

Halo nuclei from *ab initio* nuclear theory

International Symposium Commemorating the 40th
Anniversary of the Halo Nuclei (HALO-40)

Beijing, October 12-18, 2025

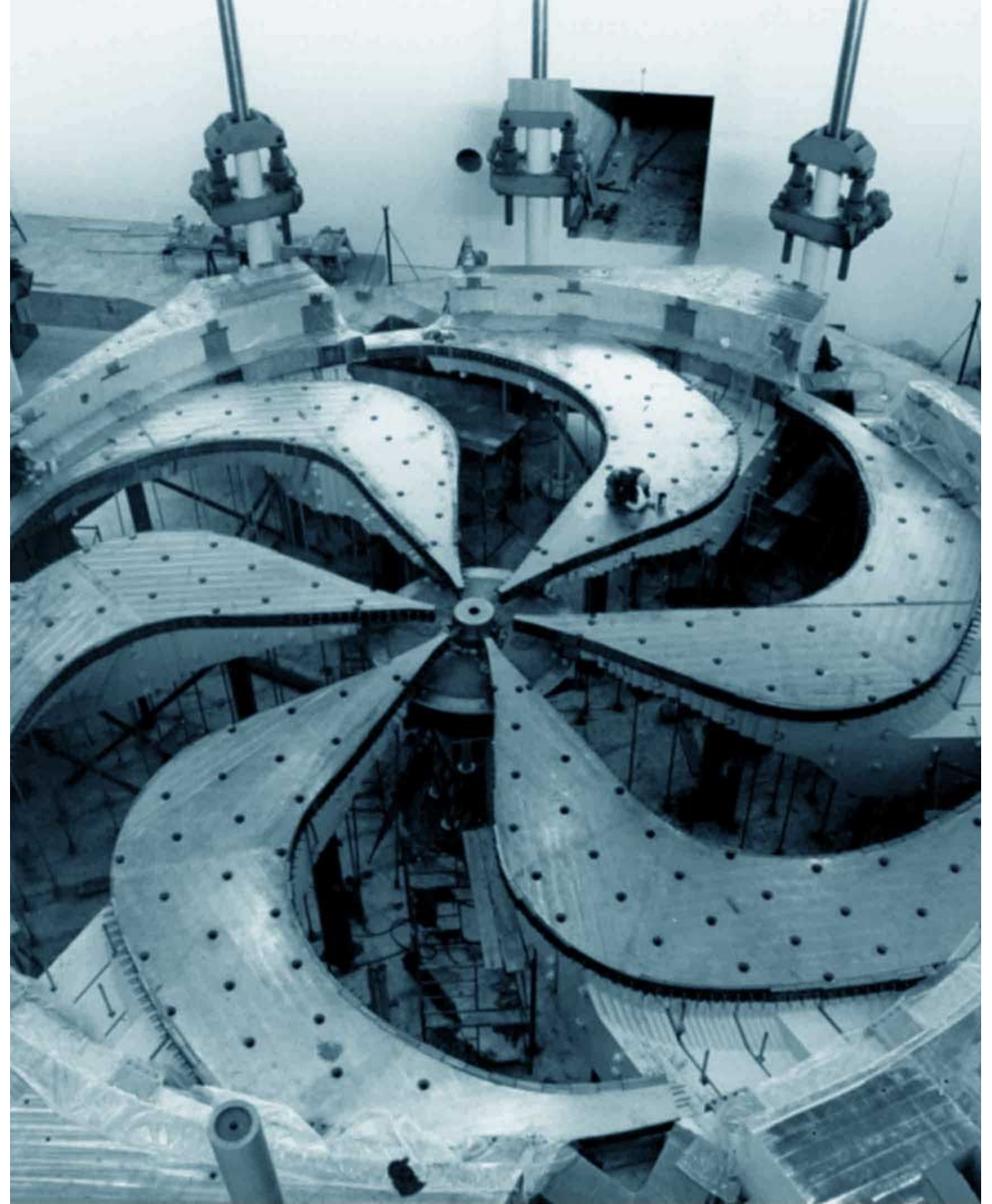
Petr Navratil

TRIUMF

Collaborators:

Sofia Quaglioni (LLNL), Kostas Kravvaris (LLNL), Guillaume
Hupin (IJCLab), Michael Gennari (Mainz),

2025-10-24

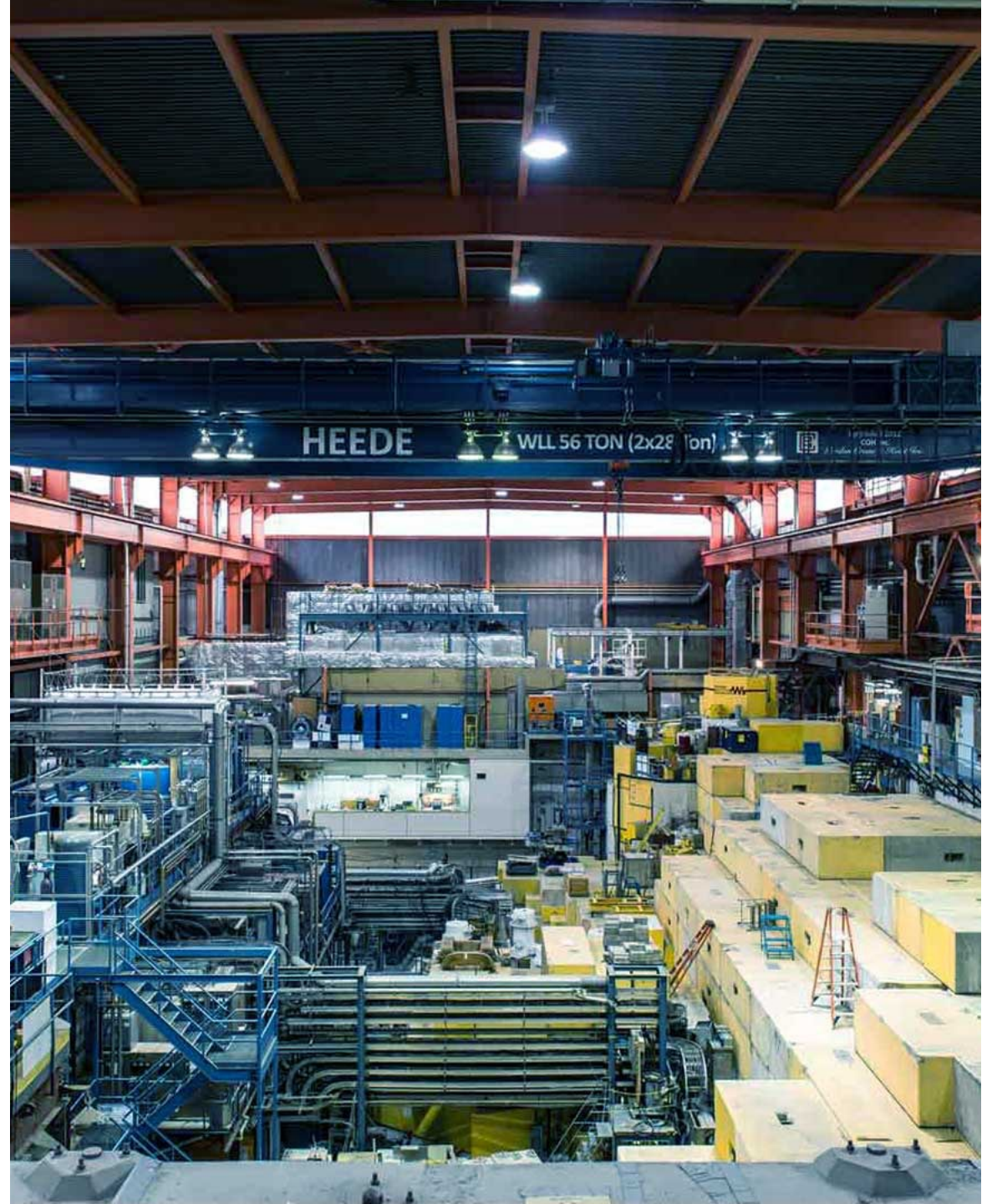


Outline

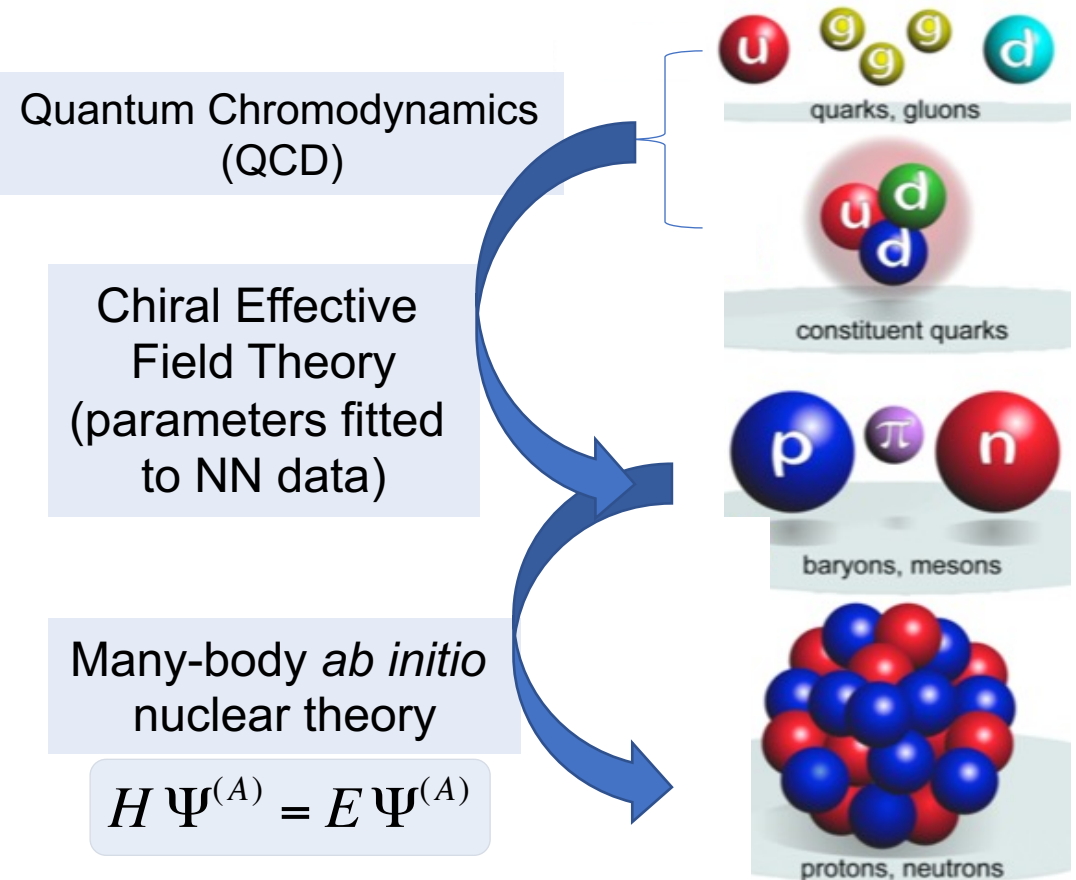
- Introduction – *Ab initio* nuclear theory – no-core shell model with continuum (NCSMC)
- Parity inversion in ^{11}Be ground state
- Halo ground state of ^{15}C
- *P*-wave halo ground state of ^8B & radiative capture of protons on ^7Be
- NCSMC for three-body cluster – $^6\text{He} \sim ^4\text{He} + n + n$
- ^{11}Li within NCSM (prerequisite for $^{11}\text{Li} \sim ^9\text{Li} + n + n$)
- Conclusions

Ab initio nuclear theory -
no-core shell model with
continuum (NCSMC)

2025-10-24



First principles or *ab initio* nuclear theory



	NN force	NNN force	NNNN force
Q^0_{LO}			
Q^2_{NLO}			
$Q^3_{N^2LO}$			
$Q^4_{N^3LO}$			
	+ ...	+ ...	+ ...



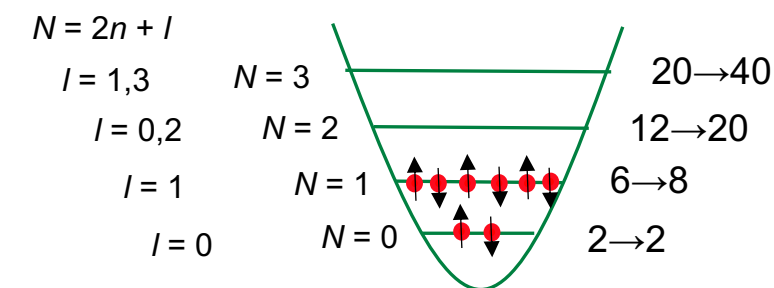
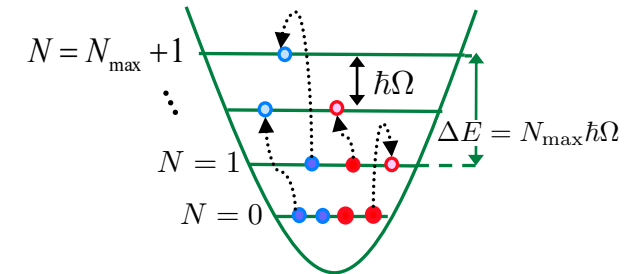
Review

Ab initio no core shell modelBruce R. Barrett^a, Petr Navrátil^b, James P. Vary^{c,*}

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Conceptually simplest *ab initio* method: No-Core Shell Model (NCSM)

- Basis expansion method (CI)
 - Harmonic oscillator (HO) basis truncated in a particular way ($N_{\max}$)
 - Why HO basis?
 - Lowest filled HO shells match magic numbers of light nuclei (2, 8, 20 – ^4He , ^{16}O , ^{40}Ca)
 - Equivalent description in relative(Jacobi)-coordinate and Slater determinant basis – nuclei self-bound, $[\mathbf{H}, \mathbf{P}_{\text{CM}}]=0$
 - Exact factorization of CM and intrinsic eigenfunctions at each $N_{\max}$
 - Good for well-bound states, approximation of weakly-bound states and narrow resonances



$$E = (2n + l + \frac{3}{2})\hbar\Omega$$

Ab initio calculations of structure including weakly bound halo states, scattering, reactions

Unified approach to bound & continuum states

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No-Core Shell Model with Continuum (NCSMC)

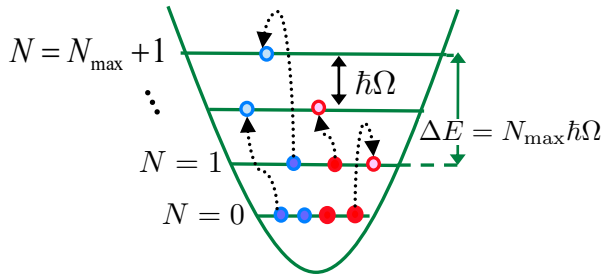
$$\Psi^{(A)} = \sum_{\lambda} c_{\lambda} \left| {}^{(A)} \text{cluster}, \lambda \right\rangle + \sum_{\nu} \int d\vec{r} \, \gamma_{\nu}(\vec{r}) \, \hat{A}_{\nu} \left| {}^{(A-a)} \text{cluster}, \nu \right\rangle$$

Ab initio calculations of structure including weakly bound halo states, scattering, reactions

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No-Core Shell Model with Continuum (NCSMC)

$$\Psi^{(A)} = \underbrace{\sum_{\lambda} c_{\lambda} \left| \begin{array}{c} (A) \\ \text{cluster} \end{array}, \lambda \right\rangle}_{\text{bound states}} + \sum_{\nu} \int d\vec{r} \gamma_{\nu}(\vec{r}) \hat{A}_{\nu} \left| \begin{array}{c} (A-a) \text{ cluster} \\ (a) \text{ particle} \end{array}, \nu \right\rangle$$


Static solutions for aggregate system,
describe all nucleons close together

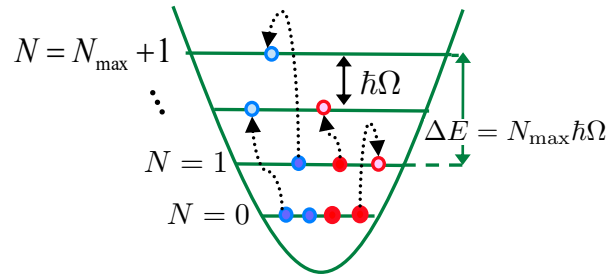
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No-Core Shell Model with Continuum (NCSMC)

$$\Psi^{(A)} = \underbrace{\sum_{\lambda} c_{\lambda} \left| \begin{array}{c} (A) \\ \text{cluster} \end{array}, \lambda \right\rangle}_{\text{Static solutions for aggregate system, describe all nucleons close together}} + \underbrace{\sum_{\nu} \int d\vec{r} \gamma_{\nu}(\vec{r}) \hat{A}_{\nu} \left| \begin{array}{c} (A-a) \\ \text{cluster} \end{array}, \nu \right\rangle}_{\text{Continuous microscopic cluster states, describe long-range projectile-target}}$$



Continuous microscopic cluster states,
describe long-range projectile-target

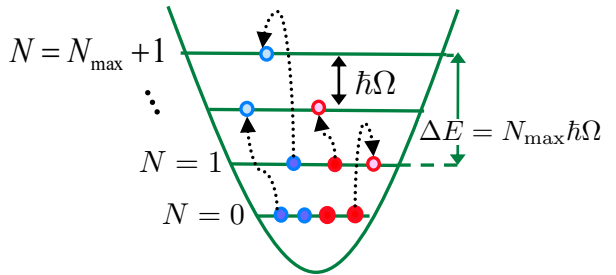
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No-Core Shell Model with Continuum (NCSMC)

$$\Psi^{(A)} = \underbrace{\sum_{\lambda} c_{\lambda} \left| \begin{array}{c} (A) \\ \text{NCSM} \end{array}, \lambda \right\rangle}_{\text{Static solutions for aggregate system, describe all nucleons close together}} + \underbrace{\sum_{\nu} \int d\vec{r} \gamma_{\nu}(\vec{r}) \hat{A}_{\nu} \left| \begin{array}{c} \text{NCSM} \\ (A-a) \quad (a) \end{array}, \nu \right\rangle}_{\text{Continuous microscopic cluster states, describe long-range projectile-target}}$$


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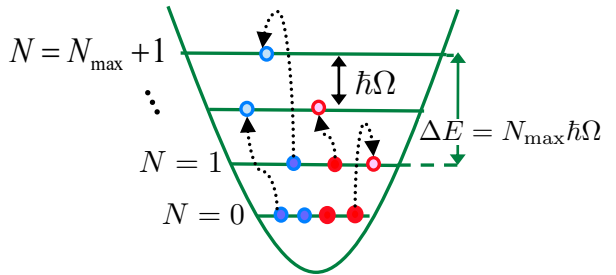
Ab initio calculations of structure including weakly bound halo states, scattering, reactions

Unified approach to bound & continuum states

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No-Core Shell Model with Continuum (NCSMC)

