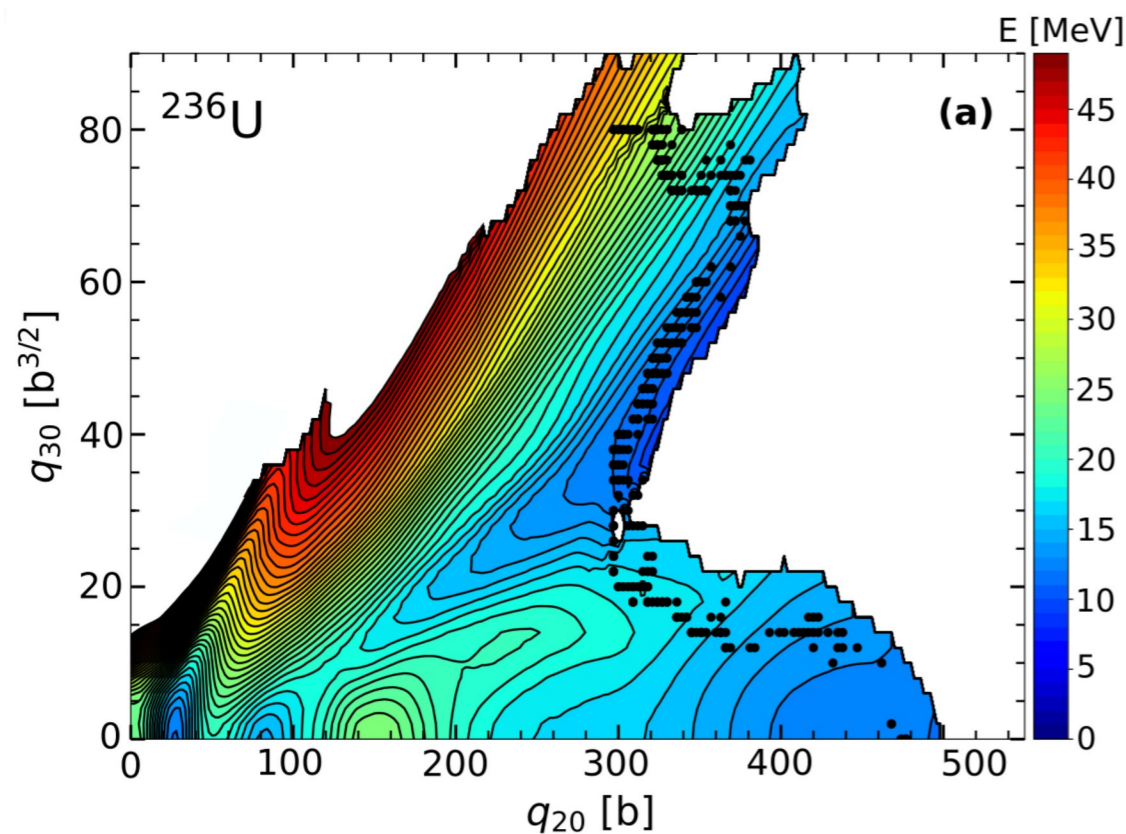
Potential Energy Surface and Scission Configurations

- **Neutron-induced fission**: even- Z and odd- N nucleus absorbs a neutron
- HFB model determines the **potential energy surface** of a fissioning system that contains **scission configurations**
 - Skyrme EDF and zero-range pairing
 - Wfs expanded in the HO basis
 - Axial and time-reversal symmetry assumed
 - Constraints imposed on q_{20} , q_{30} , and q_N values
 - Neck position is the isoscalar density minimum



HFBTHO program

P. Marević *et al.*, CPC 276 (2022) 108367



Theoretical Framework

Quantum Number Distributions in Scission Configurations

- Quantum number distributions for both FFs in each scission configuration are extracted through **projection techniques**

$$\mathbb{P}_F(J_F, N_F, Z_F, N, Z|q) = \langle \Phi_q^S | \hat{P}^{J_F} \hat{P}_{N_F, Z_F}^{N, Z} | \Phi_q^S \rangle$$

- Good particle number in the total system (N, Z)
- Good particle number in FFs (N_F, Z_F)
- Good angular momentum in FFs (J_F)

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- Good particle number in the total system (N, Z)
 - Good particle number in FFs (N_F, Z_F)
 - Good angular momentum in FFs (J_F)
- In practice, we use 1D-AMP and Fomenko expansion for PNP

$$\begin{aligned} \mathbb{P}_F(J_F, N_F, Z_F, N, Z|q) &= \frac{2J_F + 1}{2} \frac{1}{N_\varphi^4} \int_0^\pi d\beta \sin \beta d_{00}^{J_F}(\beta) \\ &\times \sum_{l_N, l_Z=1}^{N_\varphi} e^{-iN_F \varphi_{l_N}} e^{-iZ_F \varphi_{l_Z}} \sum_{k_N, k_Z=1}^{N_\varphi} e^{-iN \tilde{\varphi}_{k_N}} e^{-iZ \tilde{\varphi}_{k_Z}} \\ &\times \mathcal{N}_q^F(\beta, \varphi_{l_N}, \varphi_{l_Z}, \tilde{\varphi}_{k_N}, \tilde{\varphi}_{k_Z}) \end{aligned}$$

Theoretical Framework

Quantum Number Distributions in Scission Configurations

- Norm kernel = overlap between the initial and rotated state

$$\mathcal{N}_q^F(\beta, \vec{\varphi}) = \langle \Phi_q^S | \hat{R}_y^F(\beta) \hat{R}_{N_F, Z_F}^{N, Z}(\vec{\varphi}) | \Phi_q^S \rangle$$

- Overlaps are separable in isospin

$$\mathcal{N}_q^F(\beta, \vec{\varphi}) = \mathcal{N}_q^{F(\tau=n)}(\beta, \varphi_{l_N}, \tilde{\varphi}_{k_N}) \times \mathcal{N}_q^{F(\tau=p)}(\beta, \varphi_{l_Z}, \tilde{\varphi}_{k_Z})$$

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- They are calculated with Onishi formula (taking care of non-unitarity)

$$\mathcal{N}_q^{F(\tau)}(\beta, \varphi, \tilde{\varphi}) = \sqrt{\det[A_q^{F(\tau)}(\beta, \varphi, \tilde{\varphi})] \det[\mathcal{R}^F(\beta, \varphi, \tilde{\varphi})]}$$

$$A_q^{F(\tau)}(\beta, \varphi, \tilde{\varphi}) = U_q^{(\tau)T} [(\mathcal{R}^F(\beta, \varphi, \tilde{\varphi}))^T]^{-1} U_q^{(\tau)*} \\ + V_q^{(\tau)T} \mathcal{R}^F(\beta, \varphi, \tilde{\varphi}) V_q^{(\tau)*}$$

Theoretical Framework

Quantum Number Distributions in Scission Configurations

- Finally we extract the component with good total number of particles

$$\mathbb{P}_F(J_F, N_F, Z_F | N_0, Z_0, q) = \langle \Phi_q^S | \hat{P}^{J_F} \hat{P}_{N_F, Z_F}^{N_0, Z_0} | \Phi_q^S \rangle / \langle \Phi_q^S | \hat{P}^{N_0} \hat{P}^{Z_0} | \Phi_q^S \rangle$$

- The procedure is repeated for all configurations (hundreds of them)

Theoretical Framework

Quantum Number Distributions in Scission Configurations

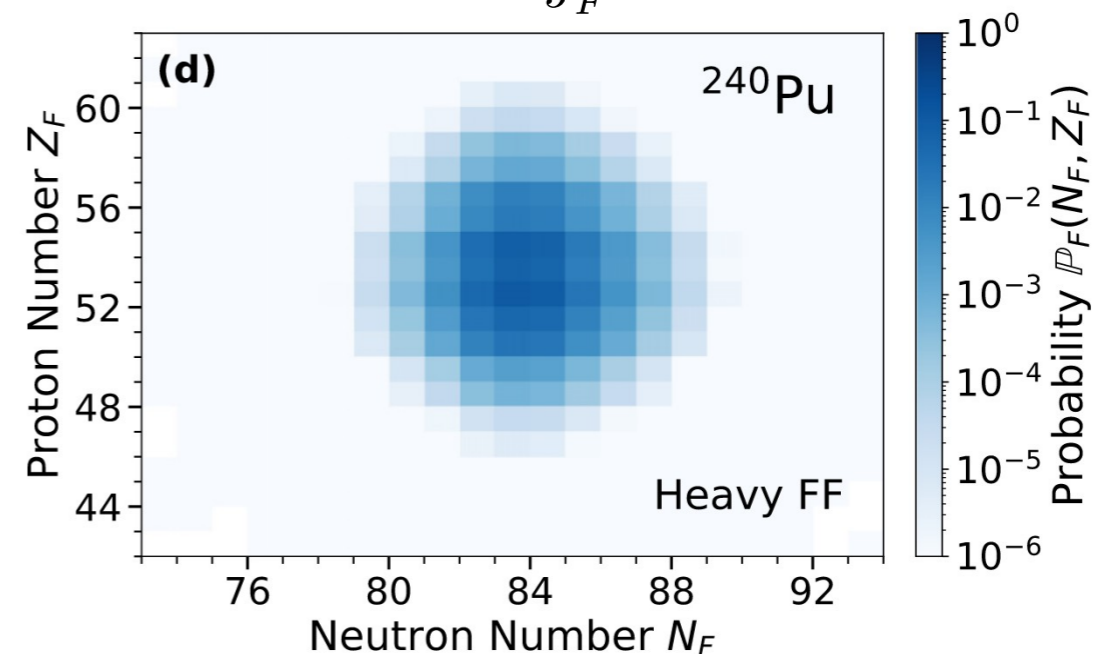
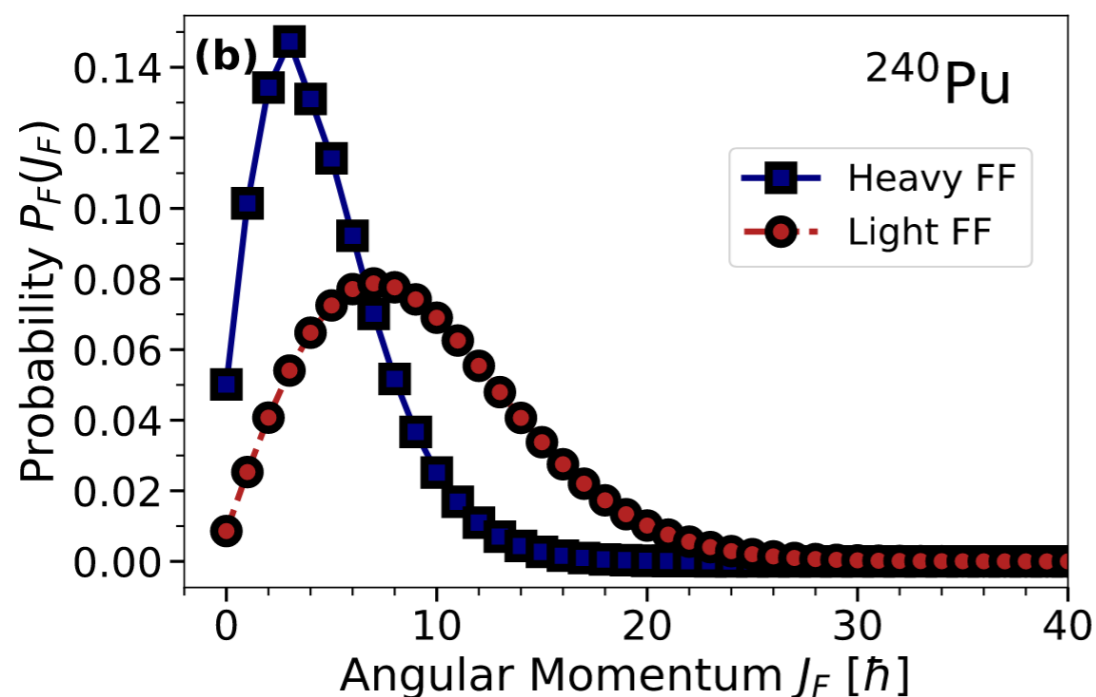
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- The procedure is repeated for all configurations (hundreds of them)
- Each scission configuration is characterized by a **full distribution**