Unknowns

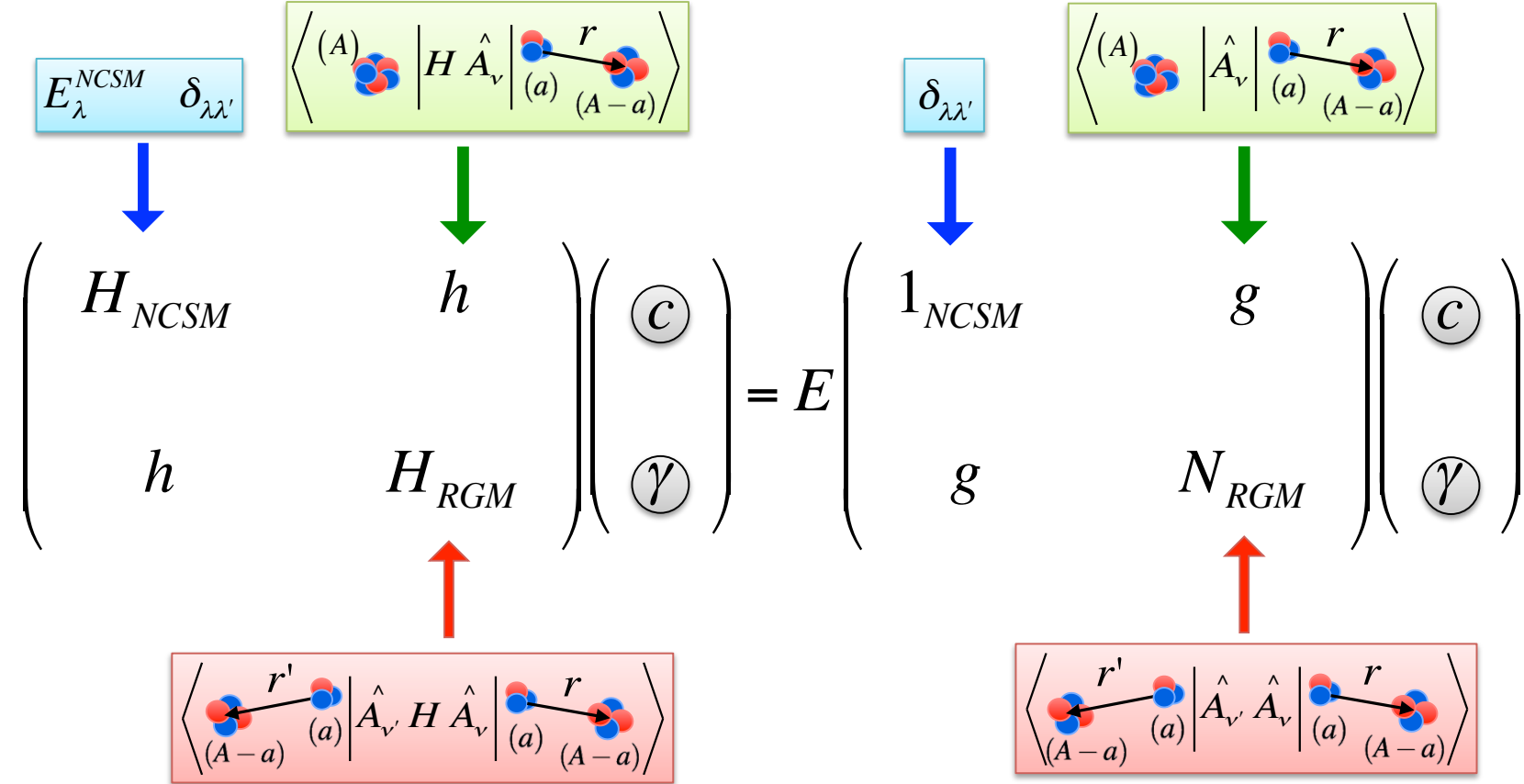
$$\Psi^{(A)} = \underbrace{\sum_{\lambda} c_{\lambda} \left| \begin{array}{c} (A) \\ \text{cluster} \\ \lambda \end{array} \right\rangle}_{\text{Static solutions for aggregate system, describe all nucleons close together}} + \underbrace{\sum_{\nu} \int d\vec{r} \gamma_{\nu}(\vec{r}) \hat{A}_{\nu} \left| \begin{array}{c} (A-a) \\ \text{cluster} \\ (a) \\ \nu \end{array} \right\rangle}_{\text{Continuous microscopic cluster states, describe long-range projectile-target}}$$


Static solutions for aggregate system,
describe all nucleons close together

Coupled NCSMC equations

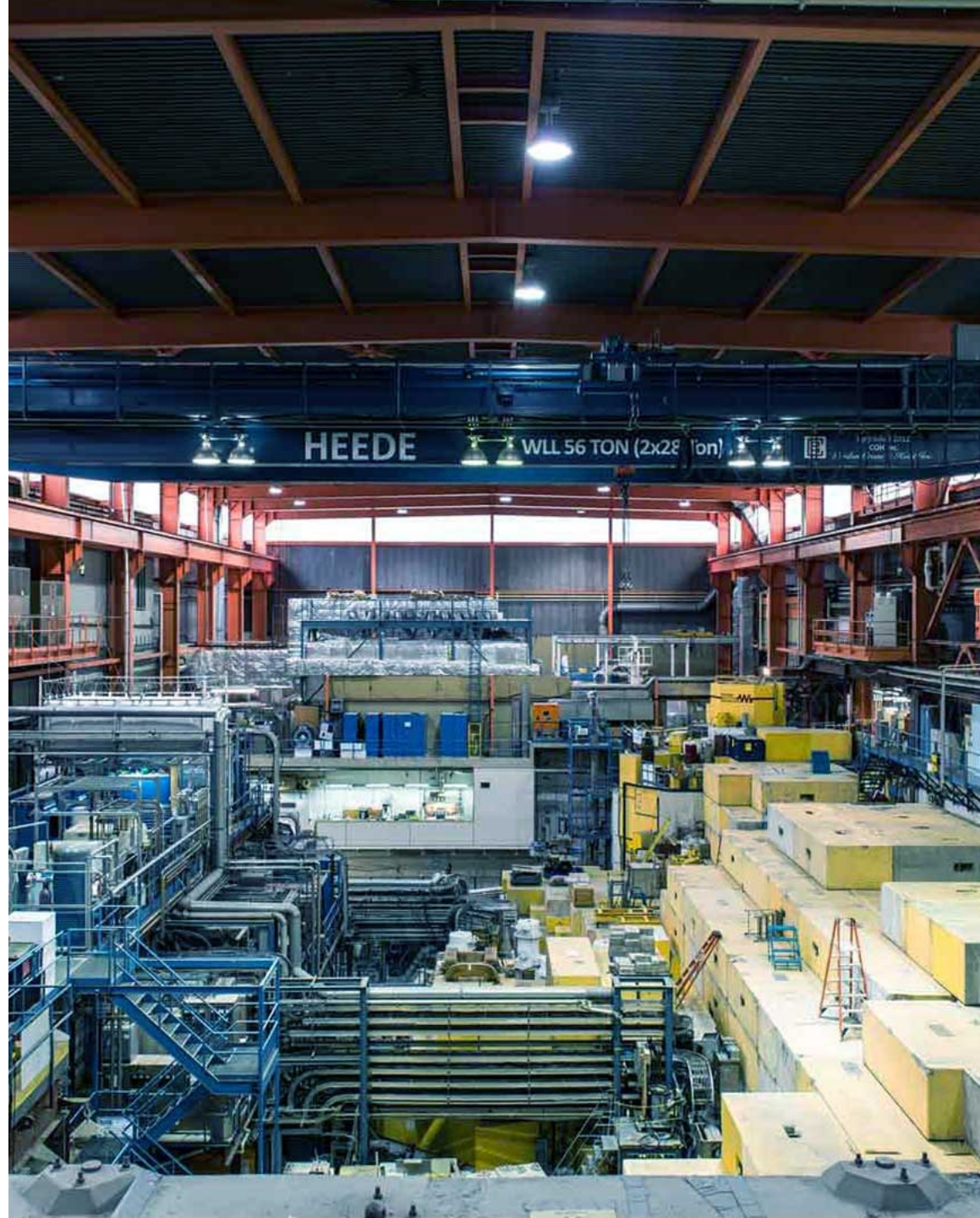
$$H \Psi^{(A)} = E \Psi^{(A)}$$

$$\Psi^{(A)} = \sum_{\lambda} c_{\lambda} \left| \begin{array}{c} (A) \\ \text{cluster} \end{array}, \lambda \right\rangle + \sum_{\nu} \int d\vec{r} \gamma_{\nu}(\vec{r}) \hat{A}_{\nu} \left| \begin{array}{c} (A-a) \\ \text{cluster} \end{array}, \nu \right\rangle$$



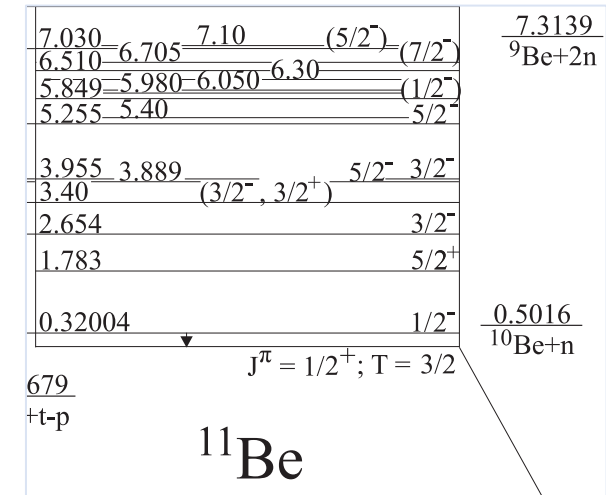
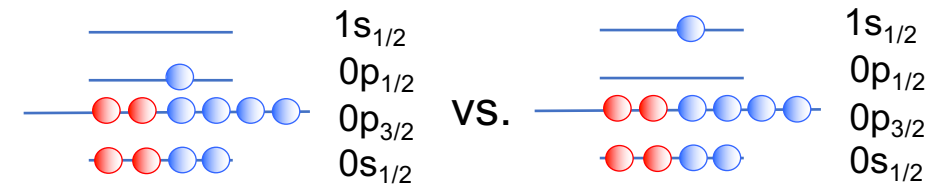
Parity inversion in ^{11}Be ground state

2025-10-24



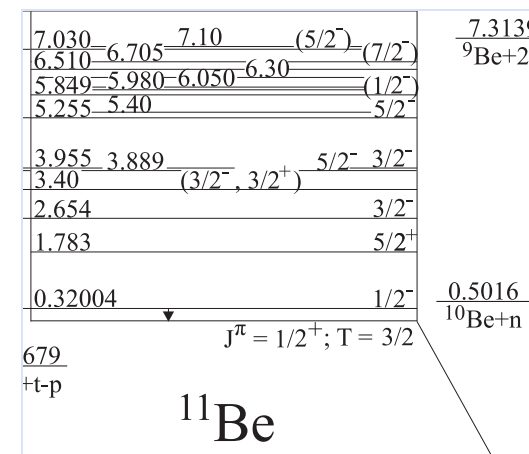
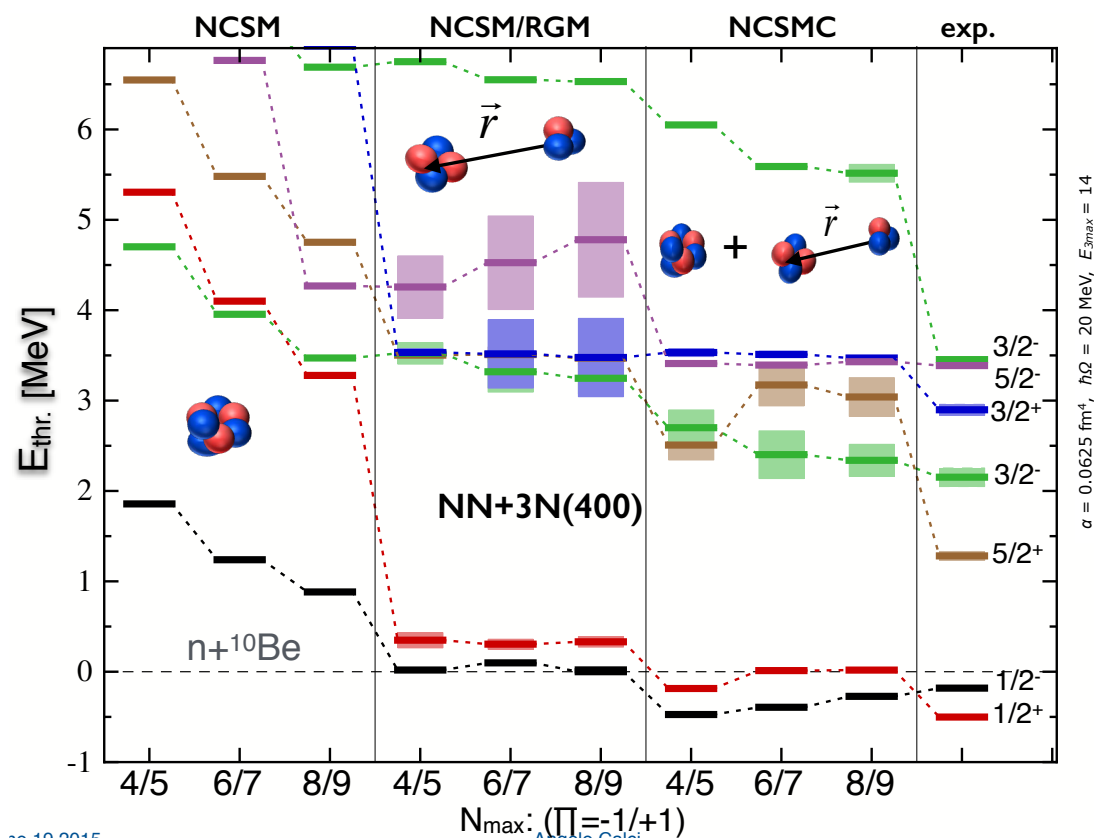
Neutron-rich halo nucleus ^{11}Be

- $Z=4, N=7$
 - In the shell model picture g.s. expected to be $J^\pi=1/2^-$
 - $Z=6, N=7$ ^{13}C and $Z=8, N=7$ ^{15}O have $J^\pi=1/2^-$ g.s.
 - In reality, ^{11}Be g.s. is $J^\pi=1/2^+$ - parity inversion
 - Very weakly bound: $E_{\text{th}}=-0.5 \text{ MeV}$
 - Halo state – dominated by $^{10}\text{Be-n}$ in the S-wave
 - The $1/2^-$ state also bound – only by 180 keV
- Can we describe ^{11}Be in *ab initio* calculations?
 - Continuum must be included
 - Does the 3N interaction play a role in the parity inversion?

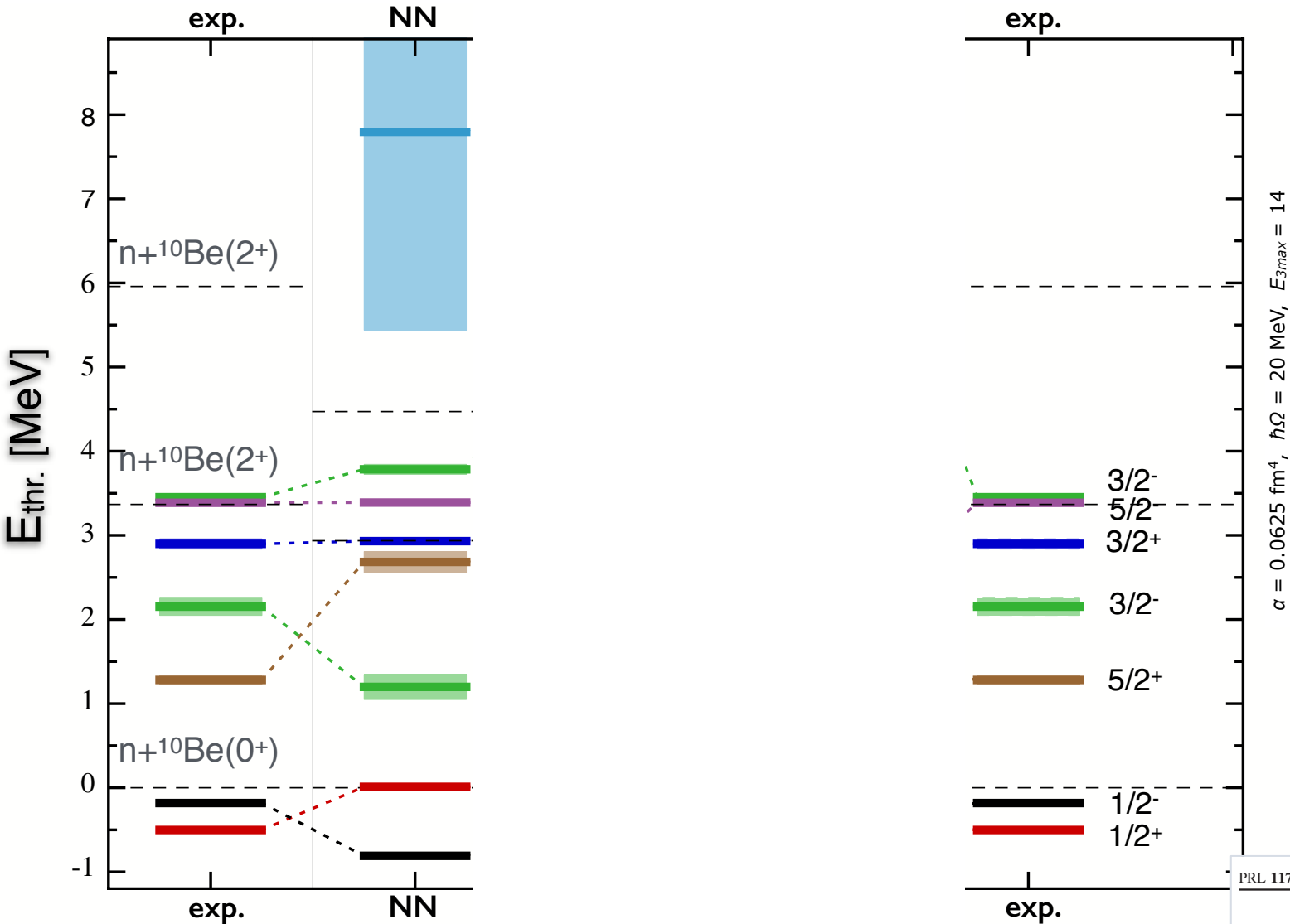


Structure of ^{11}Be from chiral NN+3N forces

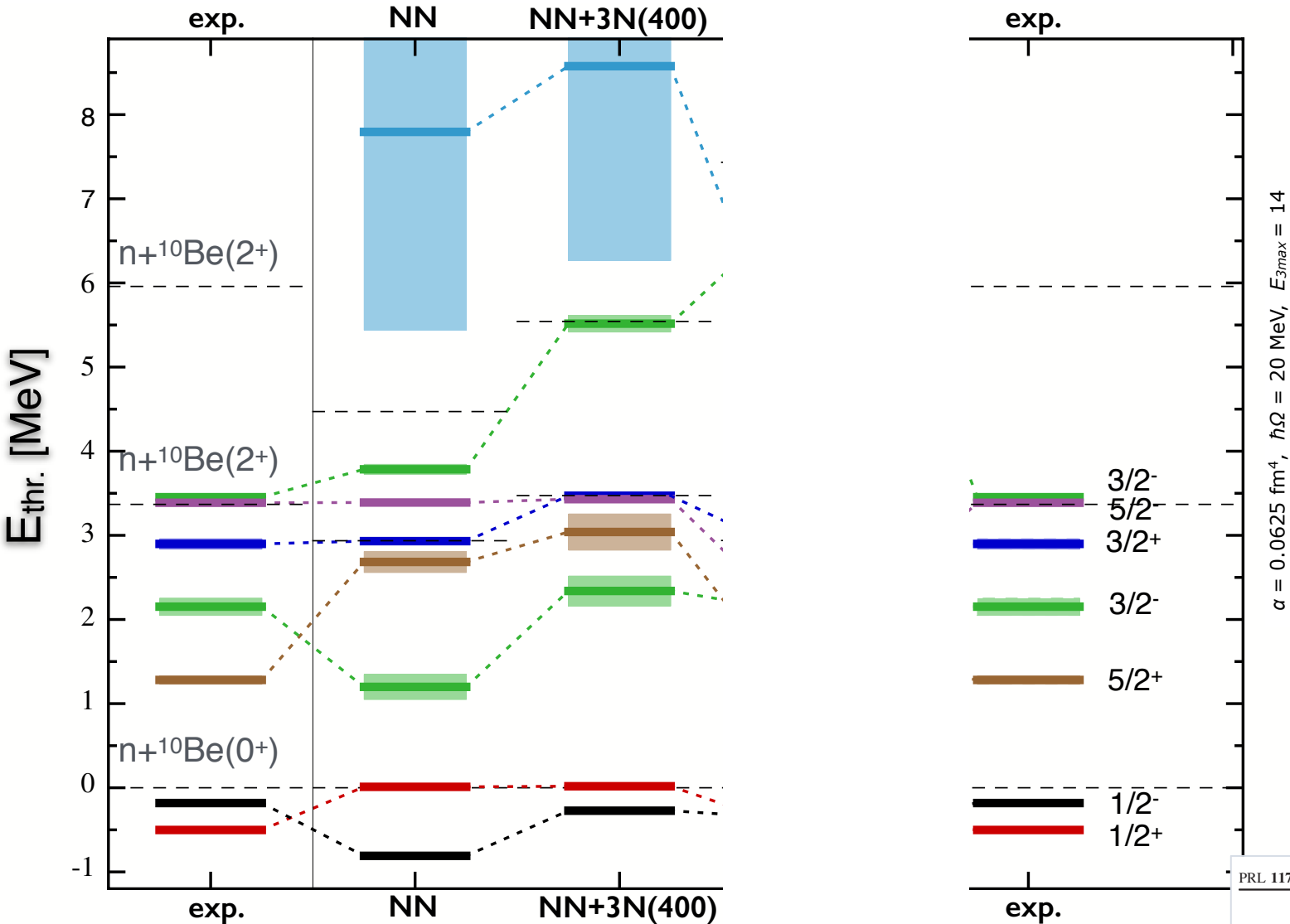
- NCSMC calculations **including chiral 3N** (NN N³LO+N²LO 3N(400), N²LO_{sat}, NN N⁴LO+3N_{lnl})
 - n-¹⁰Be + ¹¹Be
 - ¹⁰Be: 0⁺, 2⁺, 2⁺ NCSM eigenstates
 - ¹¹Be: ≥6 π = -1 and ≥3 π = +1 NCSM eigenstates



¹¹Be within NCSMC: Discrimination among chiral nuclear forces



¹¹Be within NCSMC: Discrimination among chiral nuclear forces

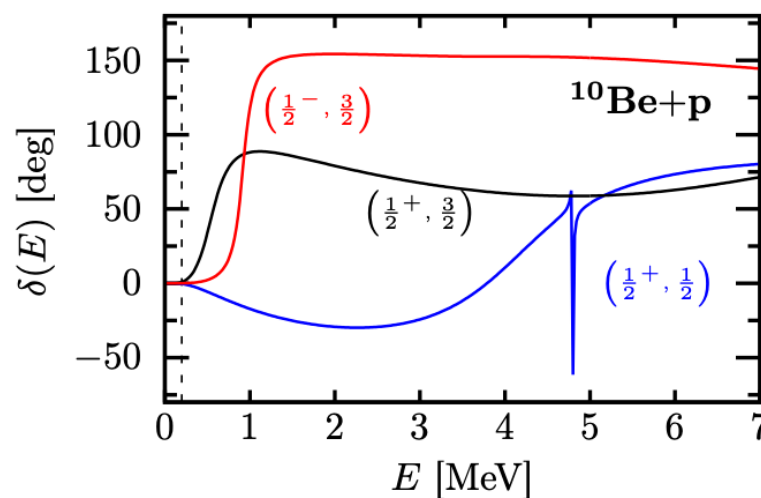
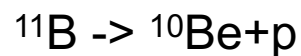


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PHYSICAL REVIEW C **105**, 054316 (2022)

***Ab initio* calculation of the β decay from ^{11}Be to a $^{10}\text{Be} + p$ resonance**

M. C. Atkinson¹, P. Navrátil¹, G. Hupin², K. Kravvaris³ and S. Quaglioni³



Parity inversion reproduced

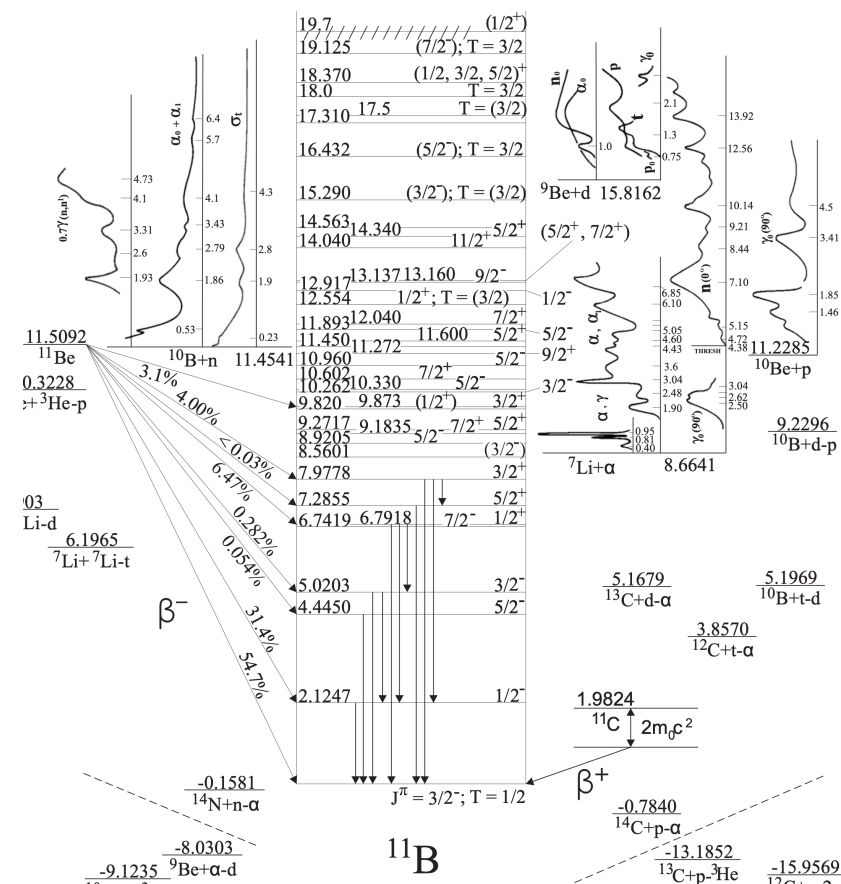
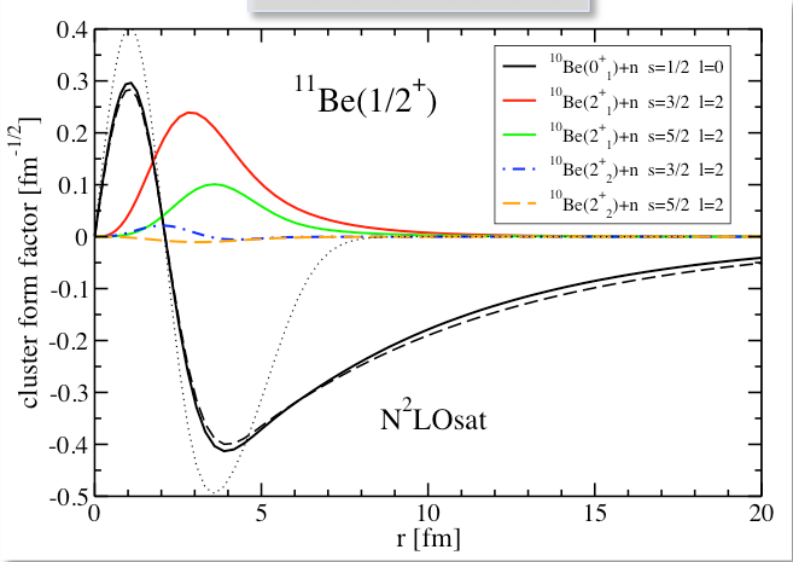


Photo-disassociation of ^{11}Be

Halo structure



cluster form factor

$$= r \langle \Phi_{vr}^{J^{\pi T}} | \hat{A}_v | \psi^{J^{\pi T}} \rangle$$

$$| \Phi_{vr}^{J^{\pi T}} \rangle = \left[\left(| ^{10}\text{Be } \alpha_1 I_1^{\pi_1 T_1} \rangle | n \frac{1}{2}^+ \frac{1}{2} \rangle \right)^{(sT)} Y_\ell(\hat{r}_{10,1}) \right]^{(J^{\pi T})} \frac{\delta(r - r_{10,1})}{rr_{10,1}}$$

PRL 117, 242501 (2016) PHYSICAL REVIEW LETTERS week ending 9 DECEMBER 2016

Can *Ab Initio* Theory Explain the Phenomenon of Parity Inversion in ^{11}Be ?

Angelo Calci,^{1,*} Petr Navrátil,^{1,†} Robert Roth,² Jérémy Dohet-Eraly,^{1,‡} Sofia Quaglioni,³ and Guillaume Hupin^{4,5}

Bound to bound	NCSM	NCSMC-phenom	Expt.
B(E1; $1/2^+ \rightarrow 1/2^-$) [$e^2 fm^2$]	0.0005	0.117	0.102(2)

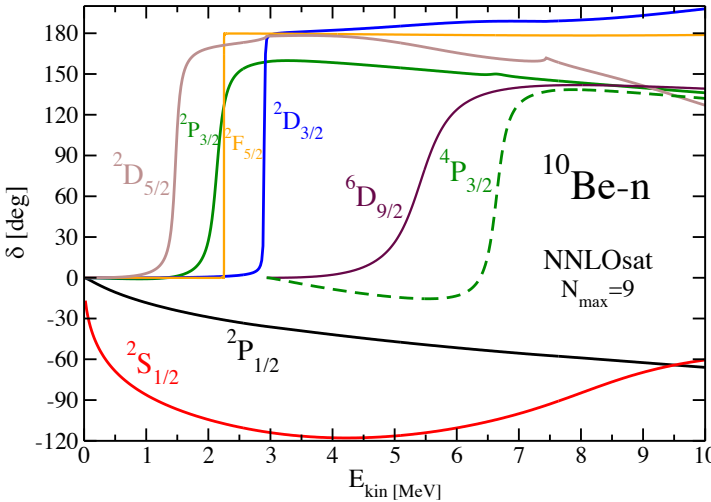
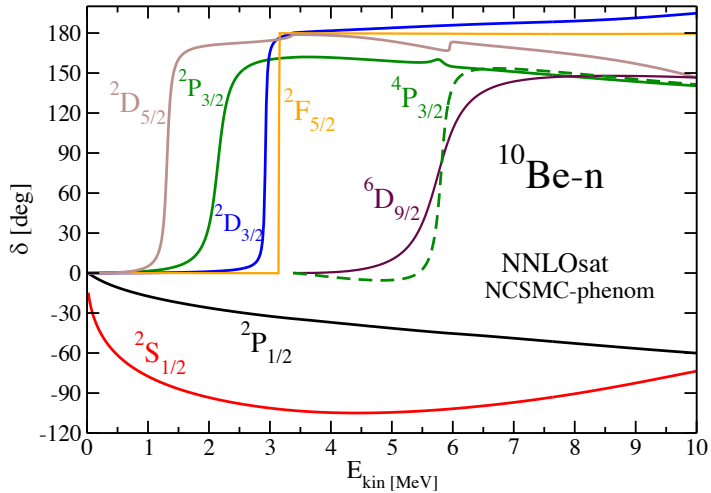
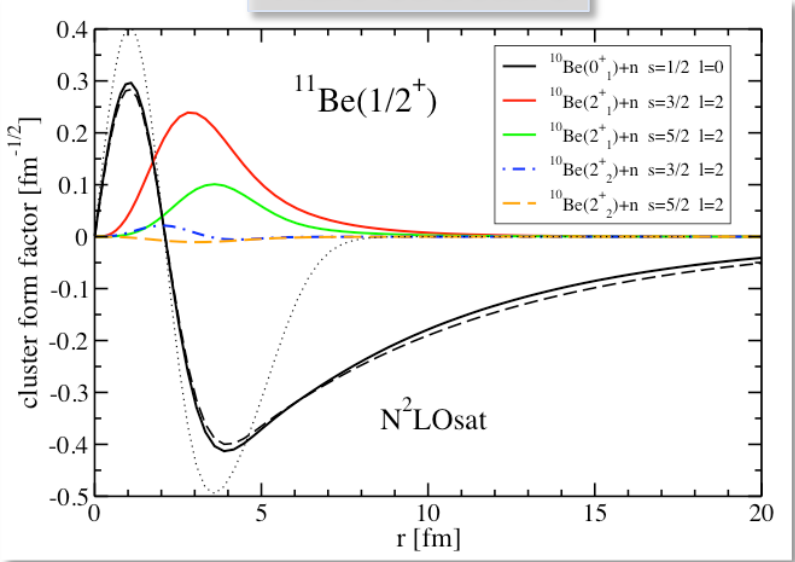


Photo-disassociation of ¹¹Be

Halo structure



cluster form factor

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$$\left| \Phi_{vr}^{J^{\pi T}} \right\rangle = \left[\left(\left| {}^{10}\text{Be } \alpha_1 I_1^{\pi_1 T_1} \right| \left| n \frac{1}{2}^+ \frac{1}{2} \right\rangle \right)^{(sT)} Y_\ell(\hat{r}_{10,1}) \right]^{(J^{\pi T})} \frac{\delta(r - r_{10,1})}{rr_{10,1}}$$

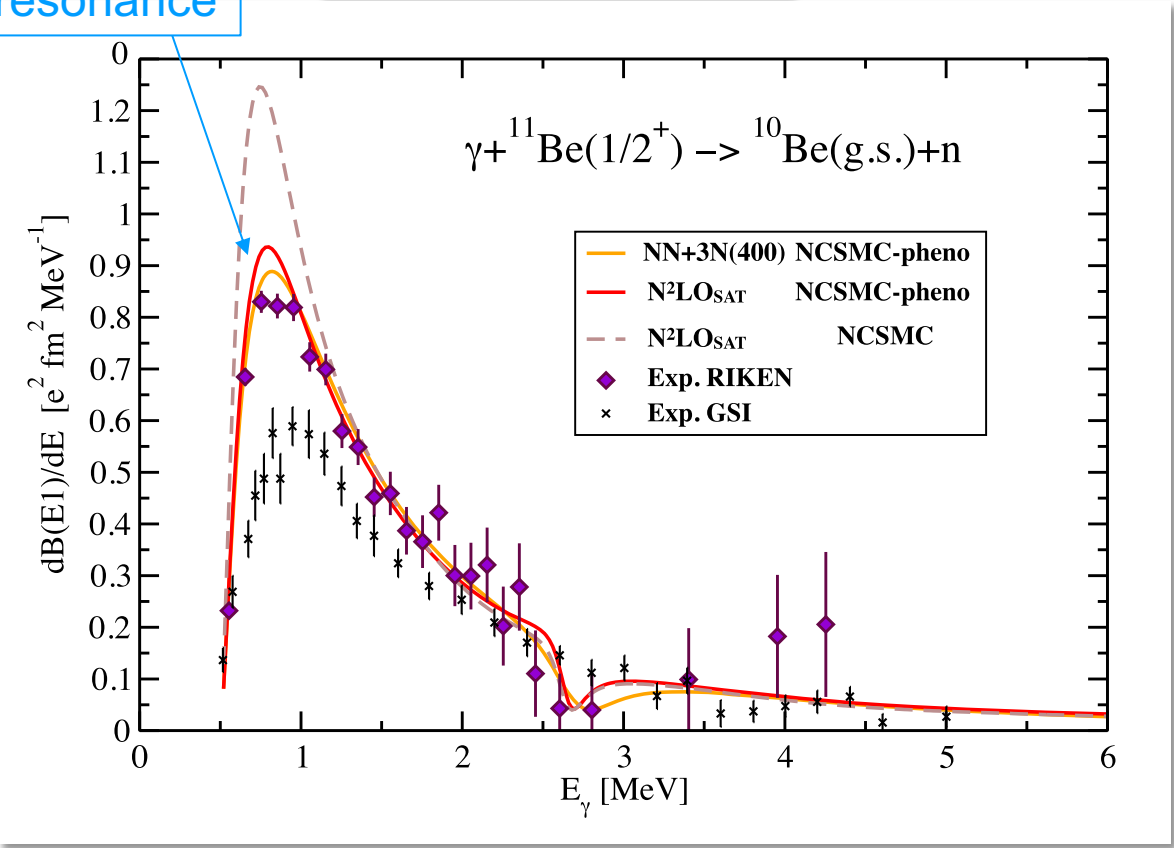
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B(E1; 1/2 ⁺ →1/2 ⁻) [e ² fm ²]	0.0005	0.117	0.102(2)