Distributions near the most likely fragmentation in $^{240}\text{Pu}^*$

$$\mathbb{P}_F(J_F | q) = \sum_{N_F, Z_F} \mathbb{P}_F(J_F, N_F, Z_F | q) \quad \mathbb{P}_F(N_F, Z_F | q) = \sum_{J_F} \mathbb{P}_F(J_F, N_F, Z_F | q)$$



Theoretical Framework

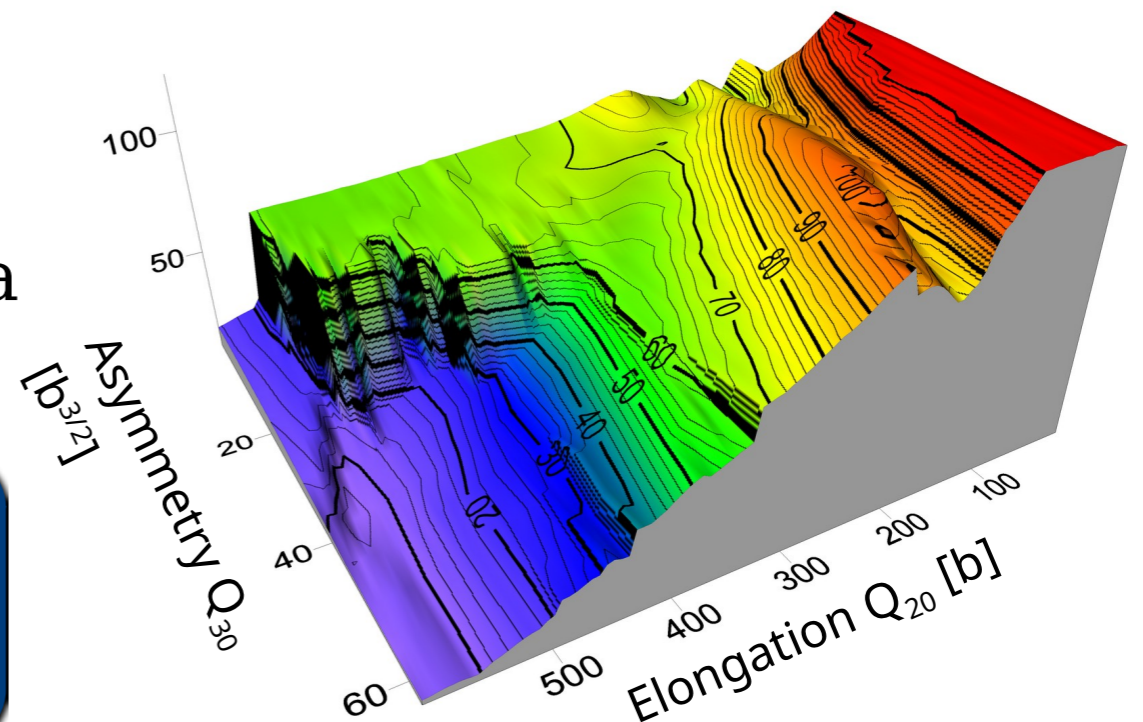
Probability of Populating Scission Configurations

- **Population probability** for each conf. is estimated with **TDGCM+GOA**
- Time-dependent wave function is a superposition of HFB states

$$|\Psi(t)\rangle = \int dq f_q(t) |\Phi_q\rangle$$

- Variational principle and GOA lead to a Schrödinger-like equation

$$i\hbar \frac{\partial}{\partial t} g_q(t) = [\mathcal{H}_q^{coll} + i\mathcal{A}_q^{coll}] g_q(t)$$



- Probability that the wave packet exits through any particular point is proportional to the time-integrated flux density

$$F(q) = \lim_{t \rightarrow \infty} \int_{\tau=0}^{\tau=t} d\tau \phi(q, \tau)$$



FELIX Finite Element Solver
D. Regnier *et al.*, CPC 225 (2018) 180-191

Theoretical Framework

Final Distribution

- Final distribution is obtained by folding the two results

$$\mathbb{P}_F(J_F, N_F, Z_F | N_0, Z_0) = \int dq F(q) \mathbb{P}_F(J_F, N_F, Z_F | N_0, Z_0, q)$$

$$\mathbb{P}(J_F, N_F, Z_F | N_0, Z_0) = \mathbb{P}_l(J_F, N_F, Z_F | N_0, Z_0) + \mathbb{P}_r(J_F, N_F, Z_F | N_0, Z_0)$$

- Distribution covers the full range of N_F and Z_F
- Fixing (N_F, Z_F) and renormalizing gives the J_F distribution in particular FF
- Marginalization over J_F gives pre-neutron mass and charge yields
- Initial conditions partially simulate different incident neutron energies

$$\mathbb{P}(J_F, N_F, Z_F | N_0, Z_0) \rightarrow \mathbb{P}(J_F, N_F, Z_F | N_0, Z_0, E_n)$$

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- Some limitations of the current model:
 - Only $K = 0$ components in AMP
 - No intrinsic excitations, only deformation energy
 - Influence of relative motion disregarded
 - No joint projection in the two FFs

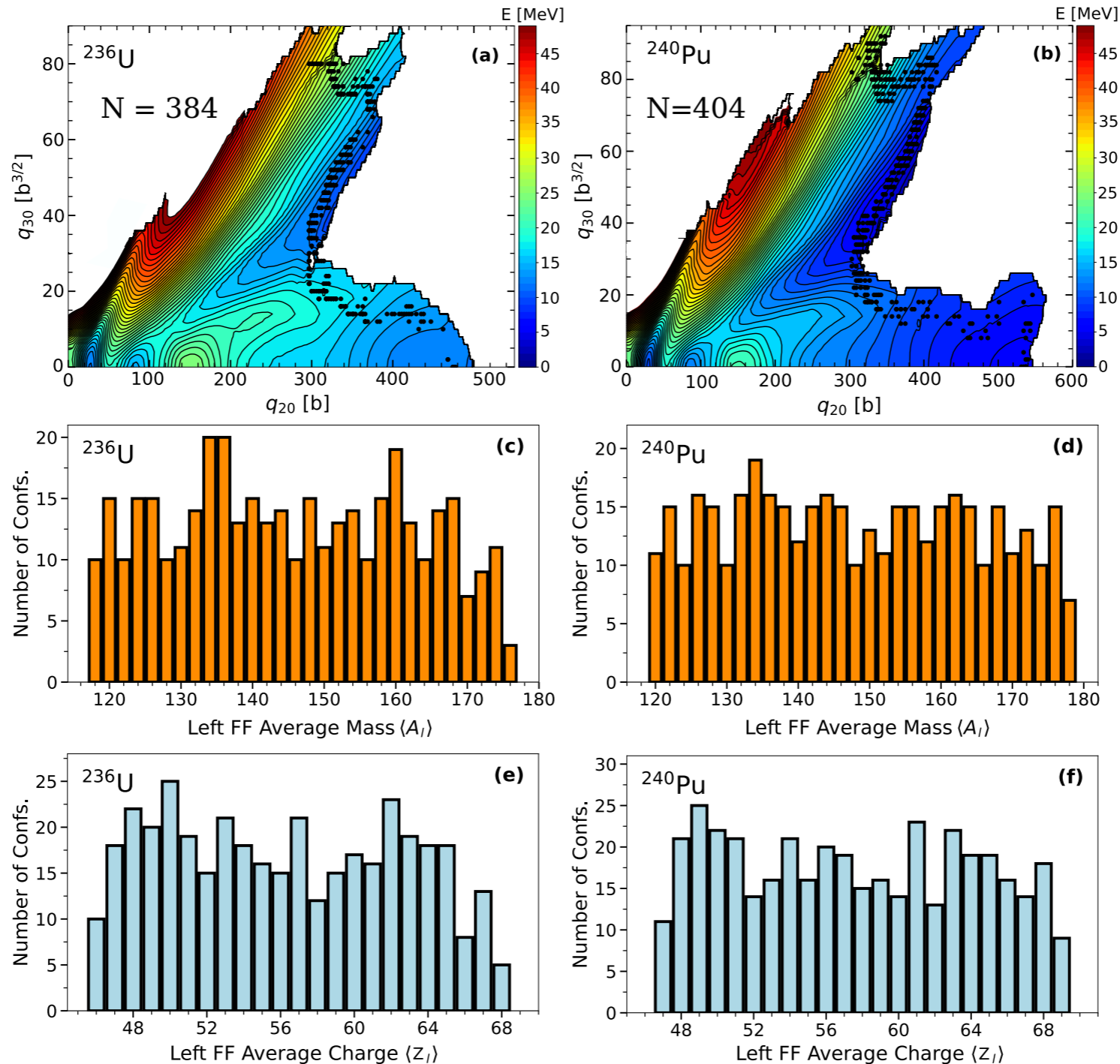
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n-induced Fission of ^{235}U and ^{239}Pu