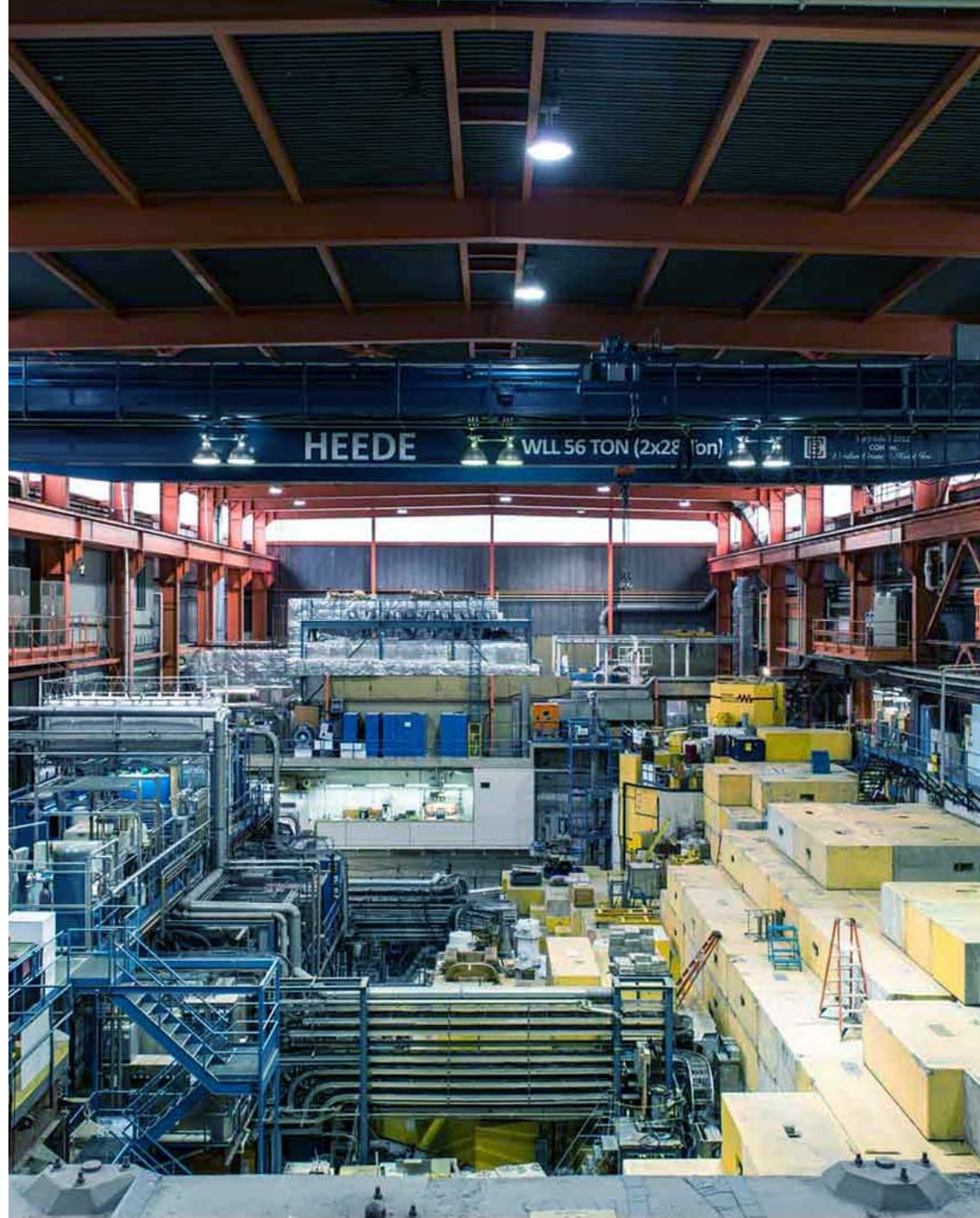
Bound to continuum

Not a resonance

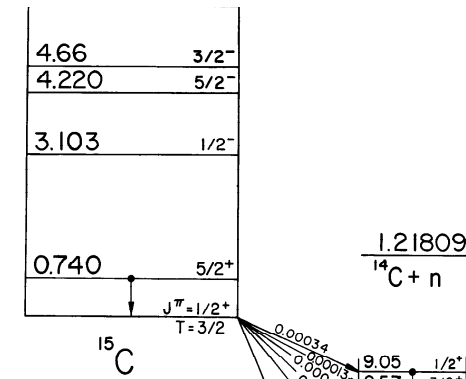


Halo ground state of ^{15}C

2025-10-24

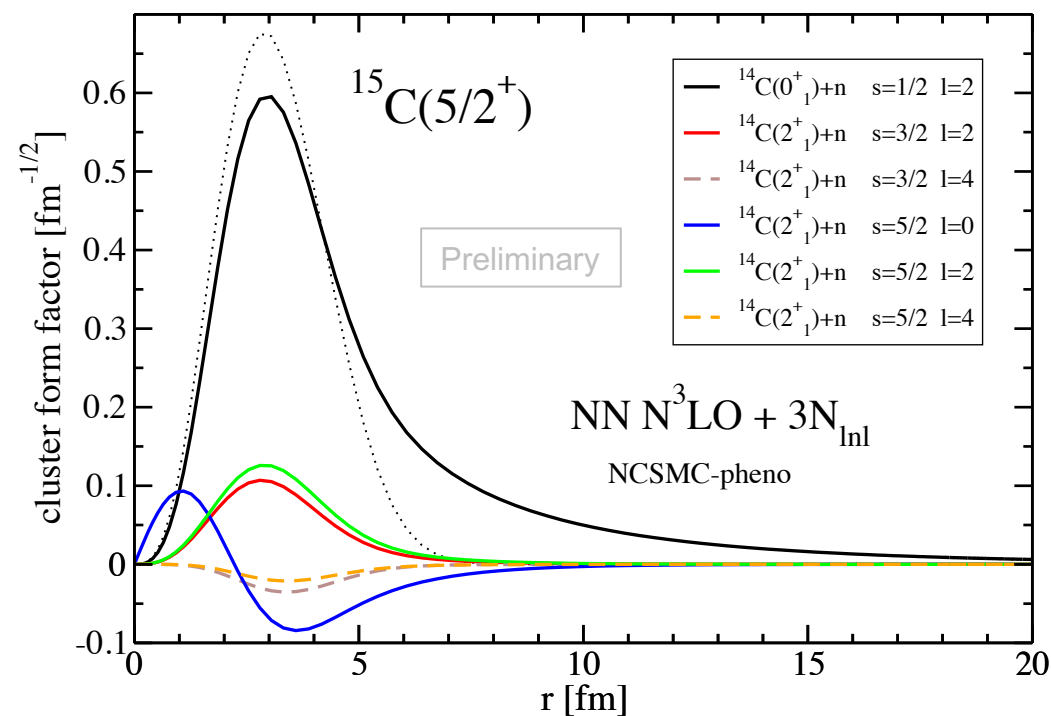
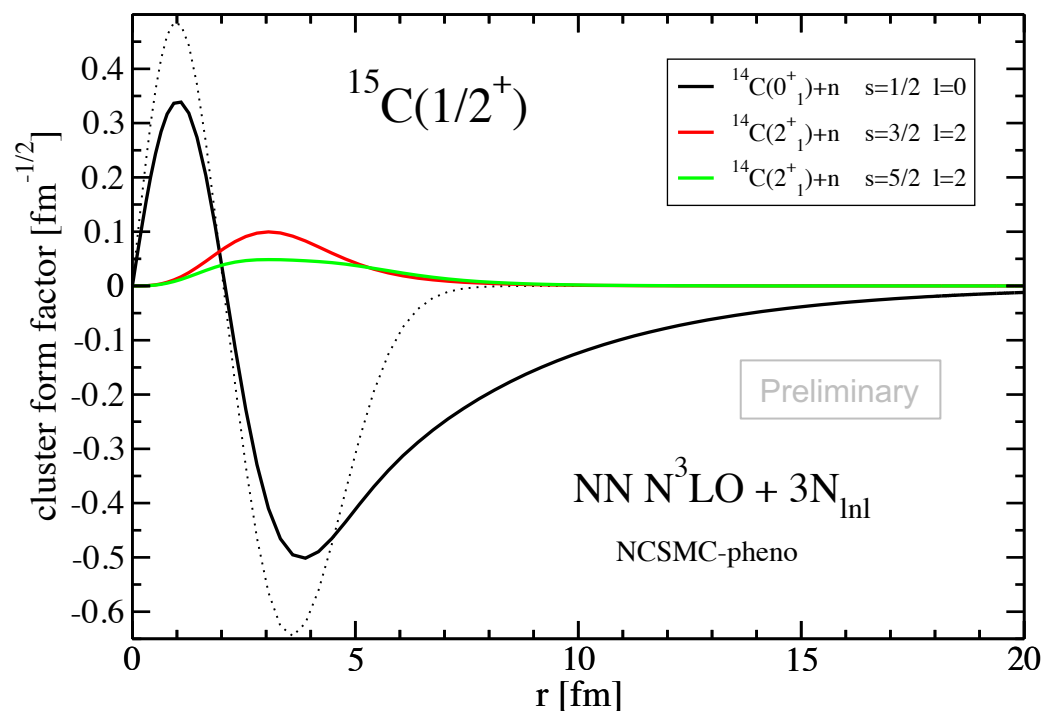


NCSMC application to $^{15}\text{C} \sim ^{14}\text{C} + n$ – cluster form factor calculations



■ $1/2^+$ S-wave and $5/2^+$ D-wave Asymptotic Normalization Coefficients

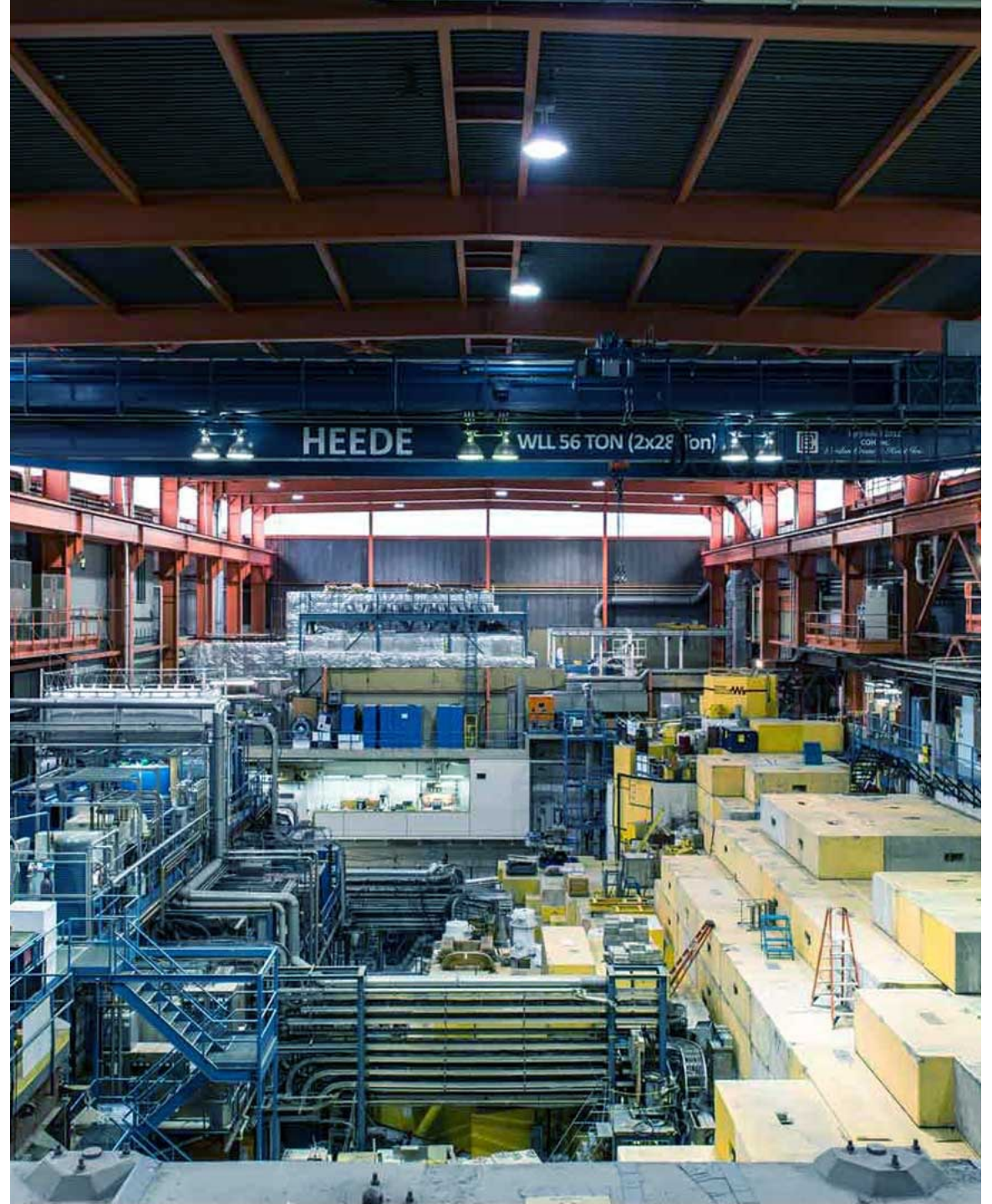
- $C_{1/2^+} = 1.282 \text{ fm}^{-1/2}$ - compare to Moschini & Capel inferred from transfer data: $1.26(2) \text{ fm}^{-1/2}$
- $C_{5/2^+} = 0.048 \text{ fm}^{-1/2}$ $0.056(1) \text{ fm}^{-1/2}$
- Spectroscopic factors: 0.96 for $1/2^+$ and 0.90 for $5/2^+$ - experiments 0.95(5) and 0.69, resp.



P -wave halo ground state of ${}^8\text{B}$

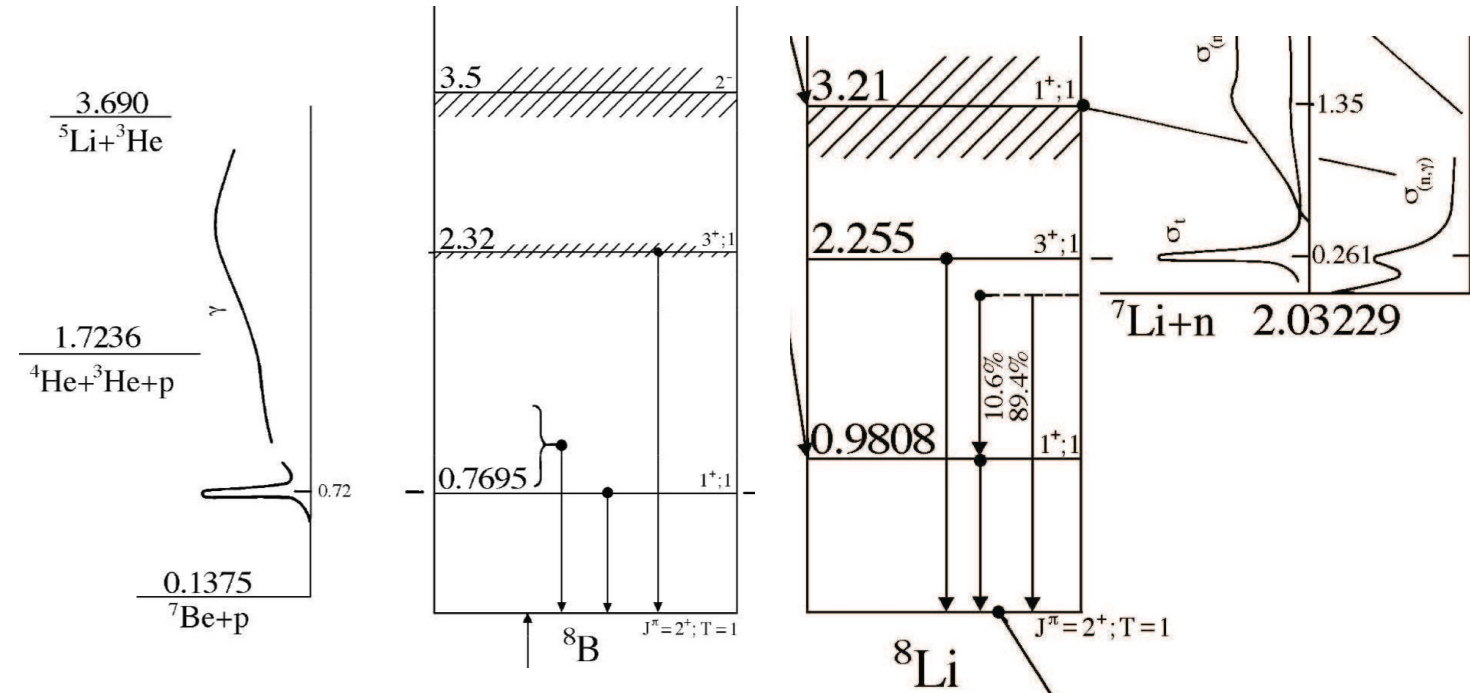
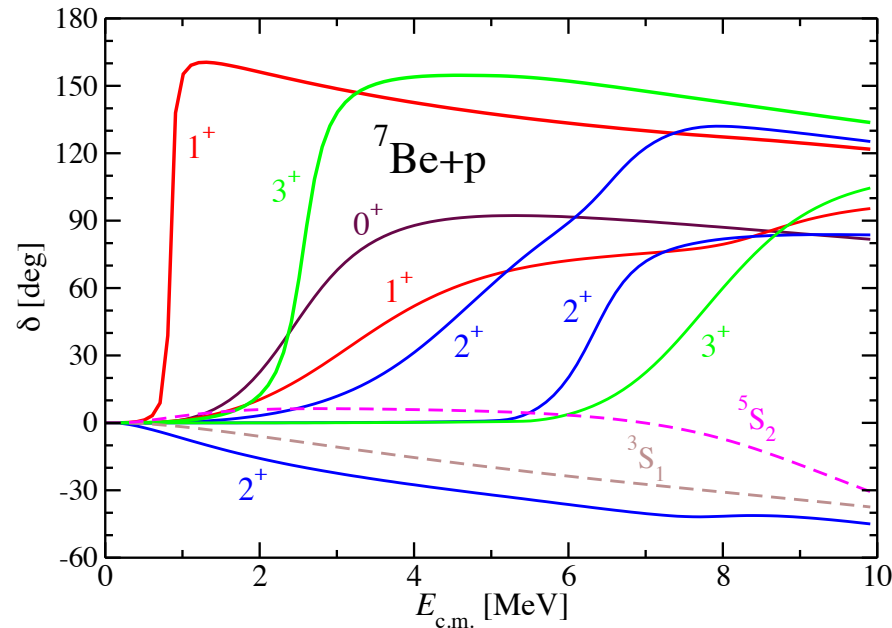
&

radiative capture of protons
on ${}^7\text{Be}$



Radiative capture of protons on ${}^7\text{Be}$

- Solar pp chain reaction, solar ^8B neutrinos
- NCSMC calculations with a set of chiral NN+3N interactions as input
- Example of ^8B structure results



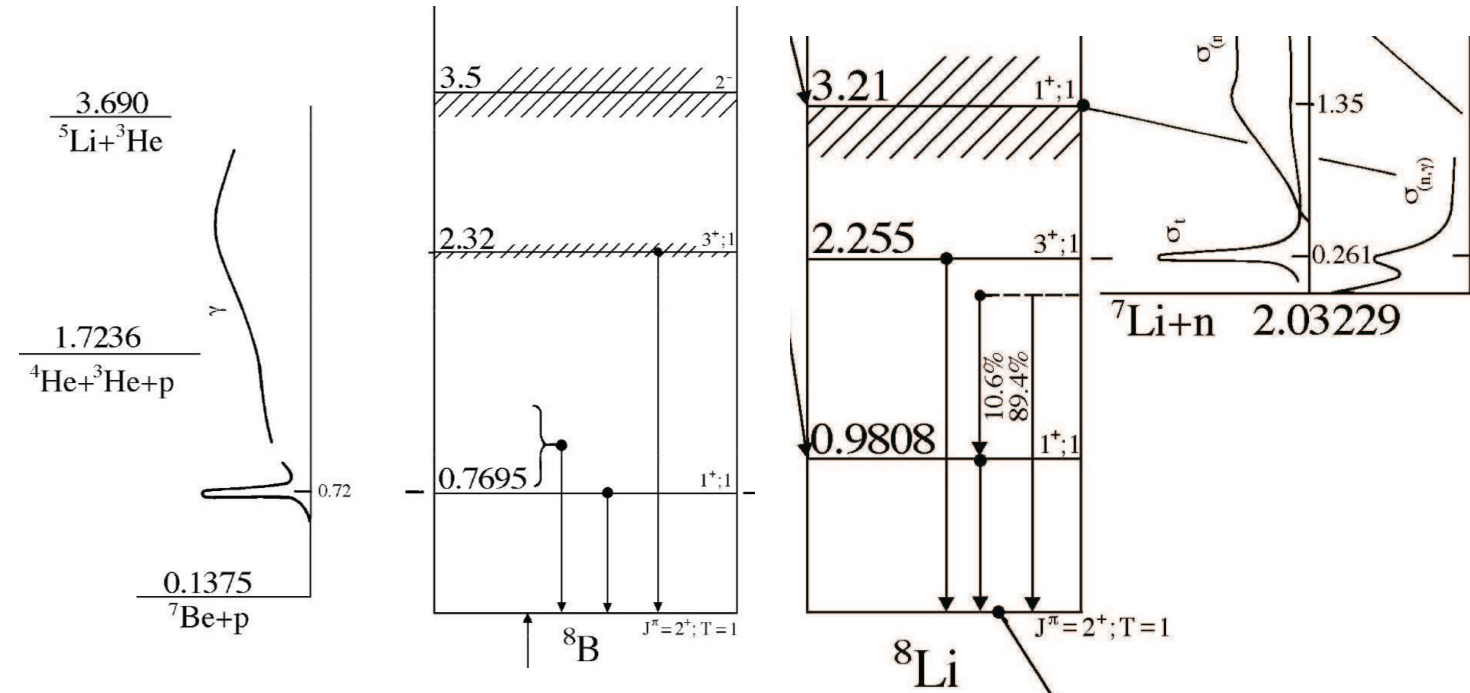
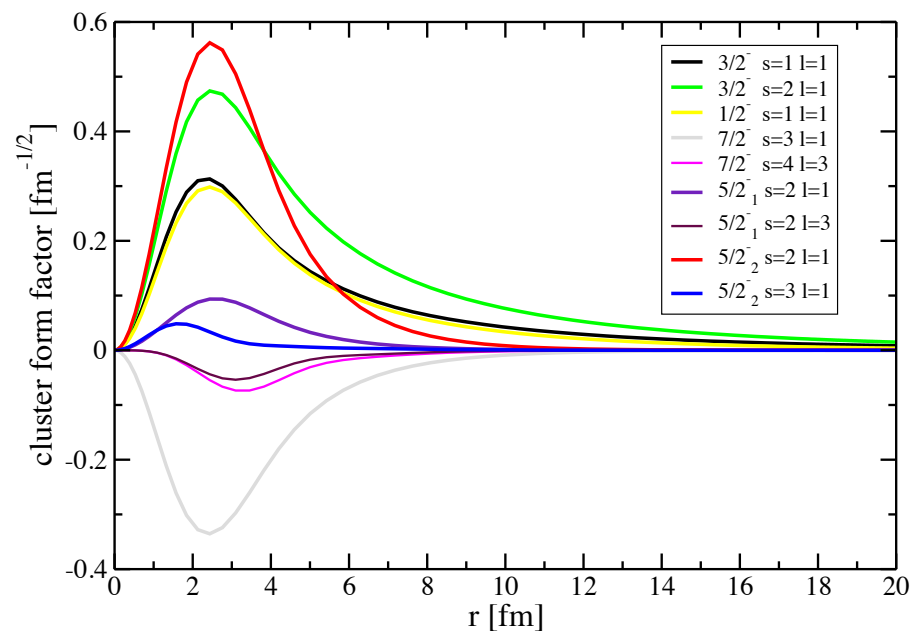


Ab initio informed evaluation of the radiative capture of protons on ^7Be

K. Kravvaris^{a,*}, P. Navrátil^b, S. Quaglioni^a, C. Hebborn^{c,a}, G. Hupin^d

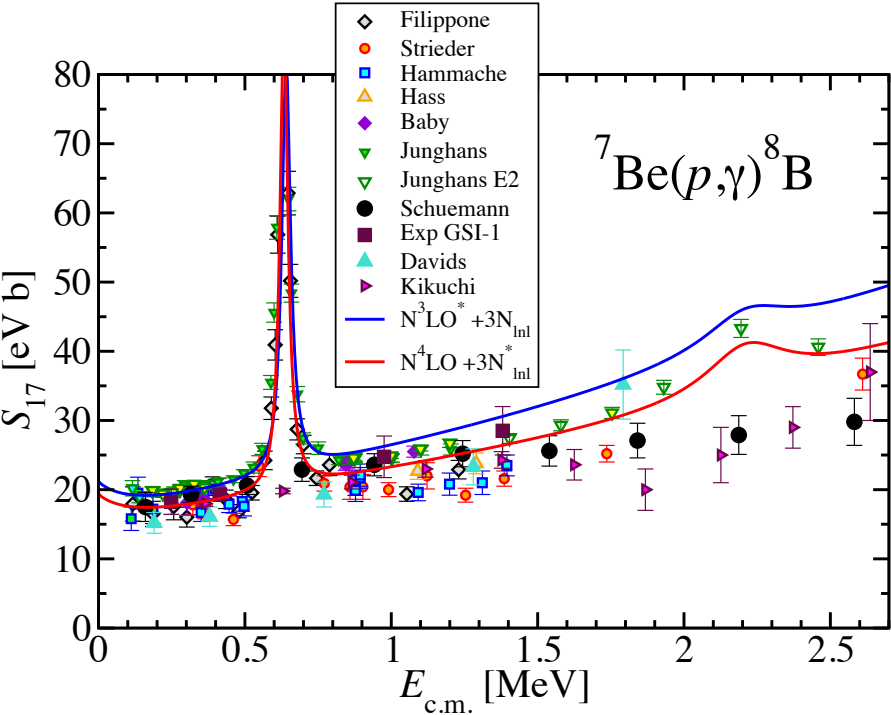
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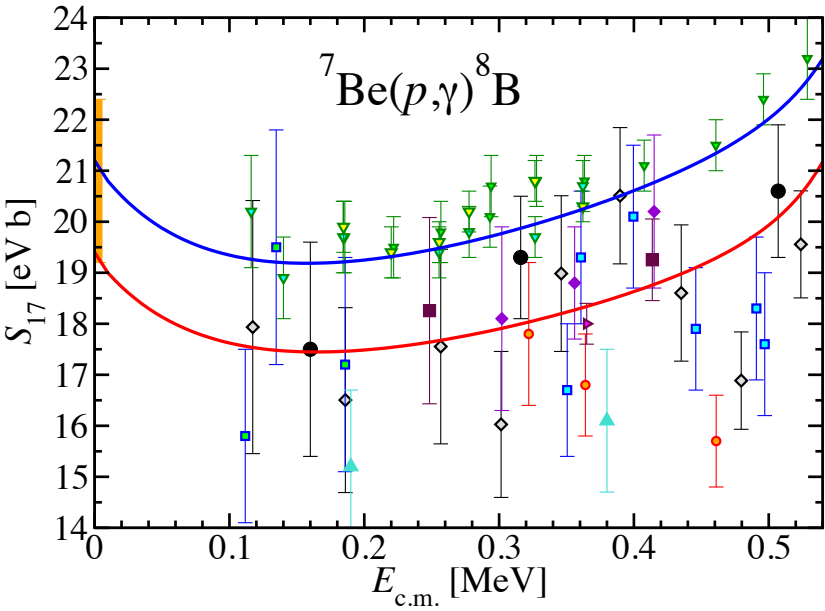
Radiative capture of protons on ⁷Be

■ NCSMC S-factor results



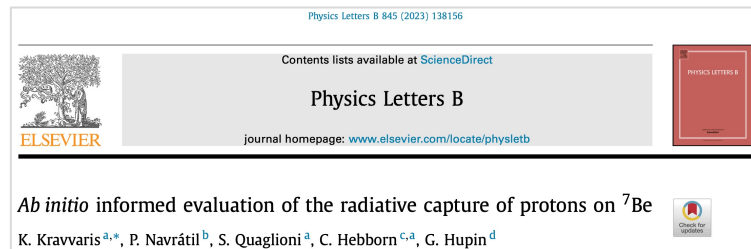
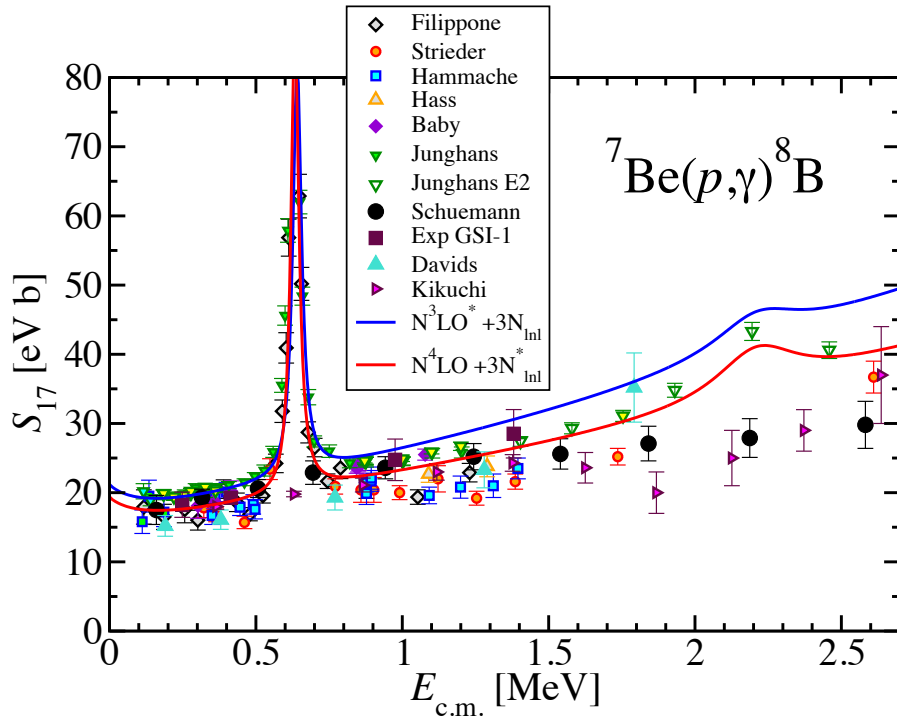
E1 non-resonant, M1/E2 at 1⁺ and 3⁺ resonances

	$C_{p_{1/2}}$	$C_{p_{3/2}}$	a_1	a_2	$S_{17}(0)$
N^2LO+3N_{lnl}	0.384	0.691	4.4(1)	-0.5(1)	23.9
N^3LO+3N_{lnl}	0.390	0.678	1.3(1)	-4.7(1)	23.5
N^4LO+3N_{lnl}	0.354	0.669	1.6(1)	-4.4(1)	22.0
$N^4LO+3N_{lnl}^*$	0.343	0.621	1.3(1)	-5.0(1)	19.3
$N^3LO^*+3N_{lnl}$	0.334	0.663	0.1(1)	-7.7(1)	21.1
$N^3LO^*+3N_{loc}$	0.308	0.584	2.5(1)	-3.6(2)	16.8
Ref. [41]	0.315(9)	0.66(2)	$17.34^{+1.11}_{-1.33}$	$-3.18^{+0.55}_{-0.50}$	

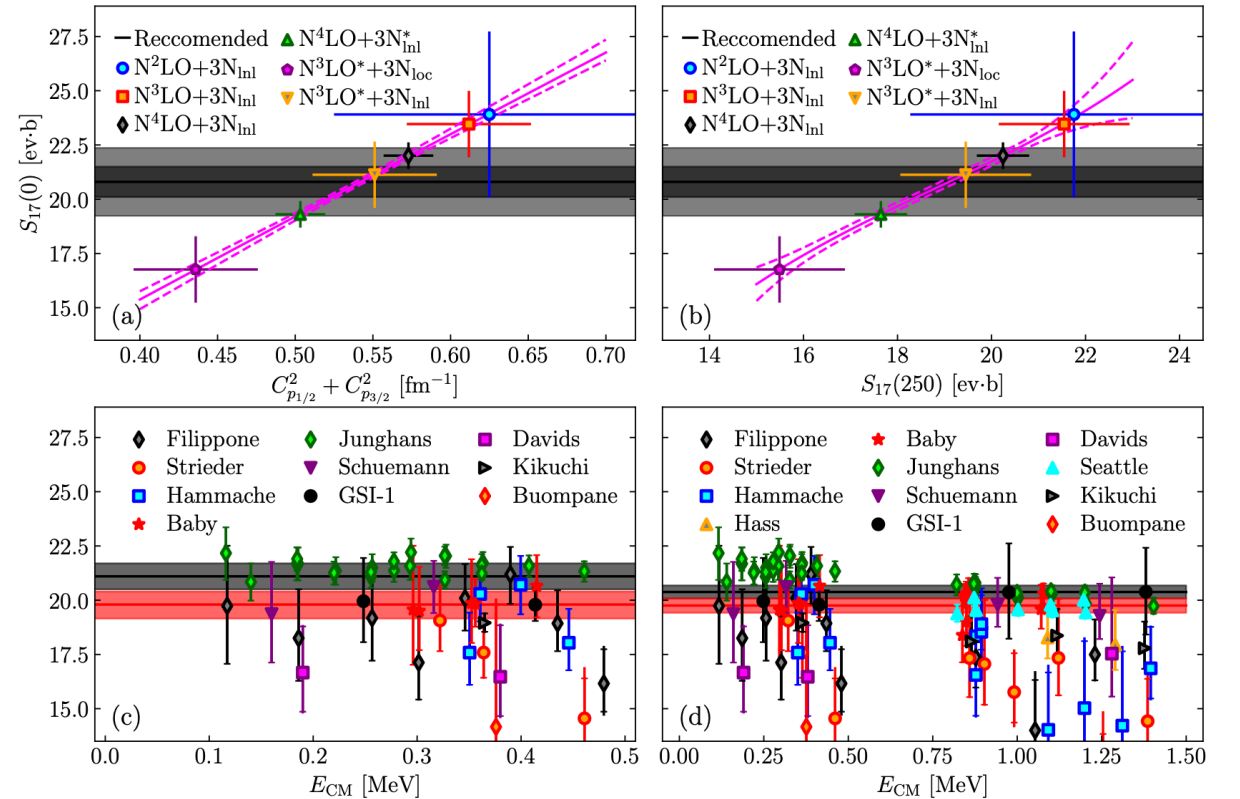


Radiative capture of protons on ^7Be

■ NCSMC S-factor results

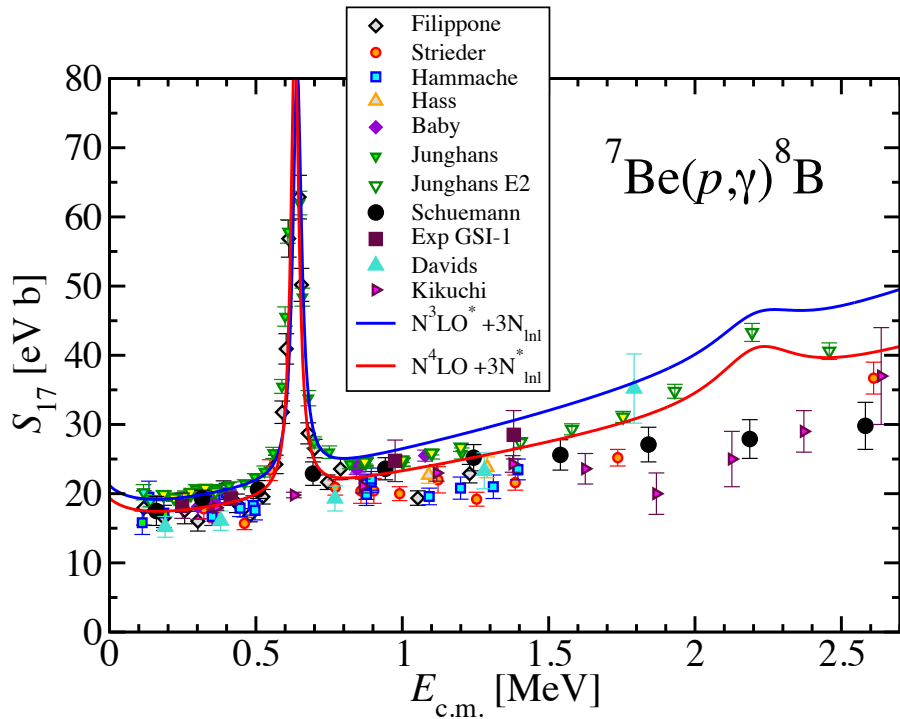


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Radiative capture of protons on ^7Be

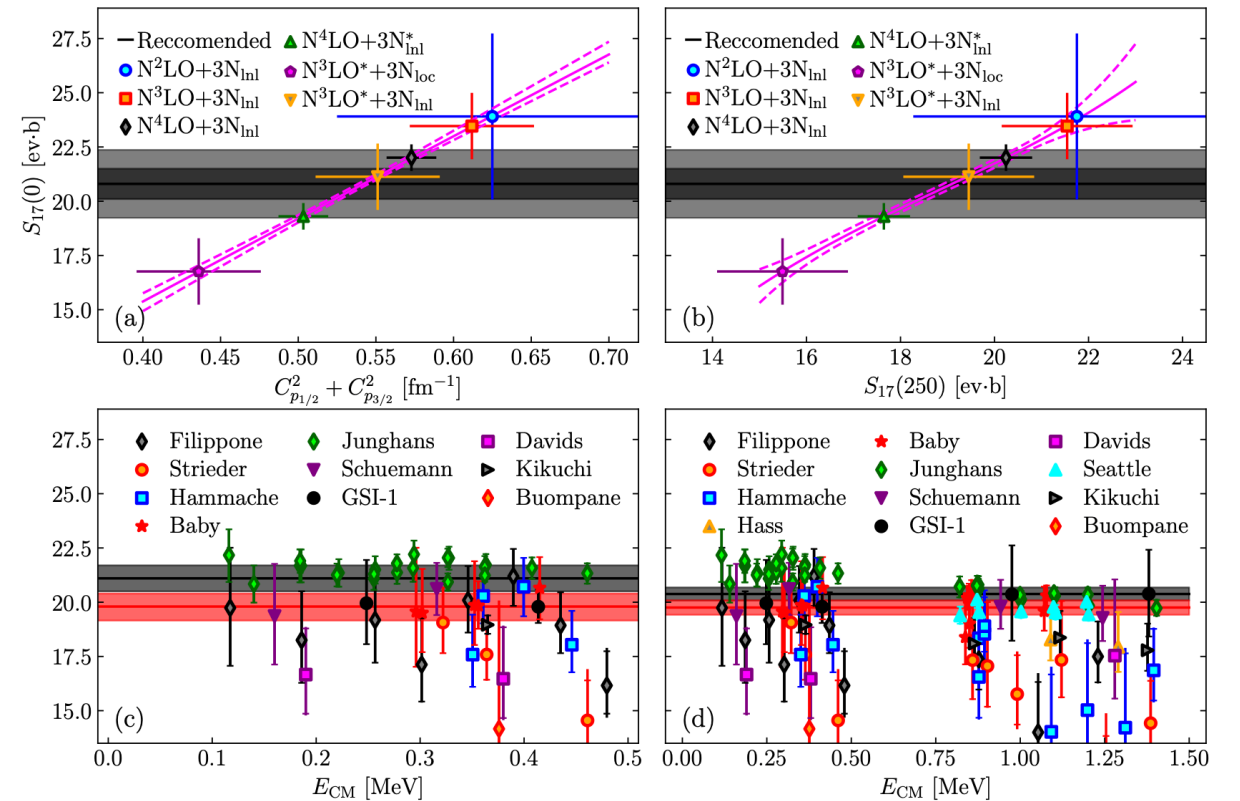
■ NCSMC S-factor results



Recommended value $S_{17}(0) \sim 19.8(3)$ eV b

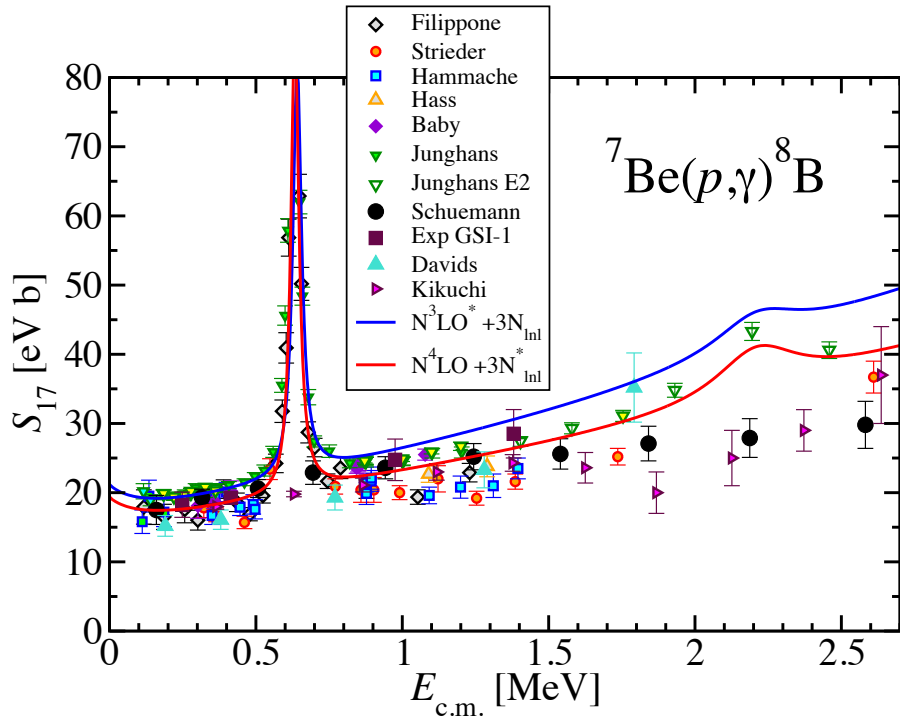
2011 evaluation in *Rev. Mod. Phys.* **83**,195–245 (2011):
 $S_{17}(0) = 20.8 \pm 0.7(\text{expt}) \pm 1.4(\text{theory})$ eV b

	$C_{p_{1/2}}$	$C_{p_{3/2}}$	a_1	a_2	$S_{17}(0)$
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$N^3\text{LO}^* + 3N_{\text{loc}}$	0.308	0.584	2.5(1)	-3.6(2)	16.8
Ref. [41]	0.315(9)	0.66(2)	$17.34^{+1.11}_{-1.33}$	$-3.18^{+0.55}_{-0.50}$	