Properties of Scission Configurations

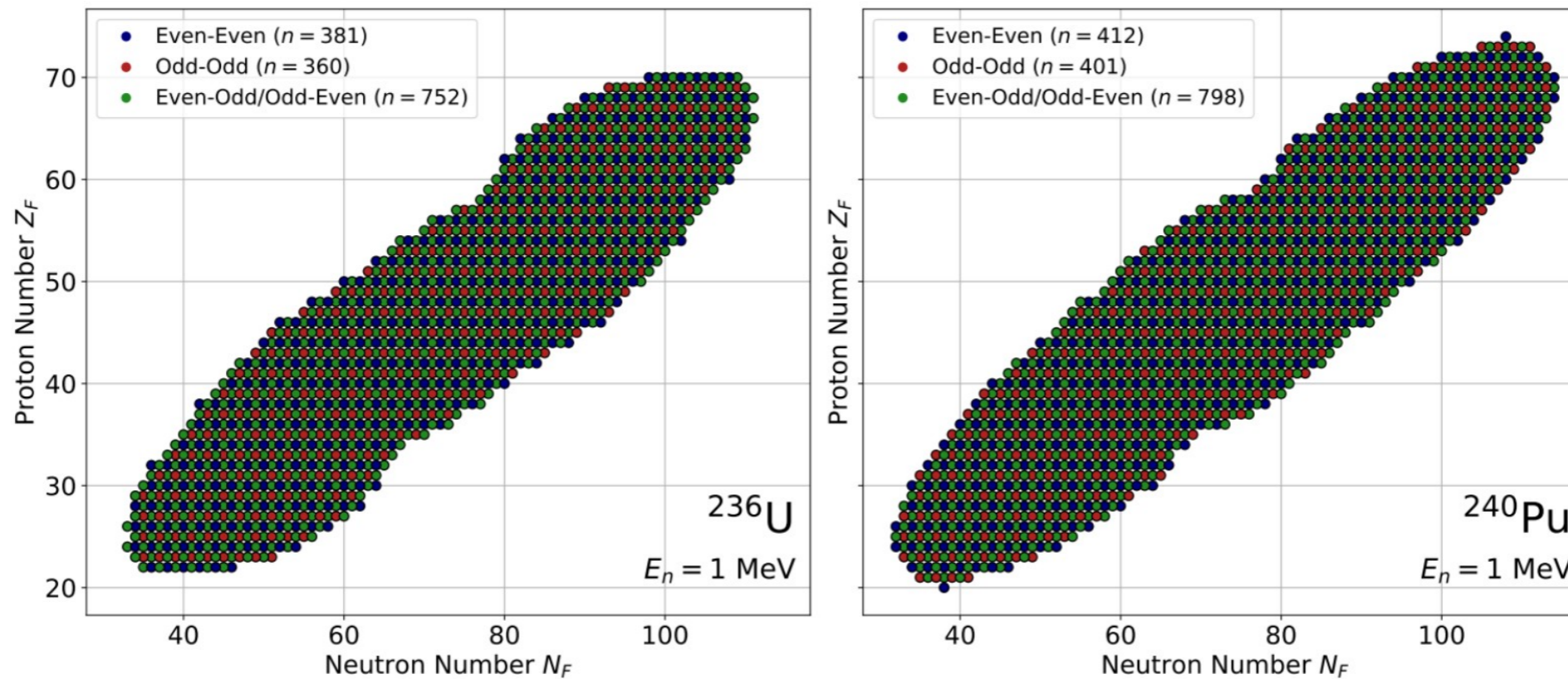
- We consider a very wide range of scission configurations



n-induced Fission of ^{235}U and ^{239}Pu

Range of Distributions

- First microscopic prediction of J_F distributions for the **full range of FFs**



PHYSICAL REVIEW C **113**, 014612 (2026)

Editors' Suggestion

Microscopic theory of angular momentum distributions across the full range of fission fragments

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Nuclear data and Theory Group, Nuclear and Chemical Science Division, [Lawrence Livermore National Laboratory](#), Livermore, California 94550, USA

Marc Verriere[‡]

Nuclear data and Theory Group, Nuclear and Chemical Science Division, [Lawrence Livermore National Laboratory](#), Livermore, California 94550, USA
and CEA, DAM, DIF, 91297 Arpajon, France

Microscopic Angular Momentum Distributions in Fragments for Neutron-Induced Fission of U-235 and Pu-239

<https://zenodo.org/records/17303186>

This dataset provides microscopic angular momentum distributions in even-even fission fragments for neutron-induced fission of ^{235}U and ^{239}Pu .

It was produced within the "Angular momentum predictions for low-energy induced fission" (AMPLIFI) project, funded by the European Union's Horizon Europe research and innovation programme under the Marie Skłodowska-Curie Actions Grant Agreement No. 101149053.

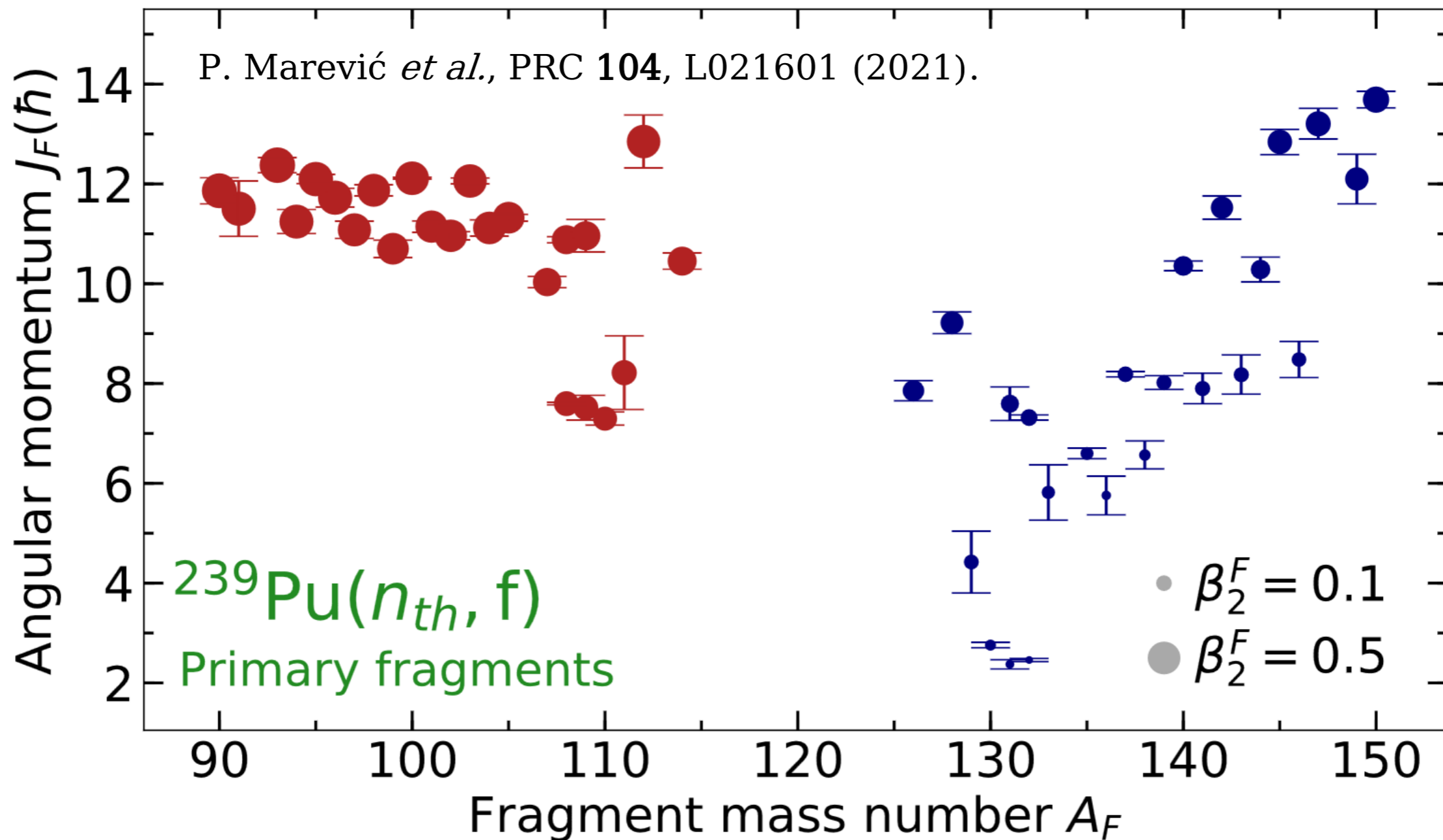
Reference: P. Marević, N. Schunck, M. Verriere, "Microscopic theory of angular momentum distributions across the full range of fission fragments", *Phys. Rev. C* **113**, 014612 (2026).

For a full description of the framework, see the [publication](#) or the [arXiv preprint](#).

n-induced Fission of ^{235}U and ^{239}Pu

The Sawtooth Pattern

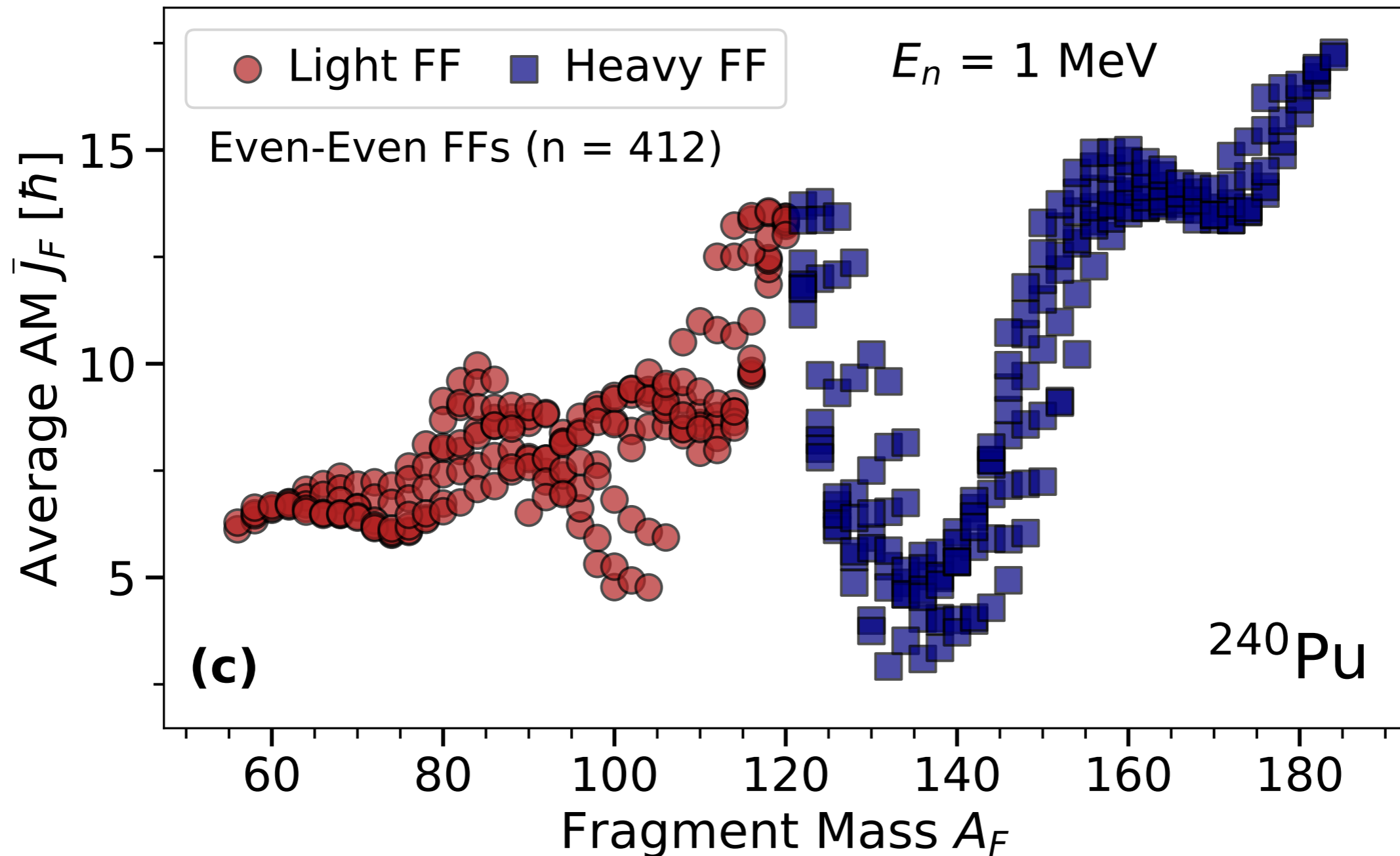
- First analysis (2021) was consistent with a sawtooth pattern