Radiative capture of protons on ^7Be

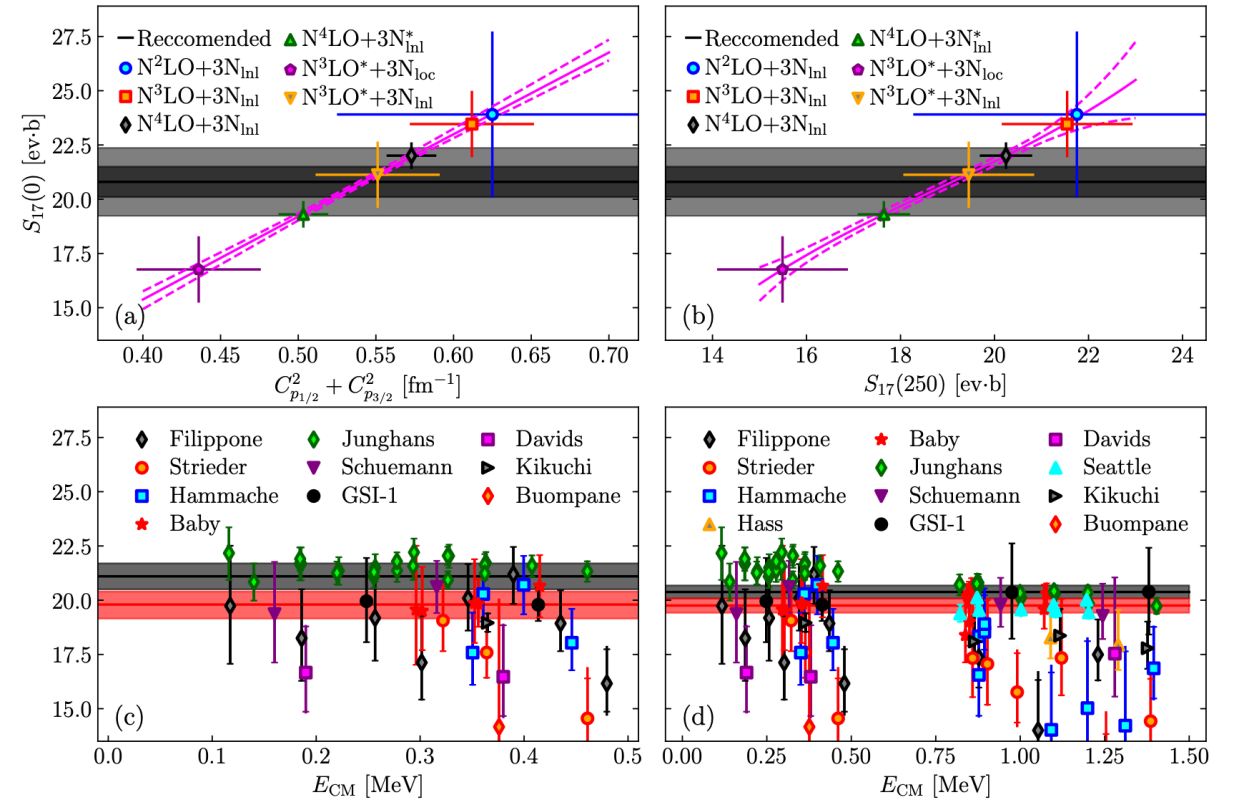
■ NCSMC S-factor results



Recommended value $S_{17}(0) \sim 19.8(3)$ eV b

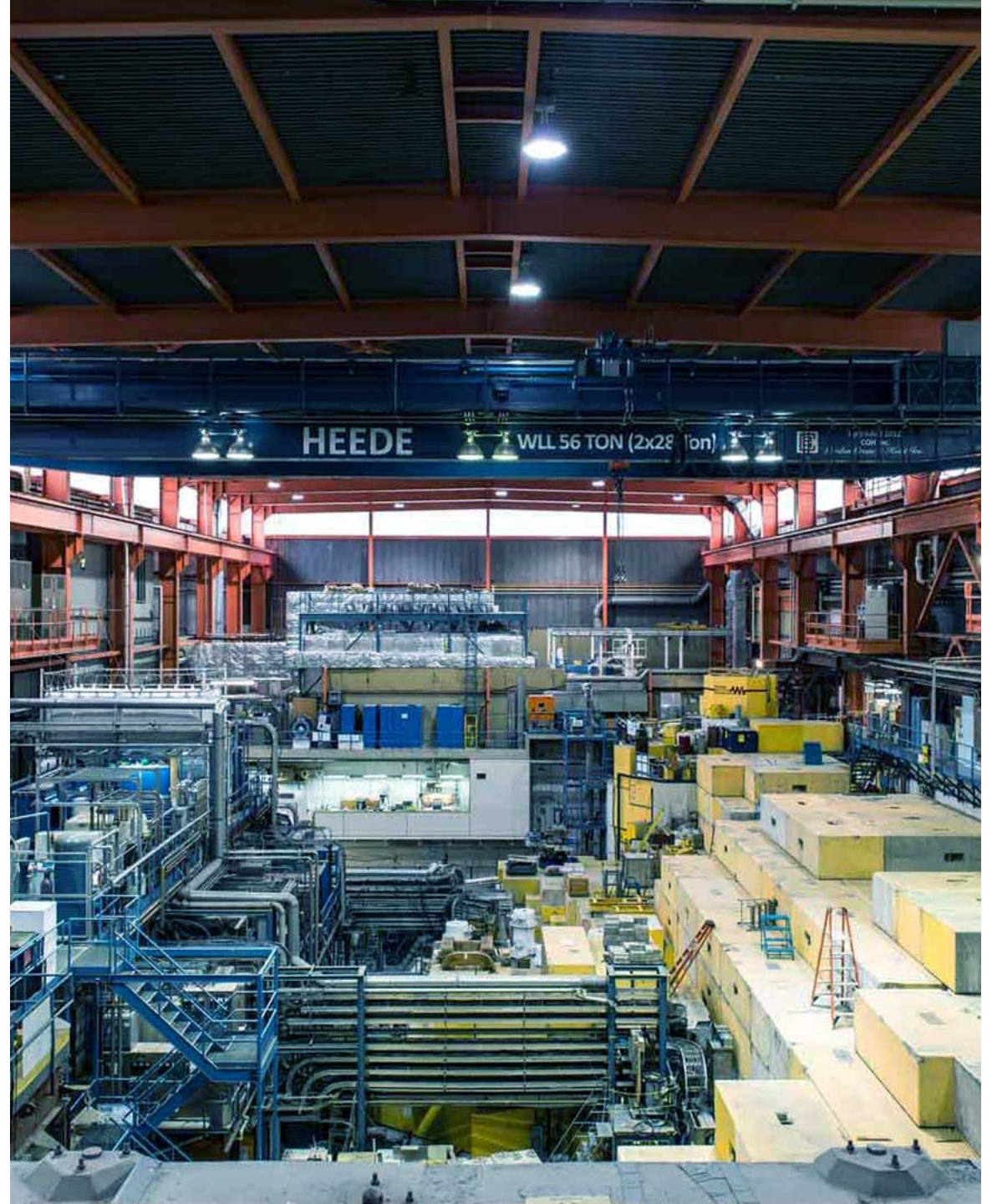
New evaluation in *Rev. Mod. Phys.* **97**, 035002 (2025):
 $S_{17}(0) = 20.5 \pm 0.7$ eV b

	$C_{p_{1/2}}$	$C_{p_{3/2}}$	a_1	a_2	$S_{17}(0)$
$N^2\text{LO} + 3N_{\text{lnl}}$	0.384	0.691	4.4(1)	-0.5(1)	23.9
$N^3\text{LO} + 3N_{\text{lnl}}$	0.390	0.678	1.3(1)	-4.7(1)	23.5
$N^4\text{LO} + 3N_{\text{lnl}}$	0.354	0.669	1.6(1)	-4.4(1)	22.0
$N^4\text{LO} + 3N_{\text{lnl}}^*$	0.343	0.621	1.3(1)	-5.0(1)	19.3
$N^3\text{LO}^* + 3N_{\text{lnl}}$	0.334	0.663	0.1(1)	-7.7(1)	21.1
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Ref. [41]	0.315(9)	0.66(2)	$17.34^{+1.11}_{-1.33}$	$-3.18^{+0.55}_{-0.50}$	



NCSMC for three-body cluster

$${}^6\text{He} \sim {}^4\text{He} + n + n$$



NCSMC for three-body clusters

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NCSM

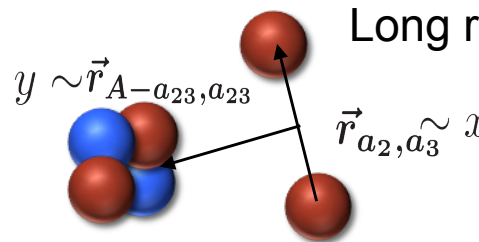
Short range description



+

NCSM/RGM-3B

Long range description



$$\Psi^{(A)} = \sum_{\lambda} c_{\lambda} \left| \begin{array}{c} \text{three-body cluster} \\ , \lambda \end{array} \right\rangle + \sum_{\nu} \int d\vec{x} d\vec{y} \gamma_{\nu}(\vec{x}, \vec{y}) \hat{A}_{\nu} \left| \begin{array}{c} \text{three-body cluster} \\ , \nu \end{array} \right\rangle$$

Unknowns

Three-cluster dynamics within the *ab initio* no-core shell model with continuum:
How many-body correlations and α clustering shape ${}^6\text{He}$

Sofia Quaglioni,^{1,*} Carolina Romero-Redondo,^{1,†} Petr Navrátil,^{2,‡} and Guillaume Hupin^{3,§}

PRL **117**, 222501 (2016)

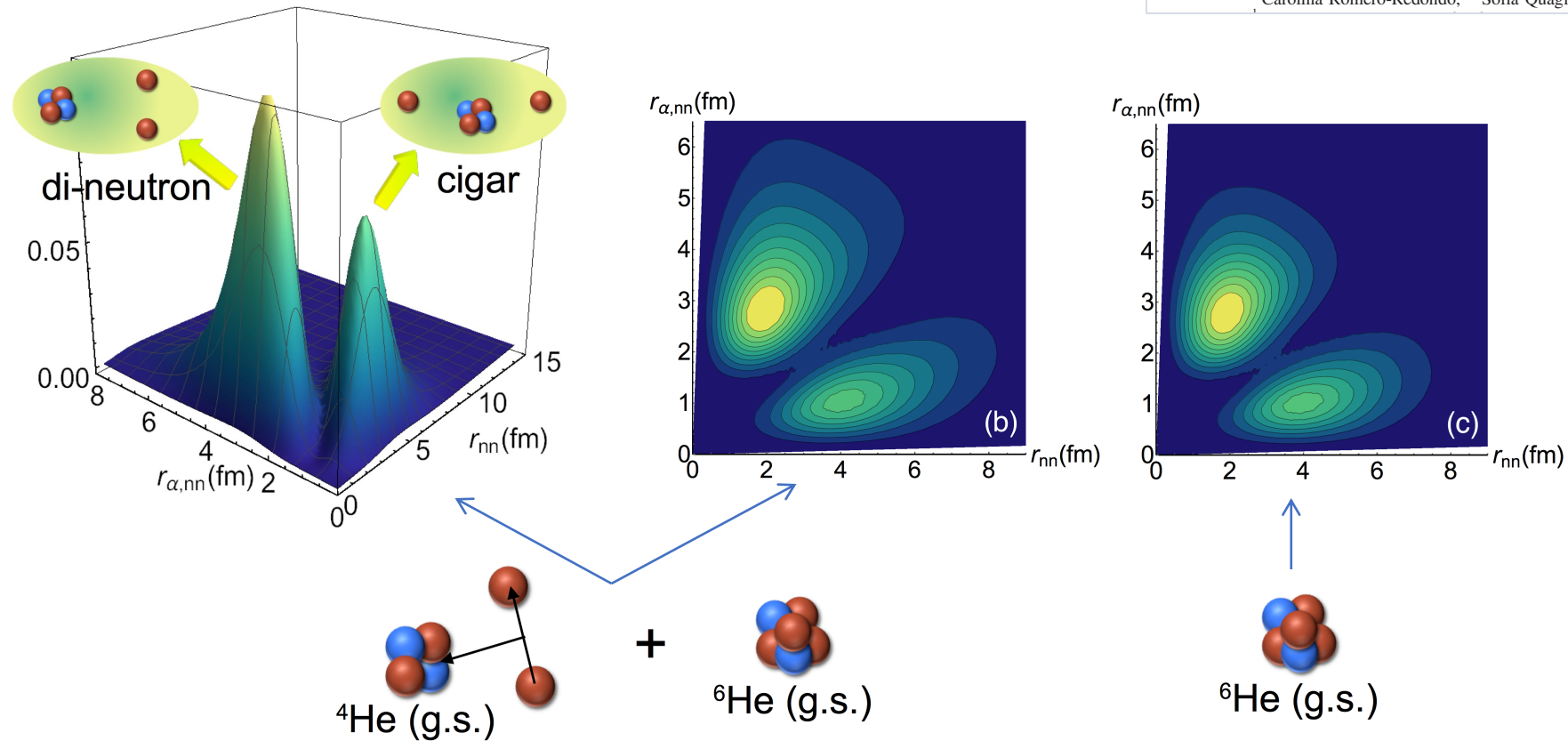
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How Many-Body Correlations and α Clustering Shape ${}^6\text{He}$

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NCSMC for three-body clusters: ${}^6\text{He} \sim {}^4\text{He} + n + n$



The probability distribution of the ${}^6\text{He}$ ground state presents two peaks corresponding to the di-neutron and cigar configurations

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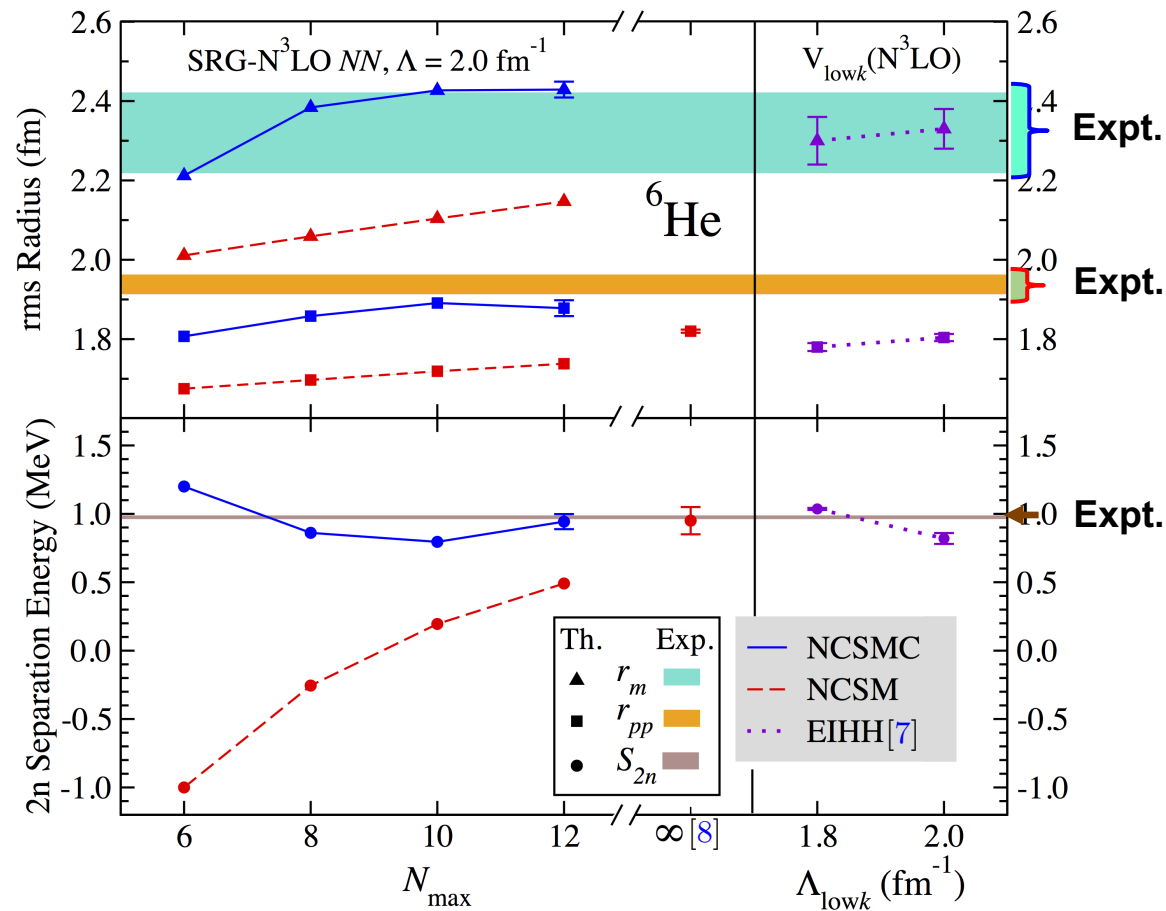
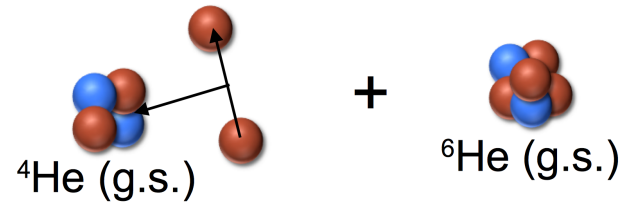
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Separation energy, point proton and matter radius simultaneously consistent with experiment