n-induced Fission of ^{235}U and ^{239}Pu

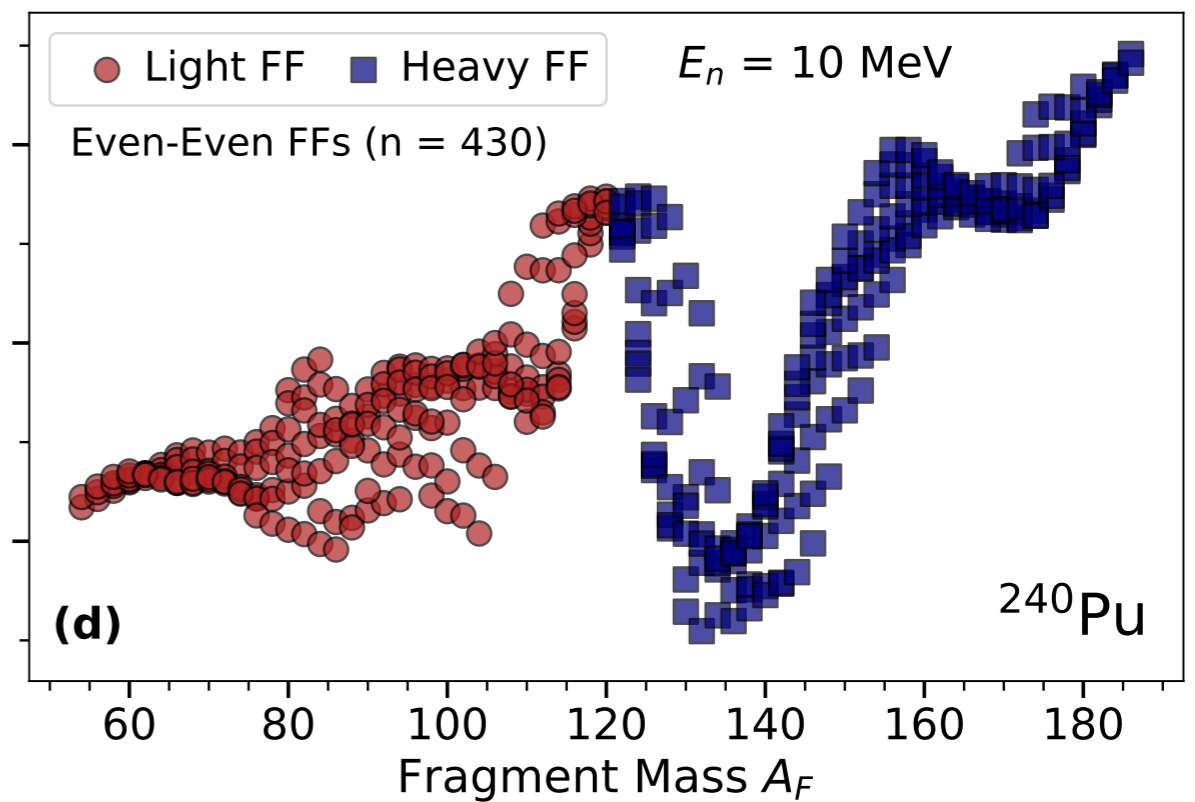
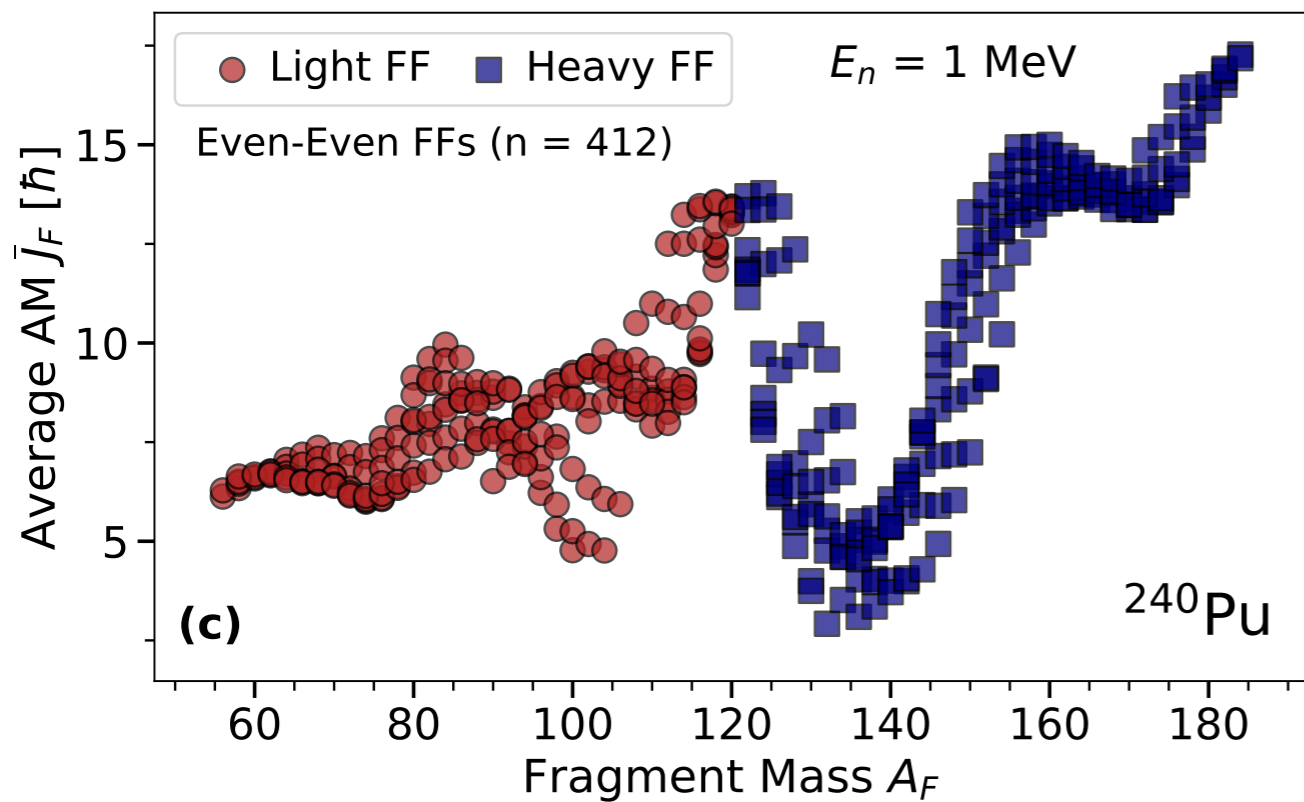
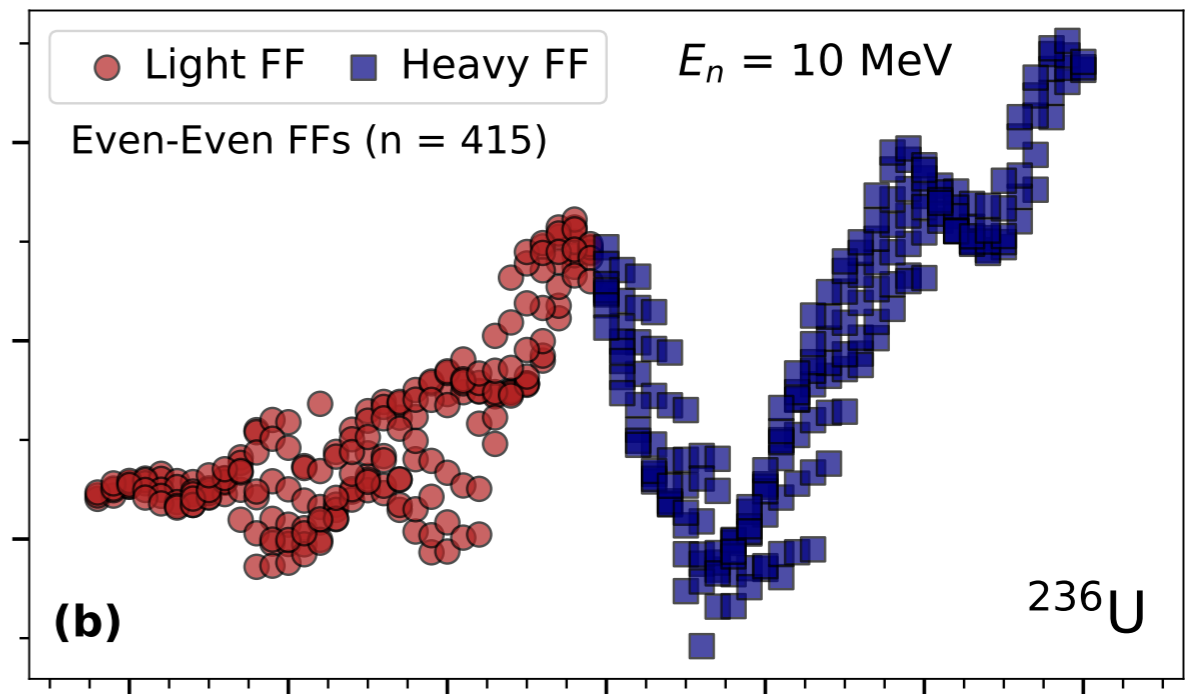
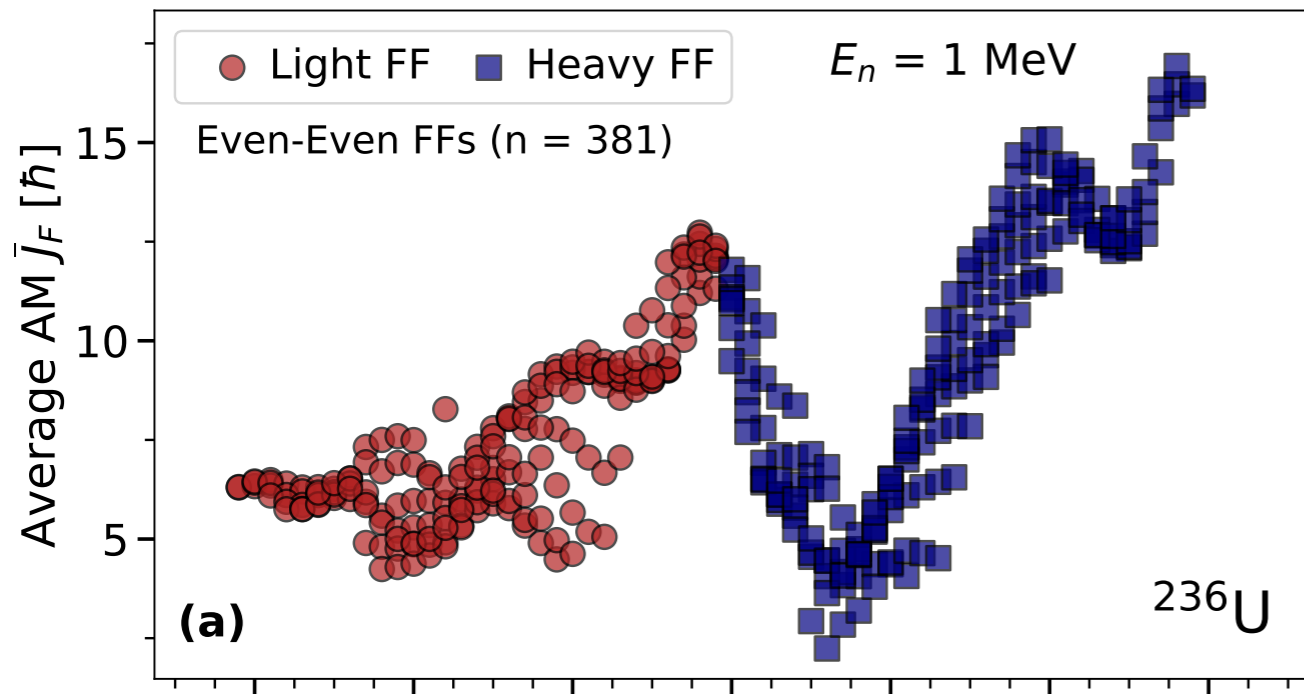
The Sawtooth Pattern