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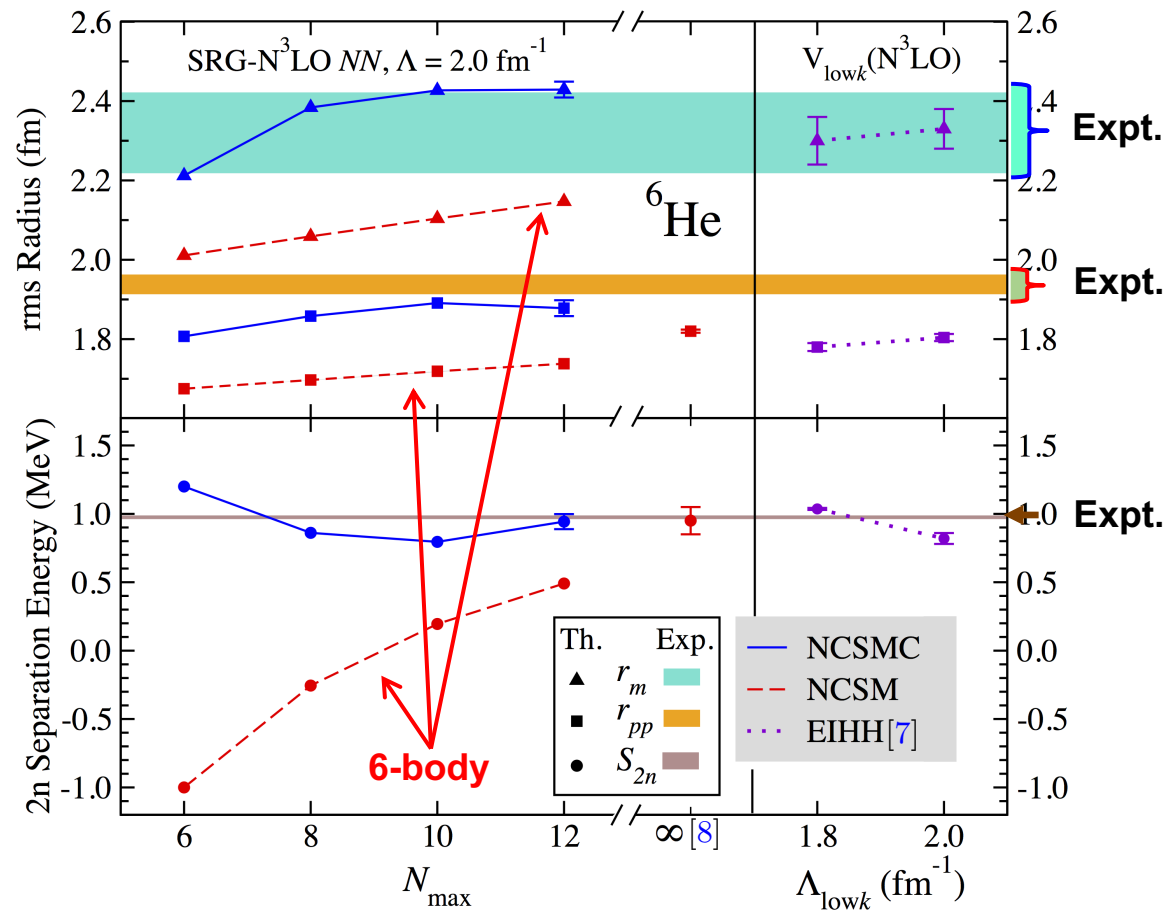
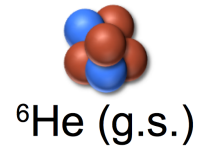
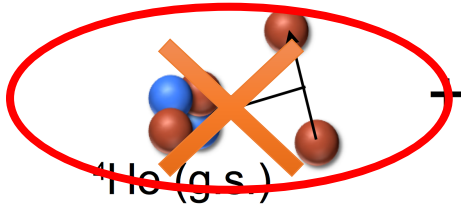
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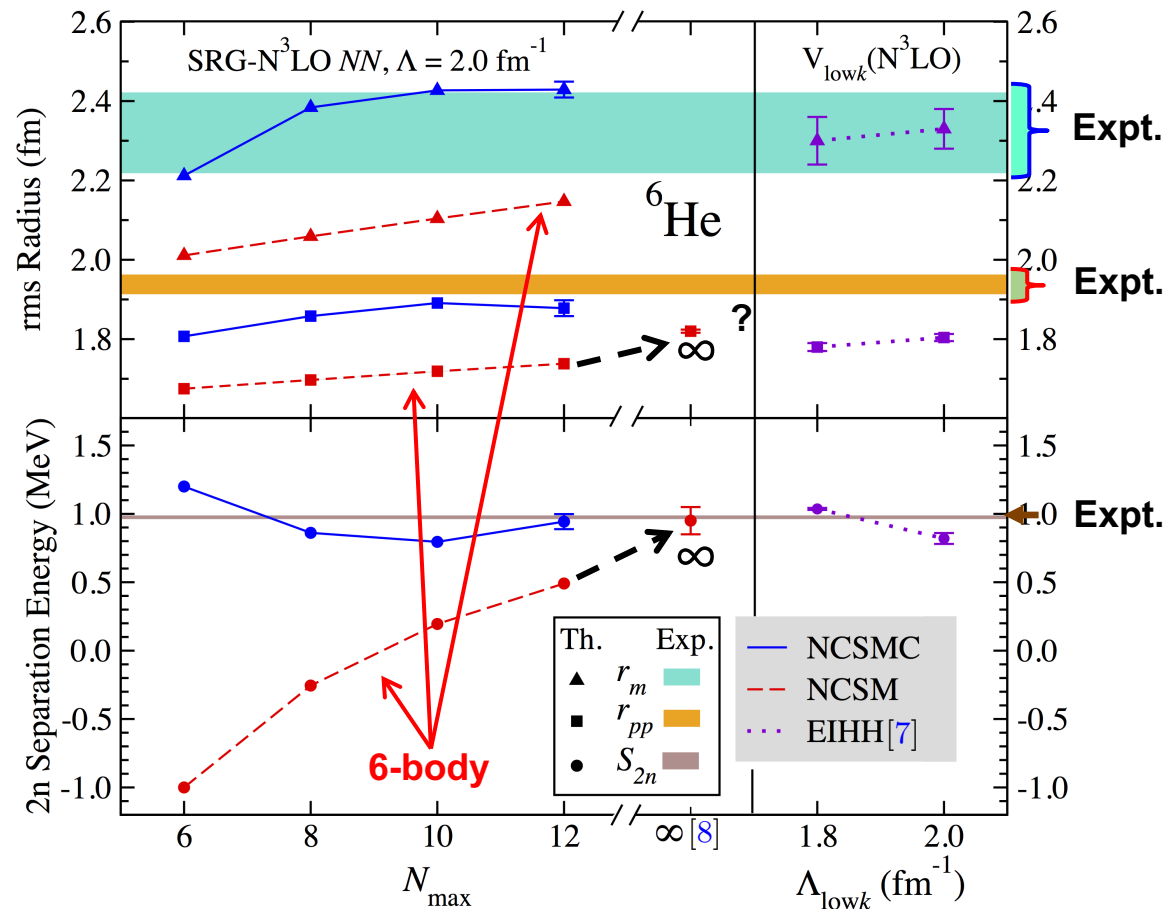
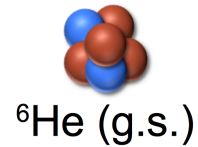
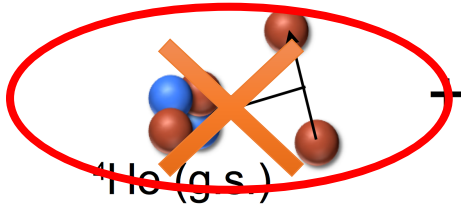
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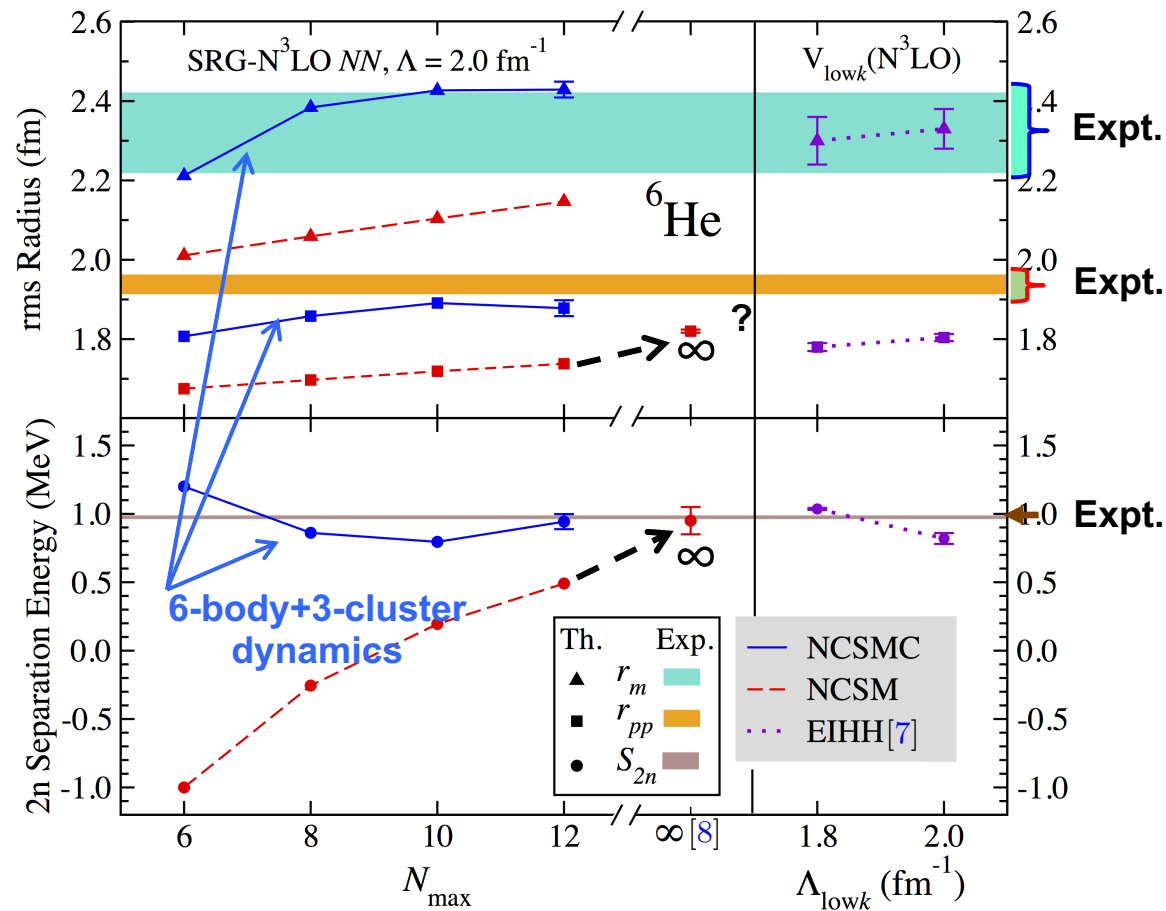
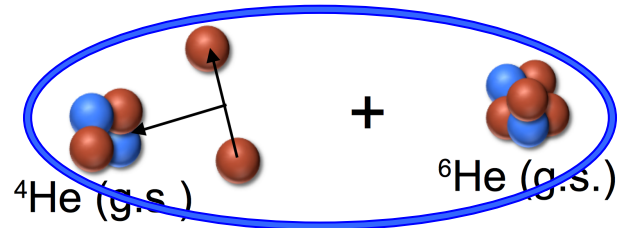
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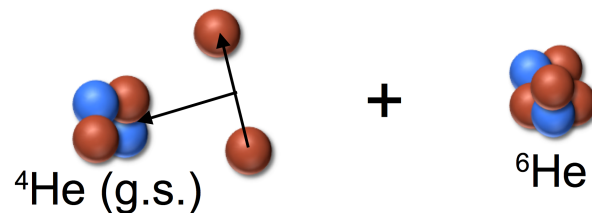
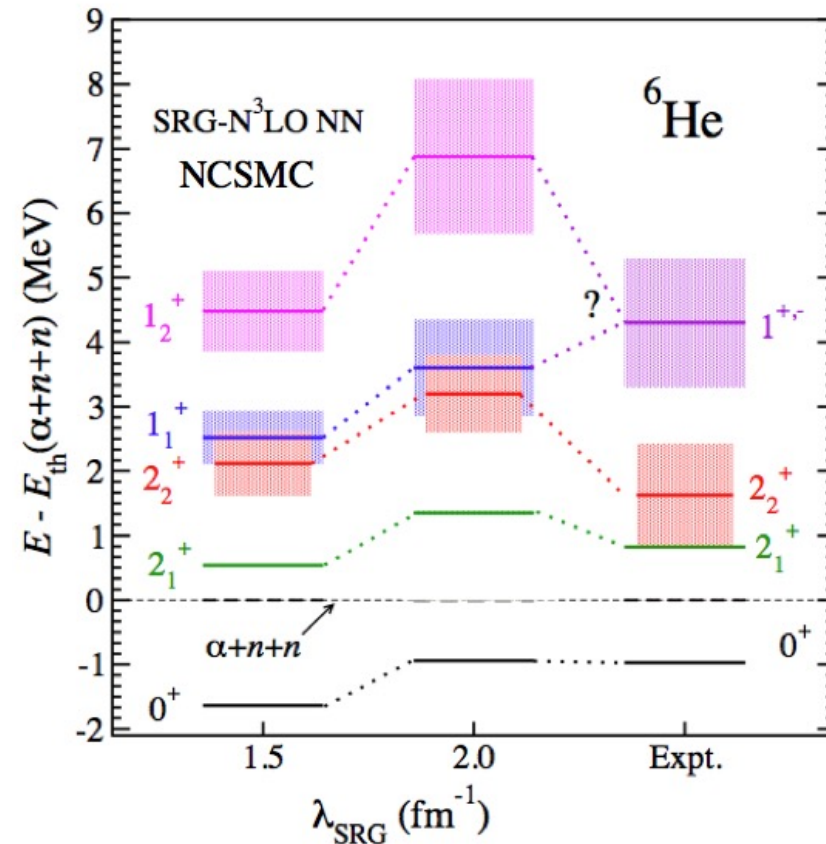
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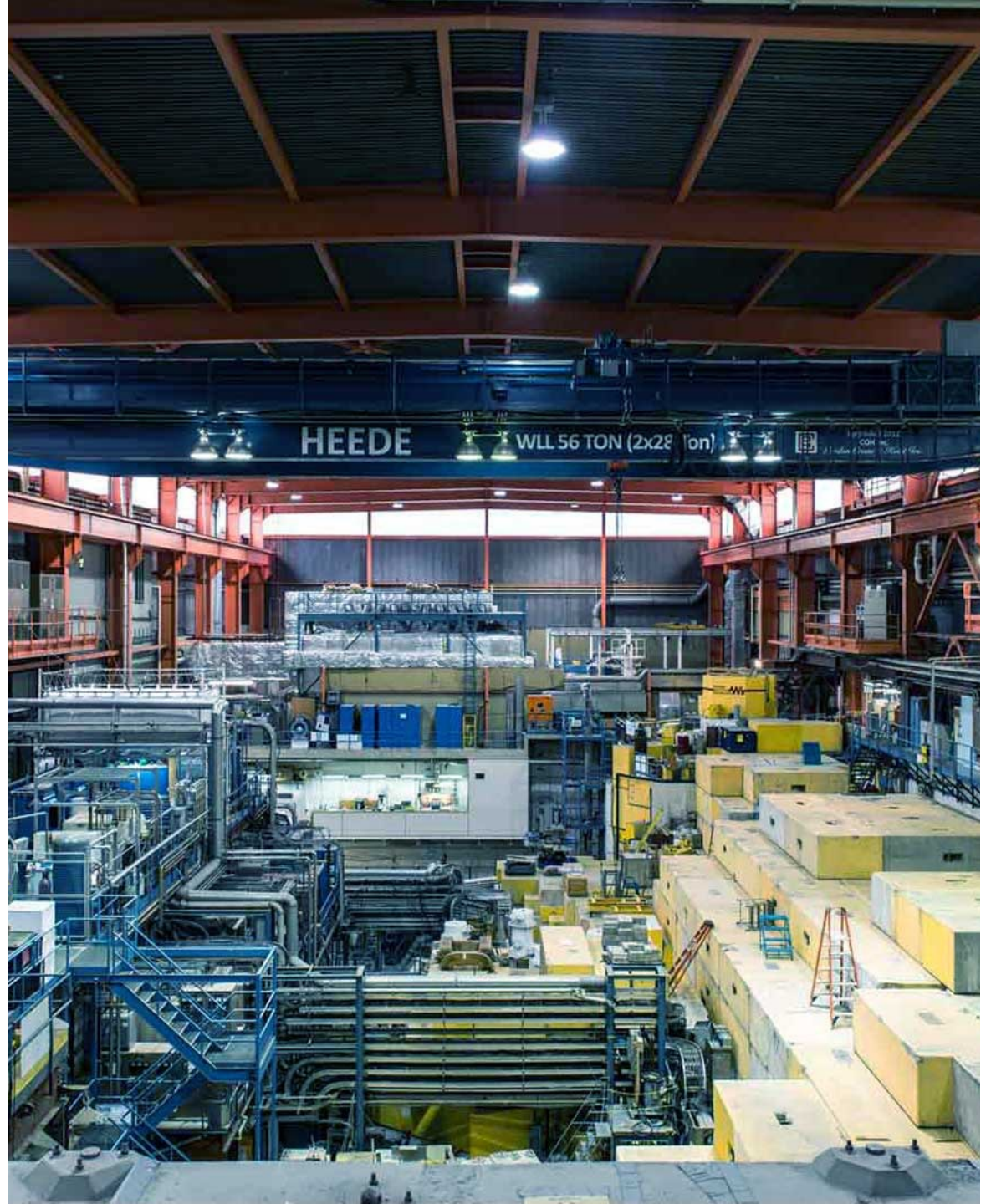
NCSMC for three-body clusters: ${}^6\text{He} \sim {}^4\text{He}+n+n$



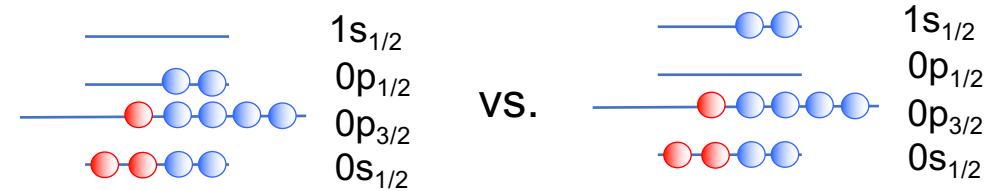
Quantitative comparison still requires
inclusion of 3N forces

^{11}Li within NCSM

(prerequisite for $^{11}\text{Li} \sim ^9\text{Li} + n + n$)

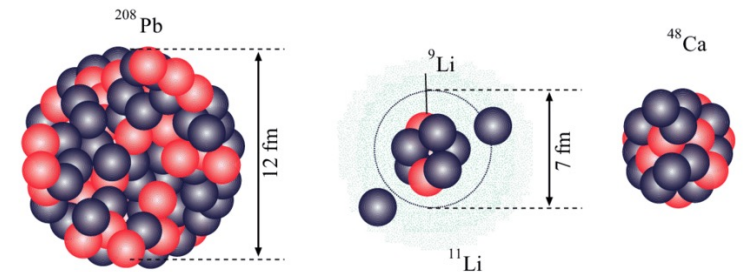
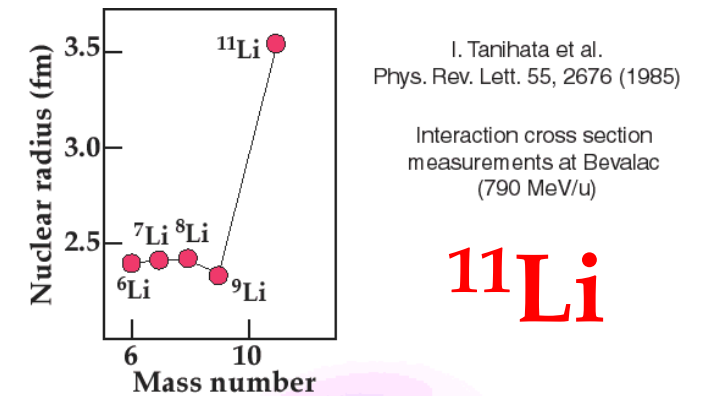


Borromean halo nucleus ^{11}Li



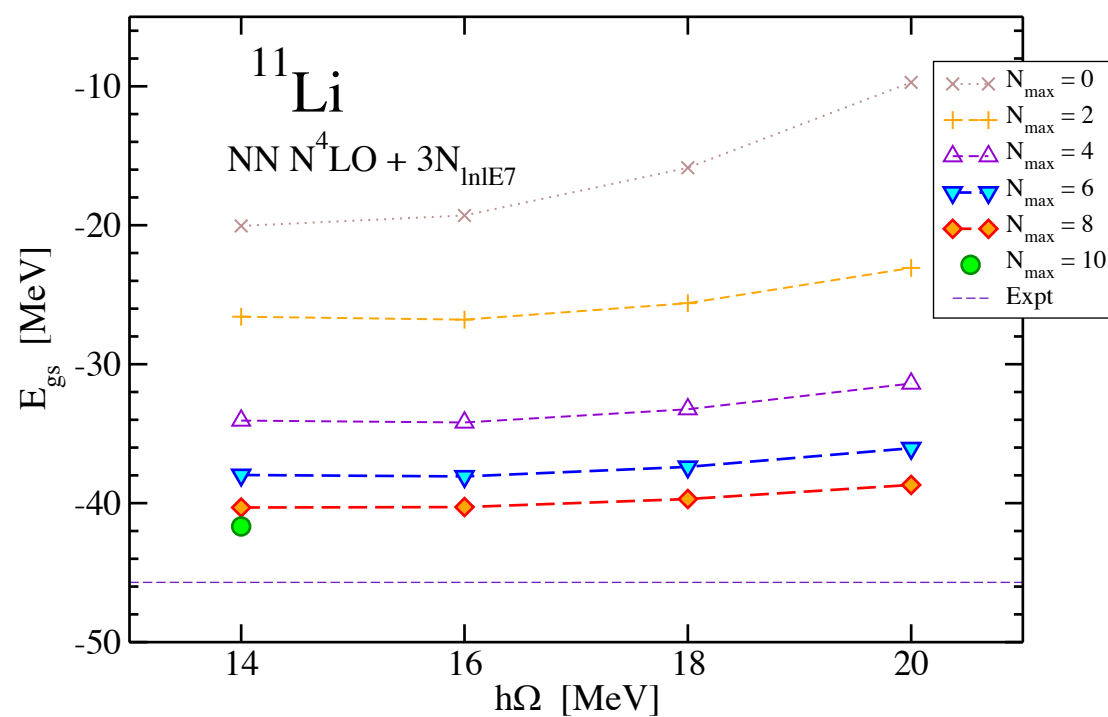
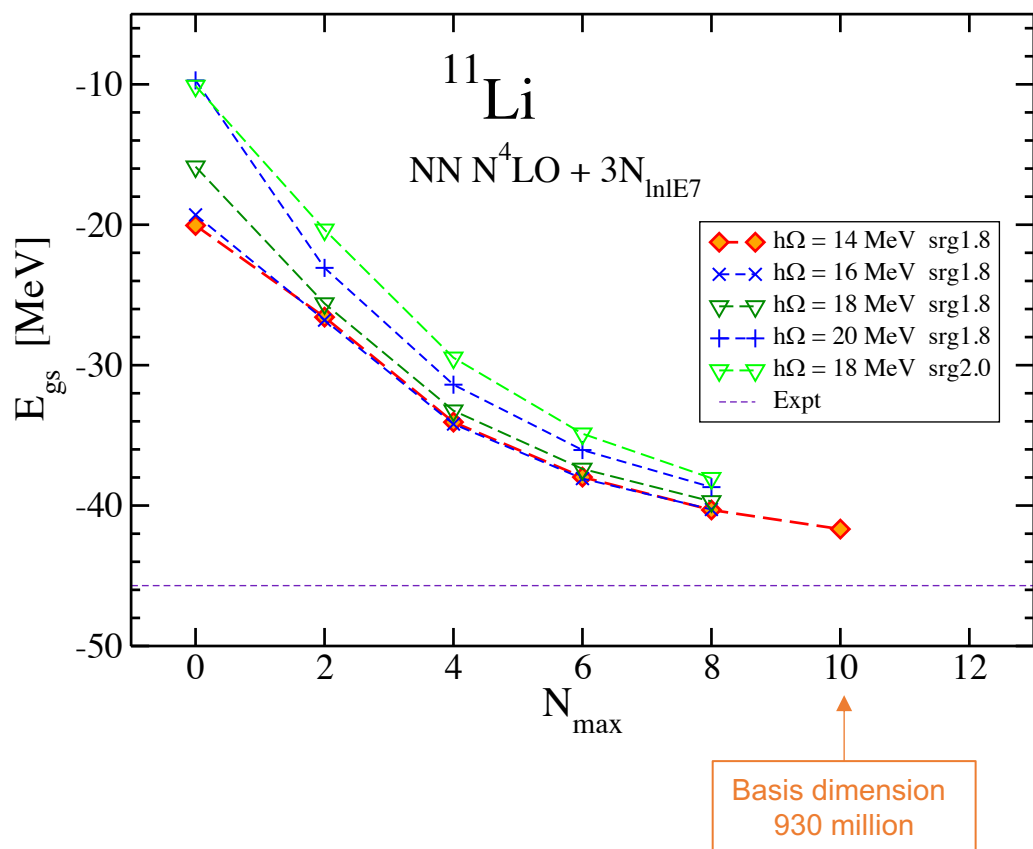
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- $Z=3, N=8$
 - Very weakly bound: $E_{\text{th}}=-0.3 \text{ MeV}$
 - Halo state – dominated by $^9\text{Li}+n+n$ in the S -wave
 - Configuration mixing, $^9\text{Li } 1/2^-$ excited state plays a role
- Can we describe ^{11}Li in *ab initio* calculations?
 - Continuum must be included
 - What role does the 3N interaction play in the configuration mixing?
 - NCSMC needs to be applied – very challenging
 - The first step – large-scale NCSM