- Unequivocal microscopic evidence of a universal sawtooth pattern



n-induced Fission of ^{235}U and ^{239}Pu

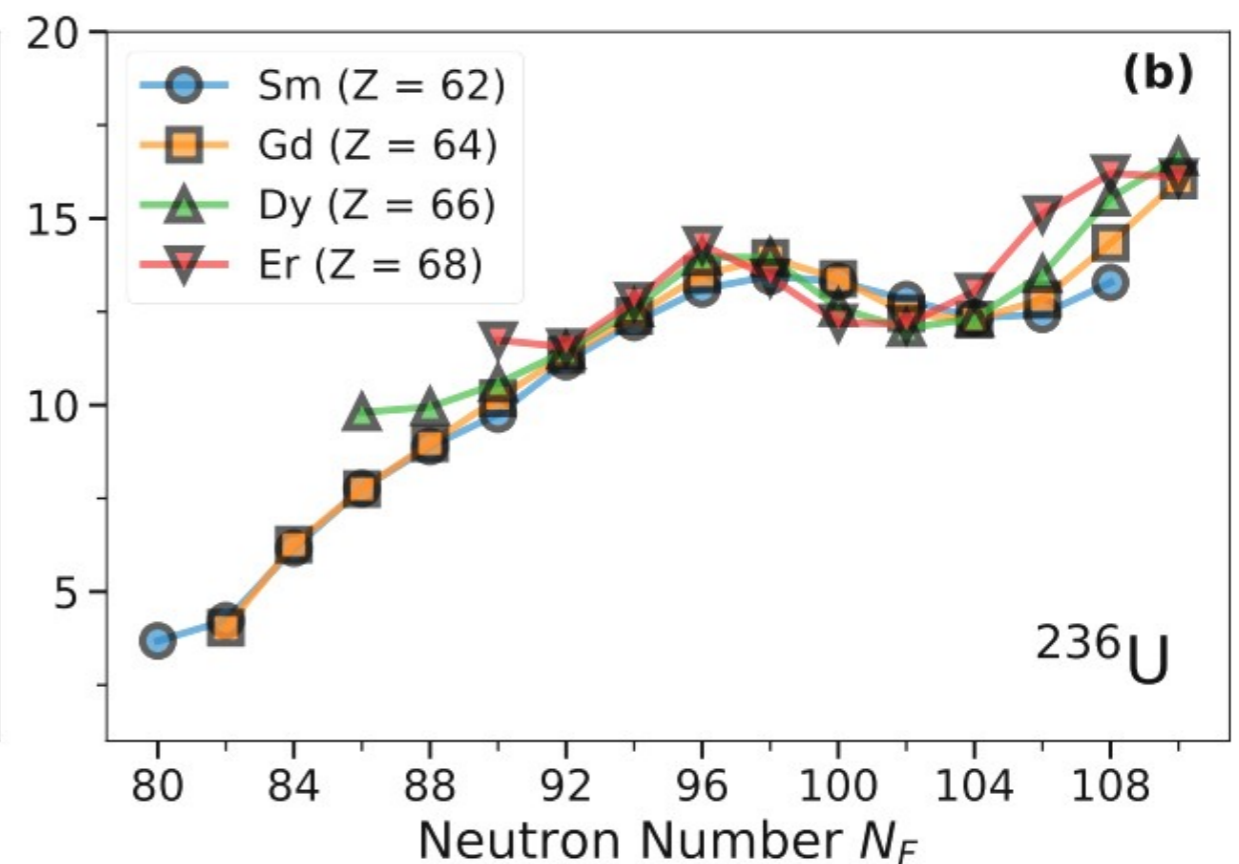
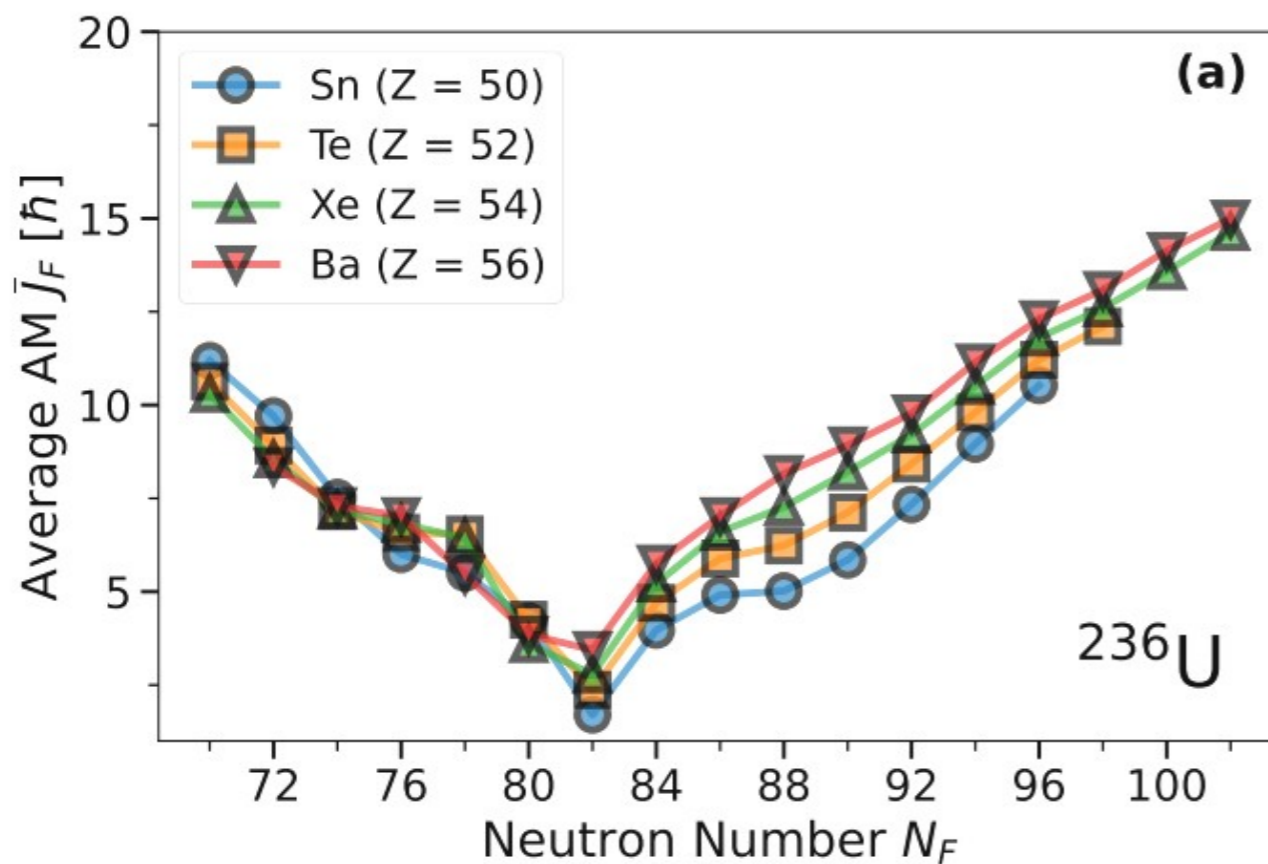
The Sawtooth Pattern



n-induced Fission of ^{235}U and ^{239}Pu

Shell Effects

- Shell effects hinder the ability of FFs to carry angular momentum
 - N=82 shell closure has a decisive impact
 - Local minimum in the N = 100-104 region → deformed shell closure?



n-induced Fission of ^{235}U and ^{239}Pu

Isobaric Dependence

Phenomenological formula:

$$p(J_F) \propto (2J_F + 1) \exp \left(- \frac{1}{2} \frac{J_F(J_F + 1)}{B^2(Z_F, A_F, T_F)} \right)$$

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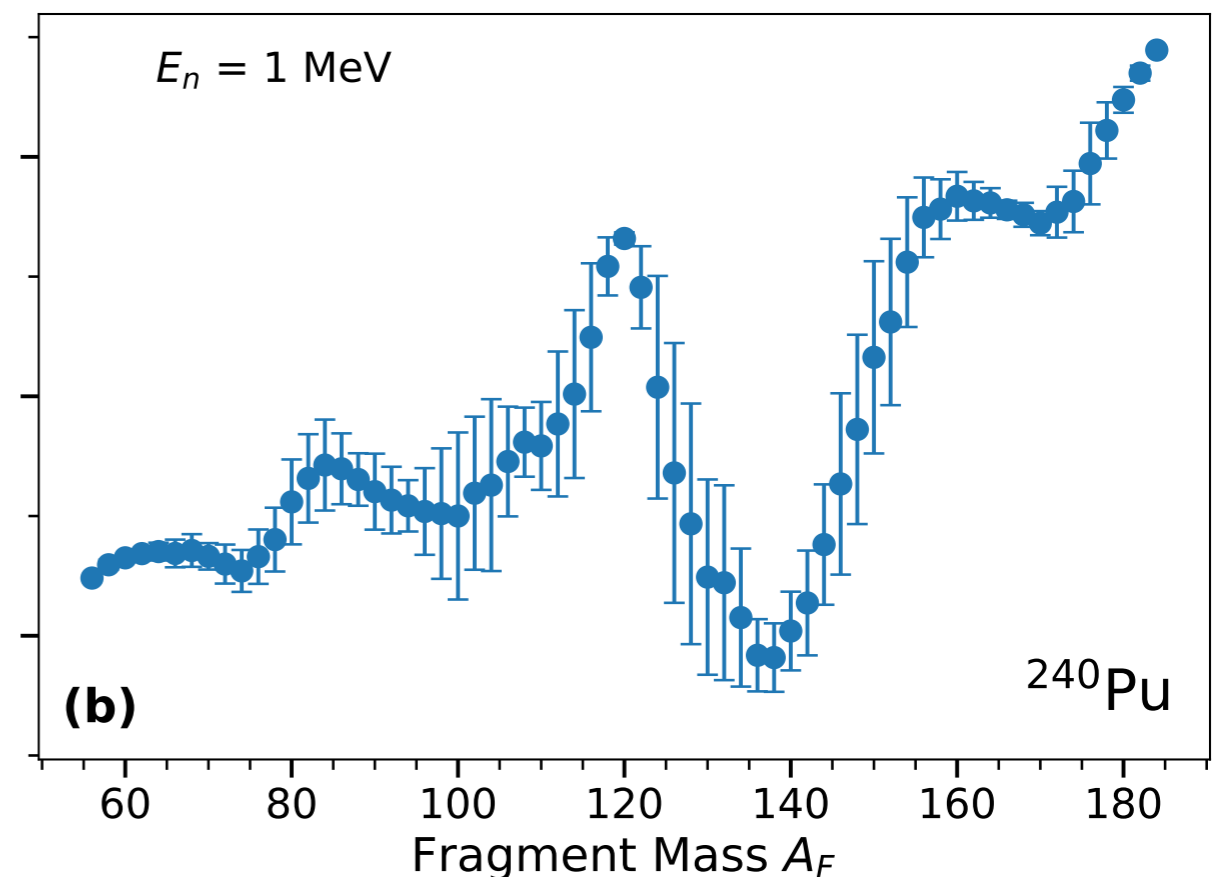
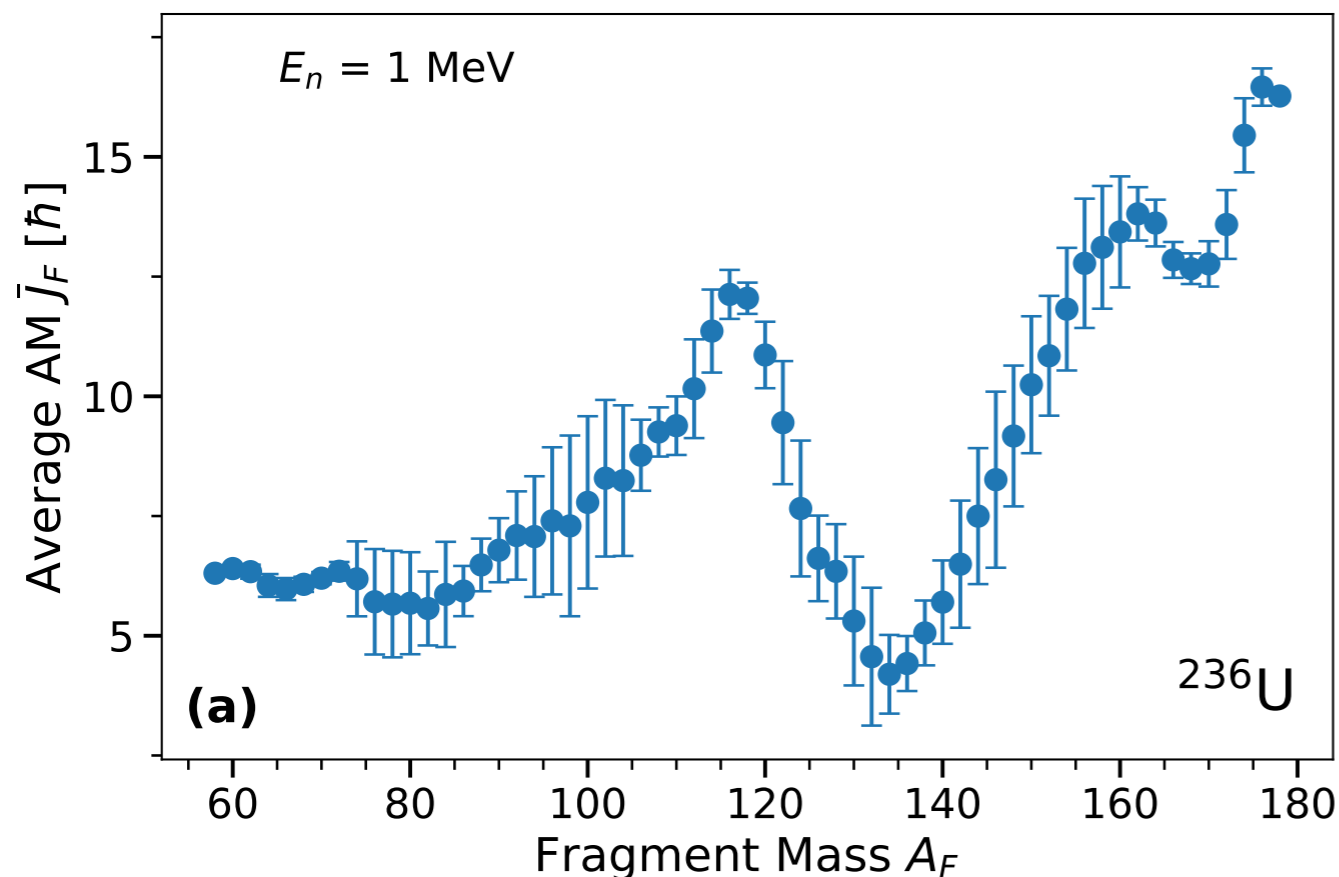
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- There is a strong **isobaric dependence** of AM distributions



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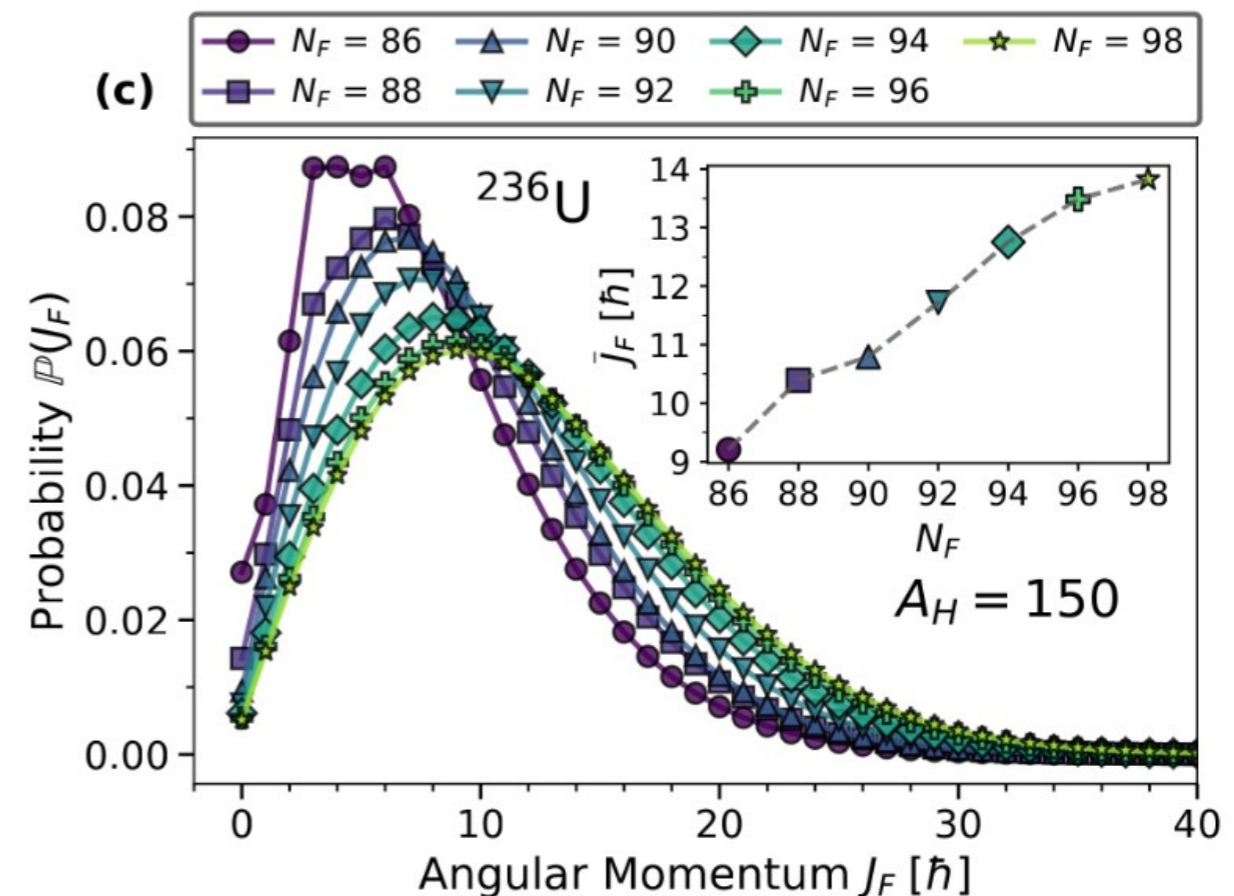
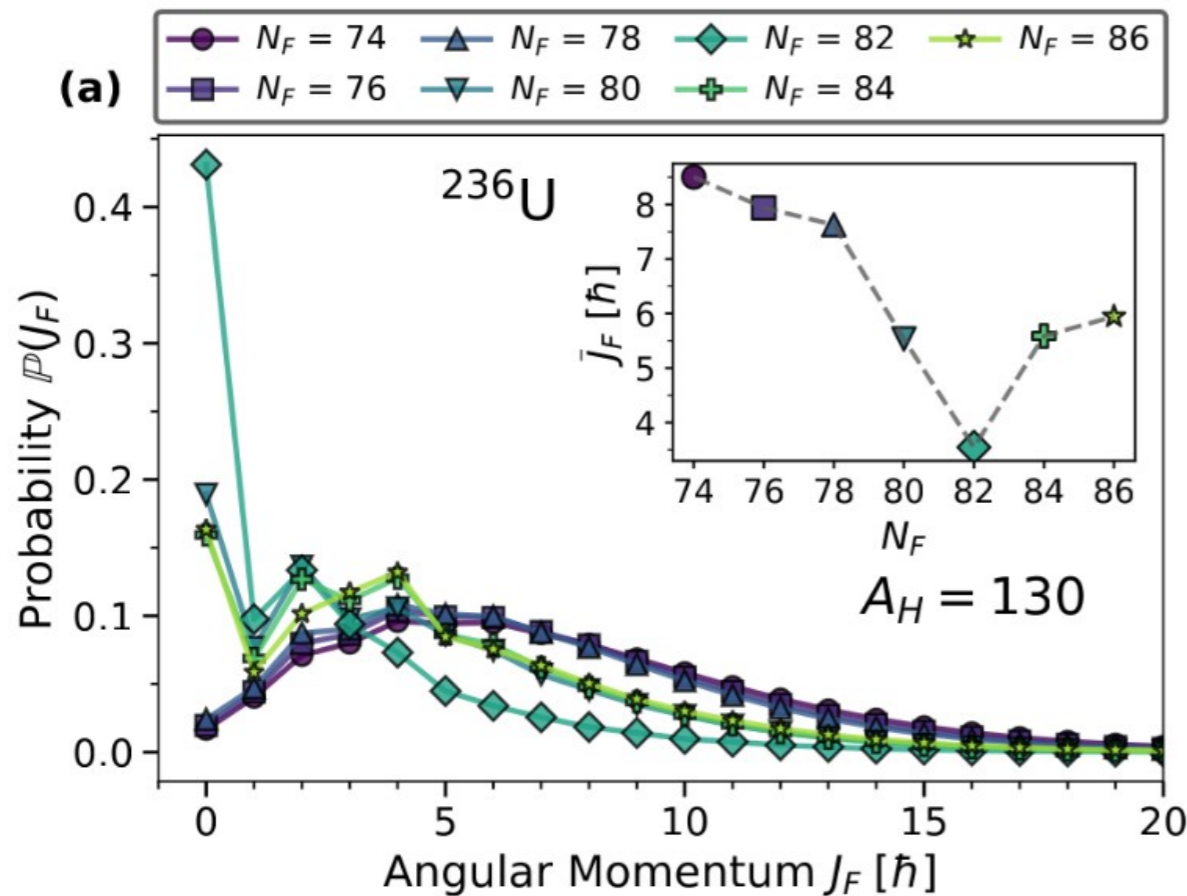
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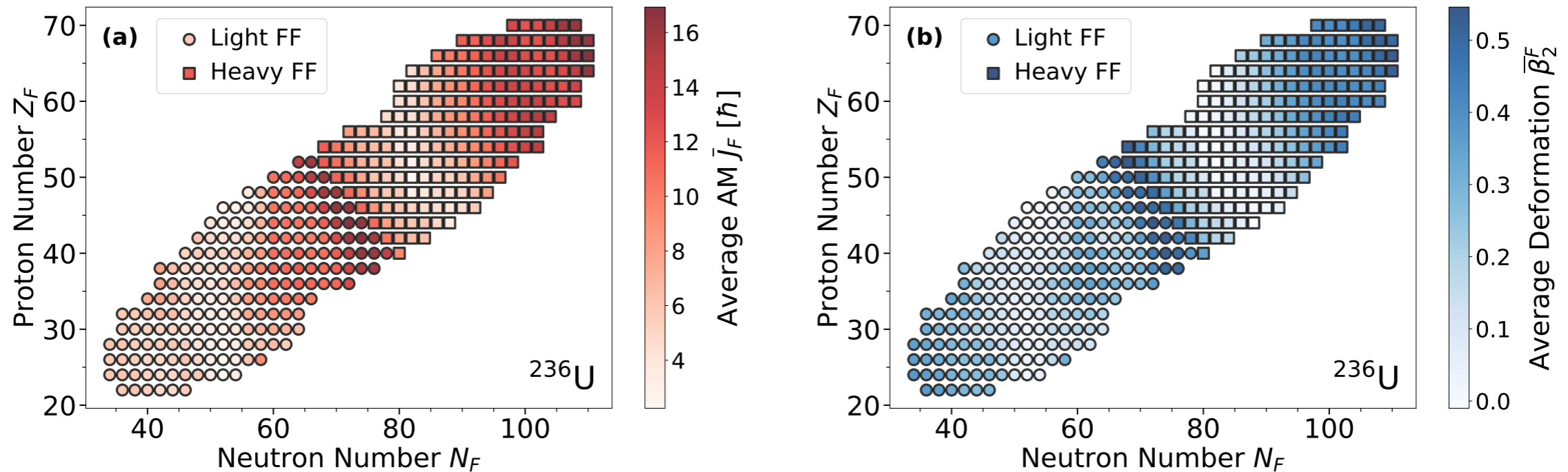
- There is a strong **isobaric dependence** of AM distributions



n-induced Fission of ^{235}U and ^{239}Pu

Fragment Deformations

- FF deformations and AM are strongly correlated ($\rho_{\text{Pear}} = 0.67$ for ^{236}U)



$$\bar{\beta}_2^F(N_F, Z_F) = \frac{\int dq F(q) \beta_2(N_F, Z_F, q)}{\int dq F(q) s(N_F, Z_F, q)}$$

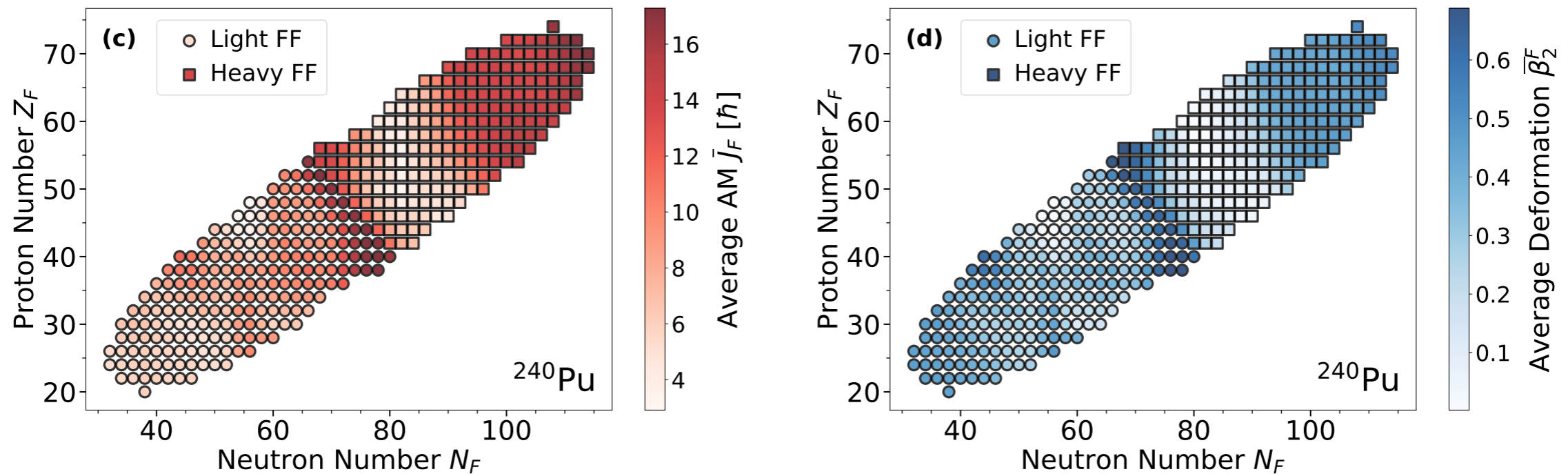
$$\beta_2(N_F, Z_F, q) = \sum_{f=l,r} \mathbb{P}_f(N_F, Z_F | N_0, Z_0, q) \beta_2^f(q)$$

$$s(N_F, Z_F, q) = \sum_{f=l,r} \mathbb{P}_f(N_F, Z_F | N_0, Z_0, q)$$

n-induced Fission of ^{235}U and ^{239}Pu

Fragment Deformations

- FF deformations and AM are strongly correlated ($\rho_{\text{Pear}} = 0.64$ for ^{240}Pu)



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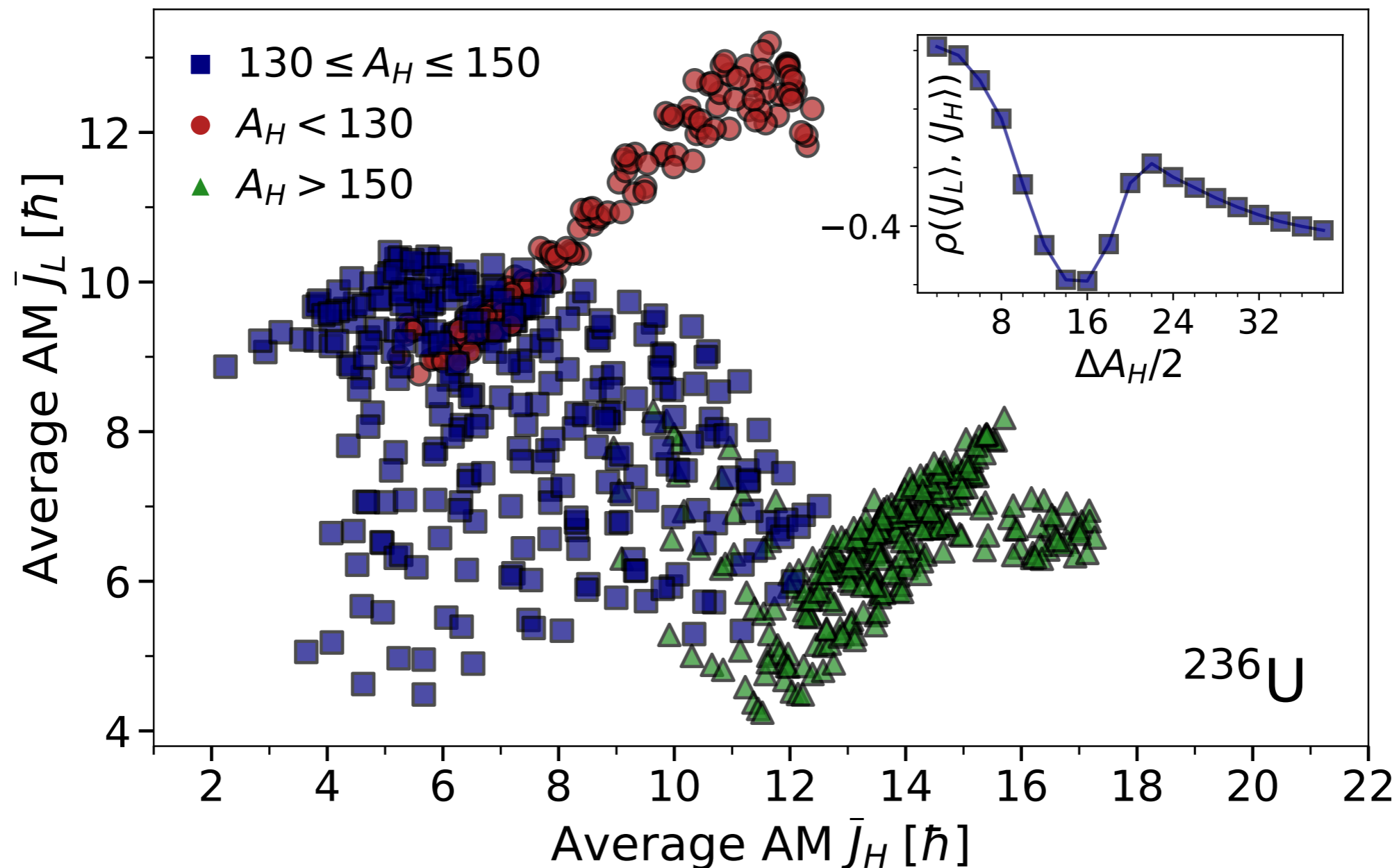
$$\beta_2(N_F, Z_F, q) = \sum_{f=l,r} \mathbb{P}_f(N_F, Z_F | N_0, Z_0, q) \beta_2^f(q)$$

$$s(N_F, Z_F, q) = \sum_{f=l,r} \mathbb{P}_f(N_F, Z_F | N_0, Z_0, q)$$

n-induced Fission of ^{235}U and ^{239}Pu

Correlation in Magnitude of FF Angular Momenta

- Weak negative correlation ($\rho_{\text{Pear}} = -0.33$) in AM magnitudes for most strongly populated FFs



n-induced Fission of ^{235}U and ^{239}Pu

Effect on Fission Spectra

$$p(J_F) \propto (2J_F + 1) \exp\left(-\frac{1}{2} \frac{J_F(J_F + 1)}{B^2(Z_F, A_F, T_F)}\right)$$
$$B^2(Z_F, A_F, T_F) = \alpha \frac{\mathcal{I}_0(Z_F, A_F)}{\hbar^2} T$$

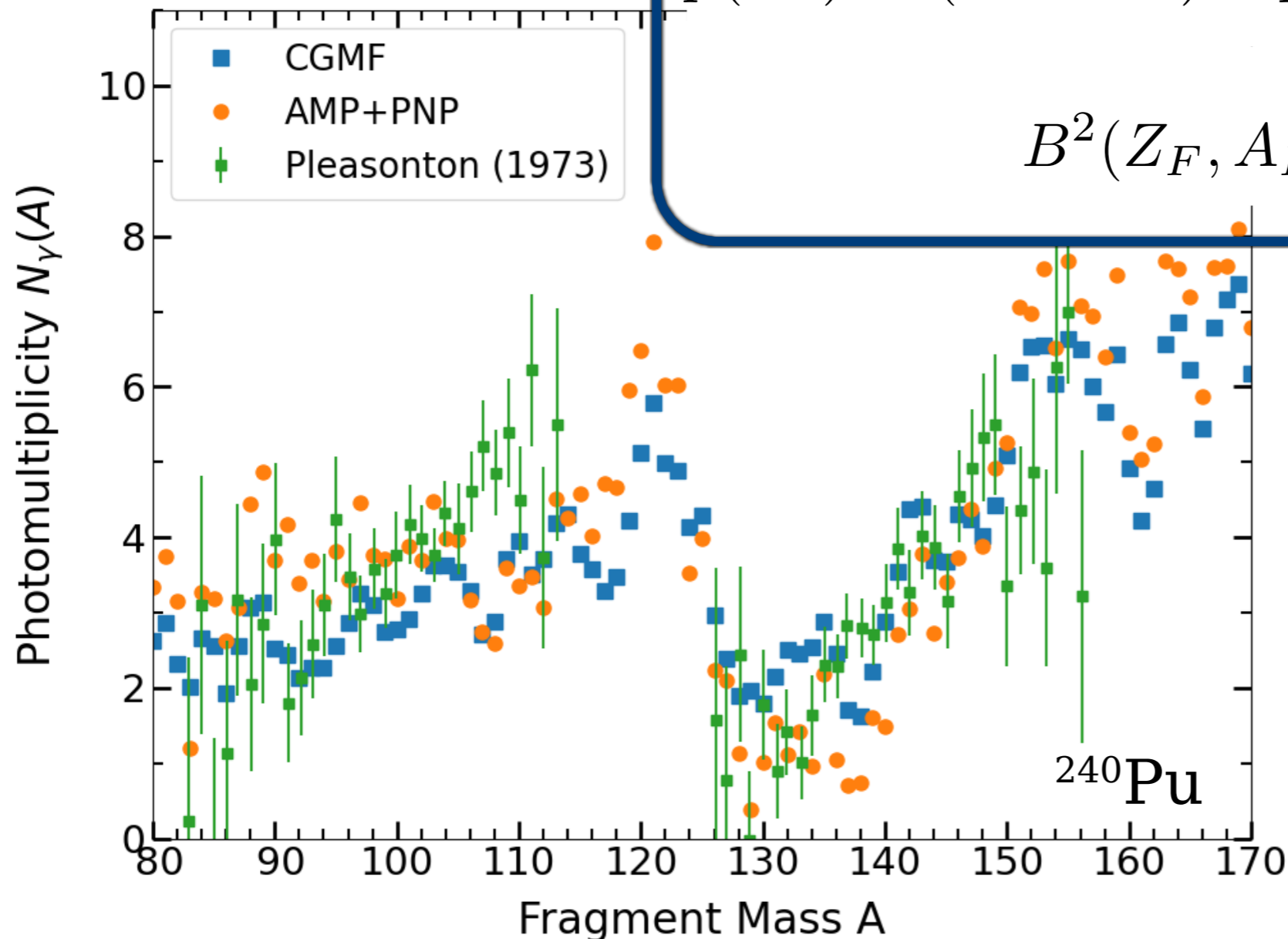
n-induced Fission of ^{235}U and ^{239}Pu

Effect on Fission Spectra

- Photon multiplicities are reproduced **without adjustable parameters** in AM distributions

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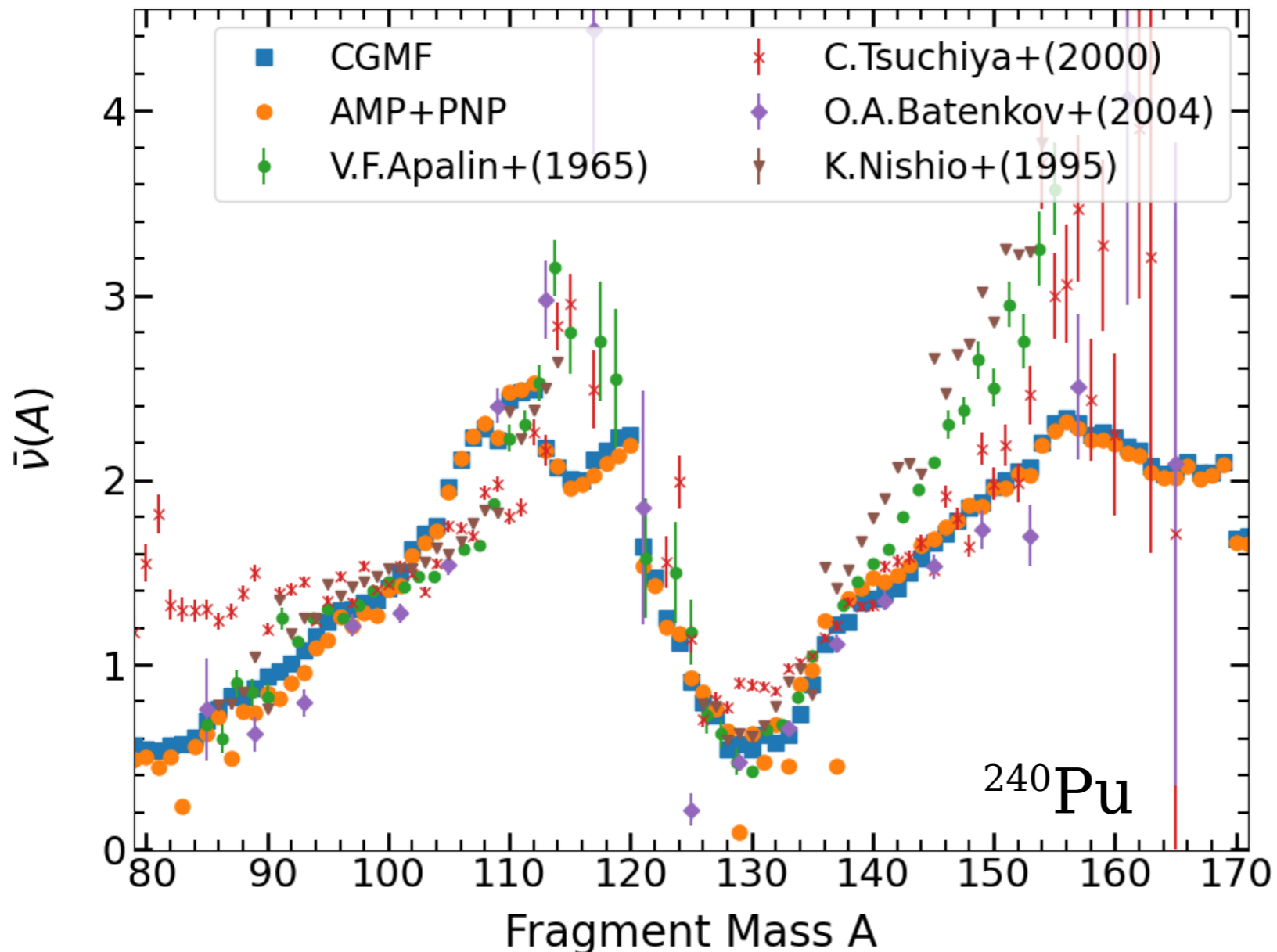
Photon multiplicity

$N_\gamma: 6.24 \rightarrow 6.08$

n-induced Fission of ^{235}U and ^{239}Pu

Effect on Fission Spectra

- Photon multiplicities are reproduced **without adjustable parameters** in AM distributions



Neutron multiplicity

$N_\nu: 2.86 \rightarrow 2.80$

Outline

1. Introduction

2. Theoretical Framework

3. n-Induced Fission of ^{235}U and ^{239}Pu

4. Conclusion

Conclusion

1. Microscopic fission theory is becoming quantitatively competitive with phenomenological models.

Publication: PM, N. Schunck, M. Verriere, “*Microscopic theory of angular momentum distributions across the full range of fission fragments*”, Phys. Rev. C **113**, 014612 (2026).

Database: “*Microscopic Angular Momentum Distributions in Fragments for Neutron-Induced Fission of U-235 and Pu-239*”, <https://zenodo.org/records/17303186>



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Conclusion

1. Microscopic fission theory is becoming quantitatively competitive with phenomenological models.

2. We can now predict microscopic AM distributions for all FFs.

- ✓ Sawtooth pattern
- ✓ Isobaric dependence
- ✓ Photon and neutron multiplicities

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Conclusion

1. Microscopic fission theory is becoming quantitatively competitive with phenomenological models.

2. We can now predict microscopic AM distributions for all FFs.

- ✓ Sawtooth pattern
- ✓ Isobaric dependence
- ✓ Photon and neutron multiplicities

3. Exciting developments are underway.

- Refinement and extension of AM models
- Predicting other ingredients for FFs decay modeling

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