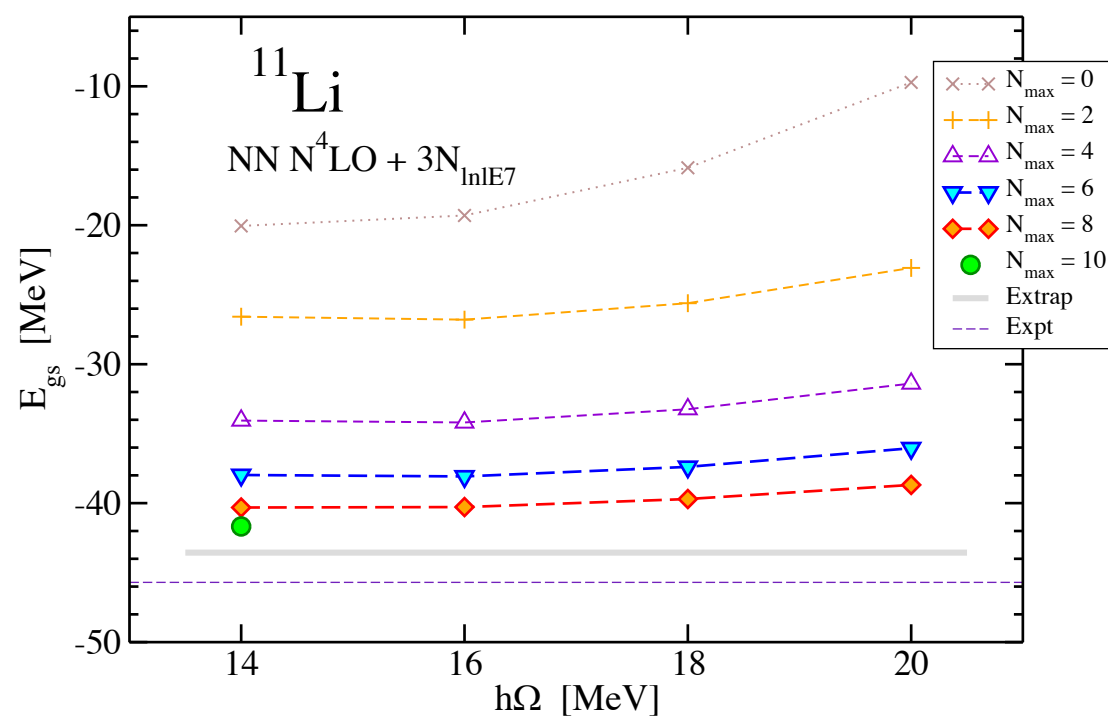
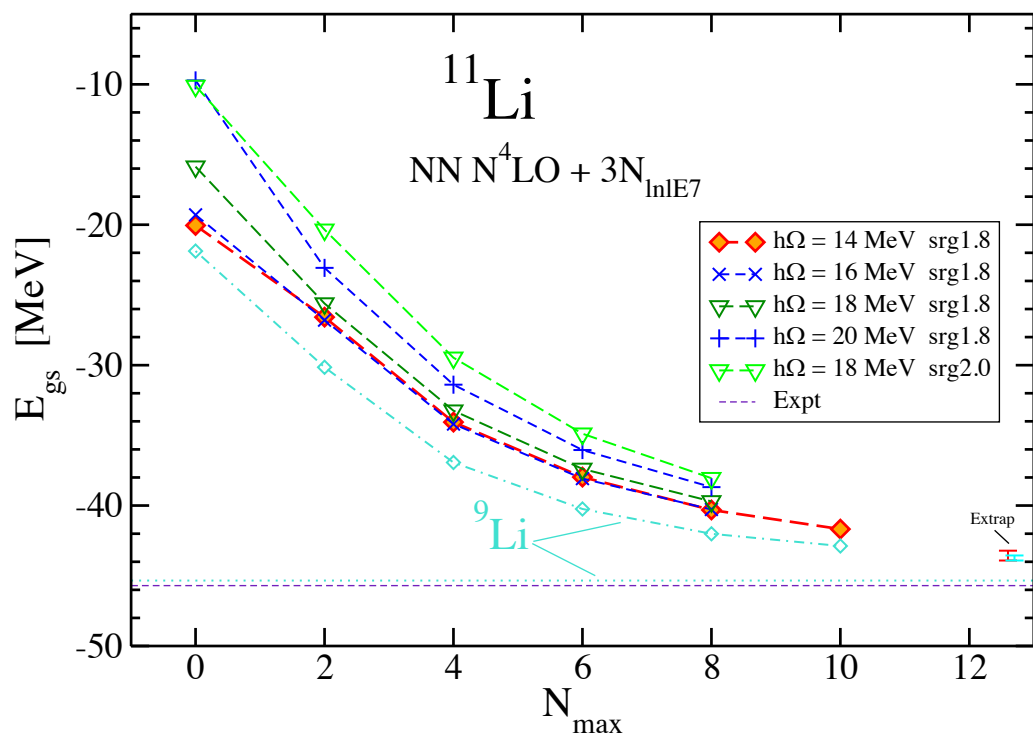
Large-scale NCSM calculations of ^{11}Li

40



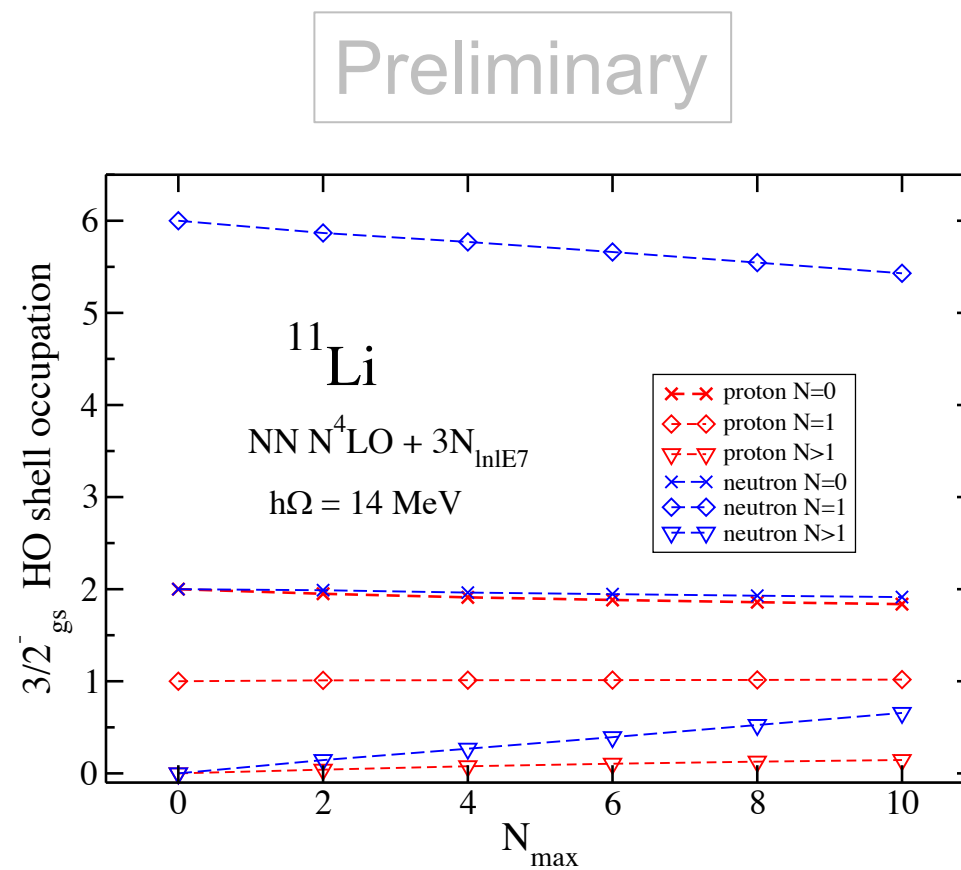
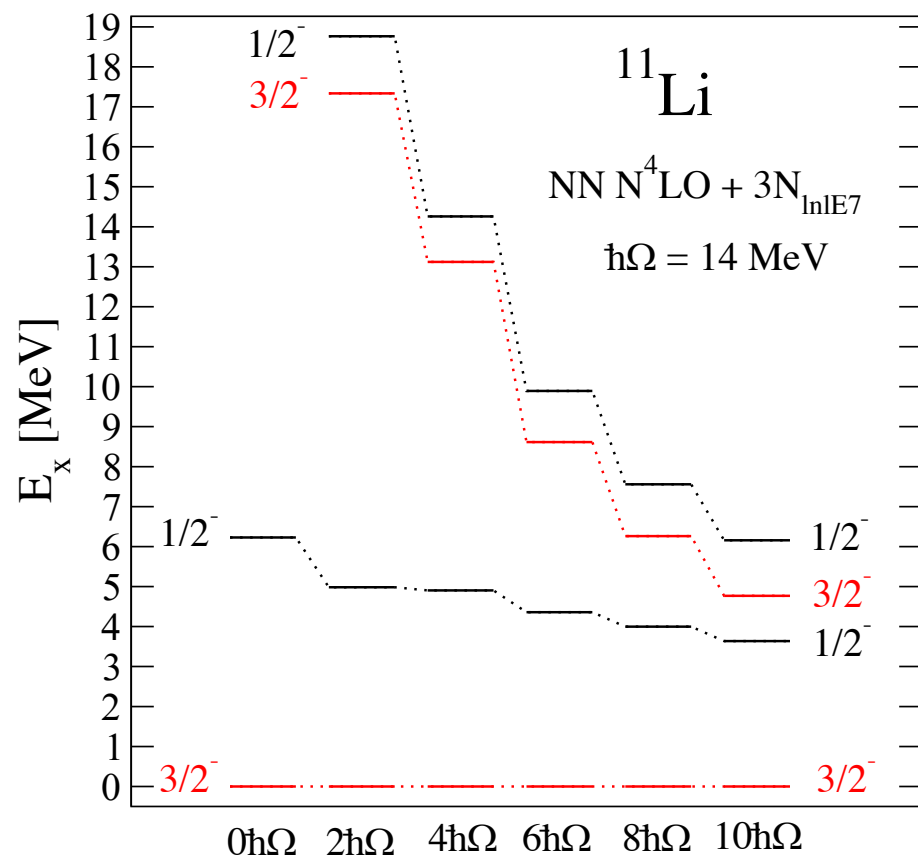
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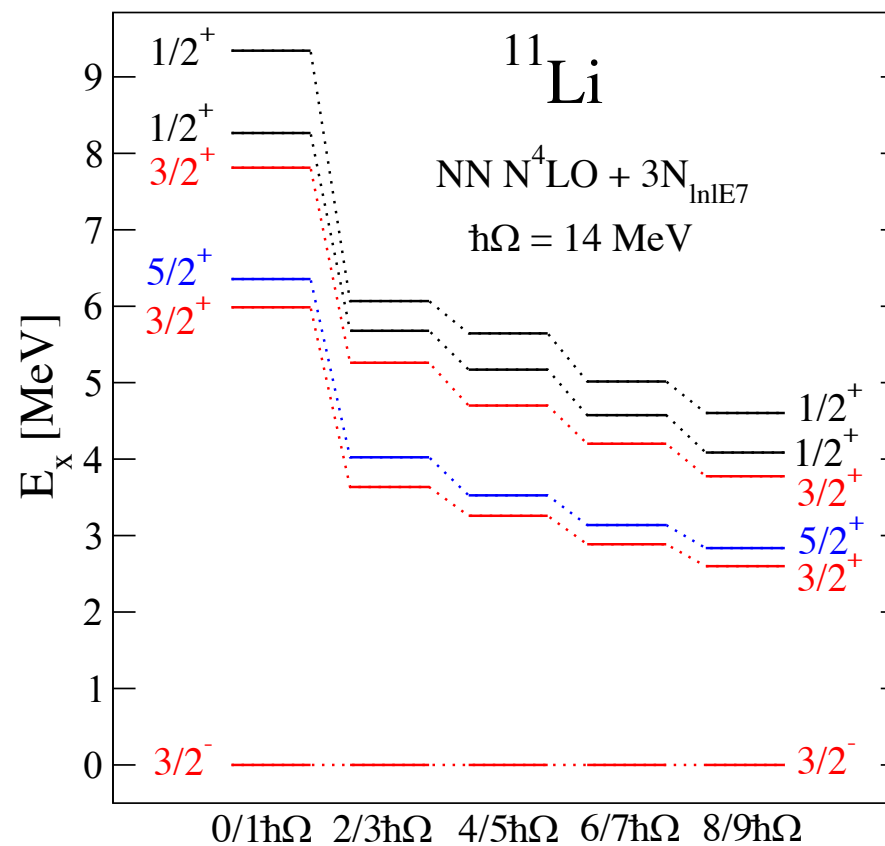
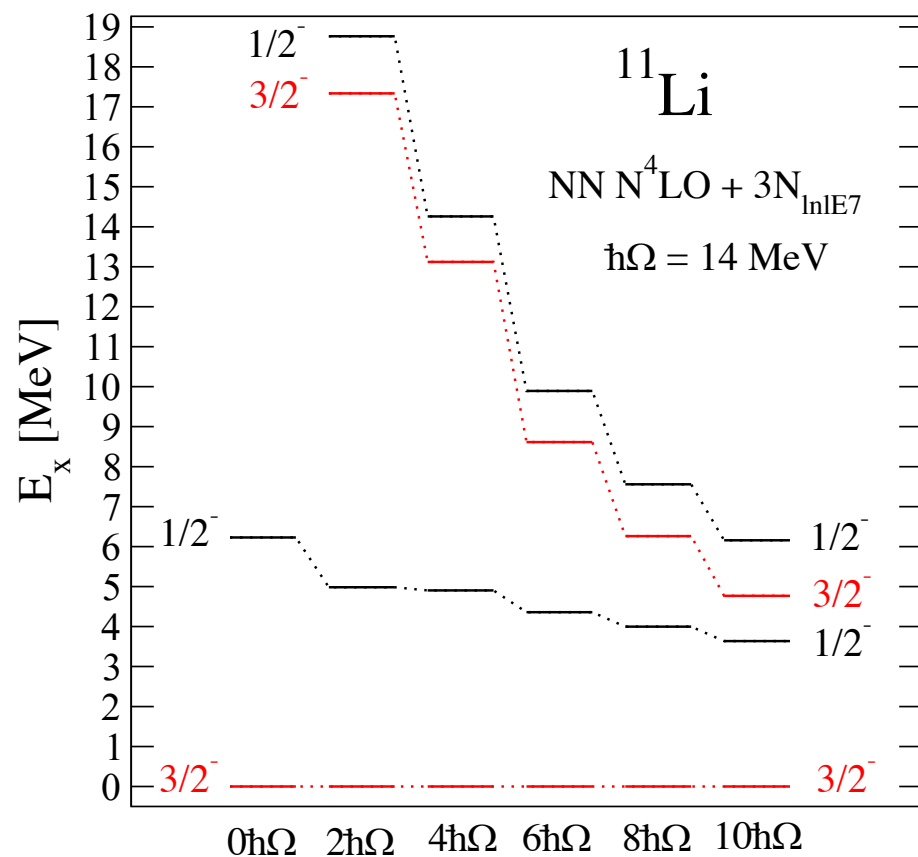
Large-scale NCSM calculations of ^{11}Li

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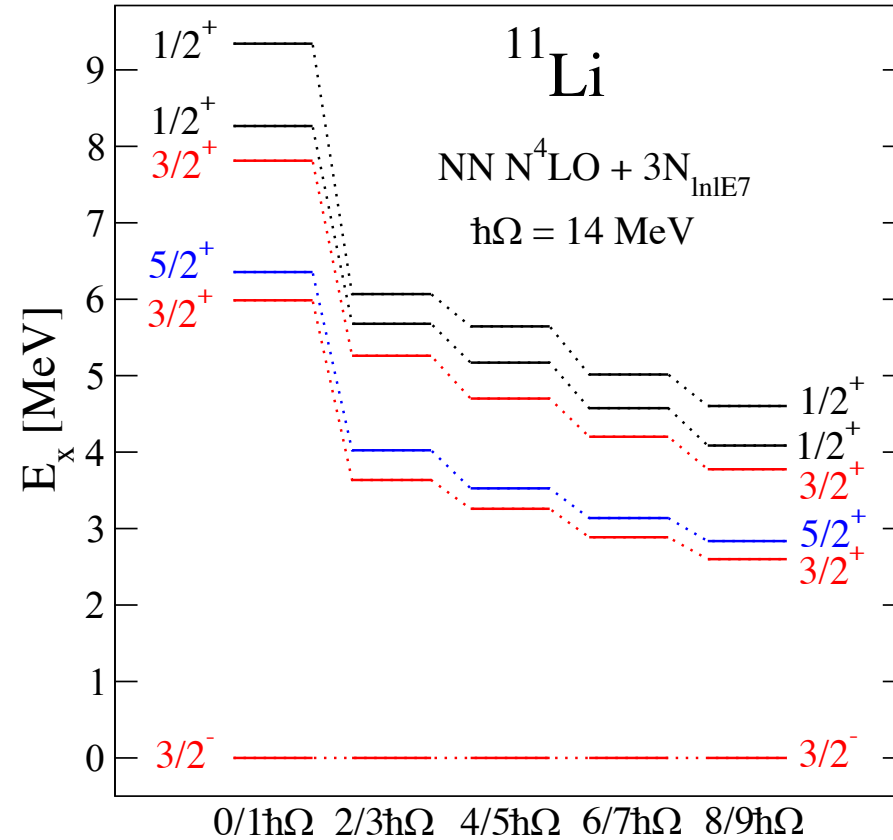
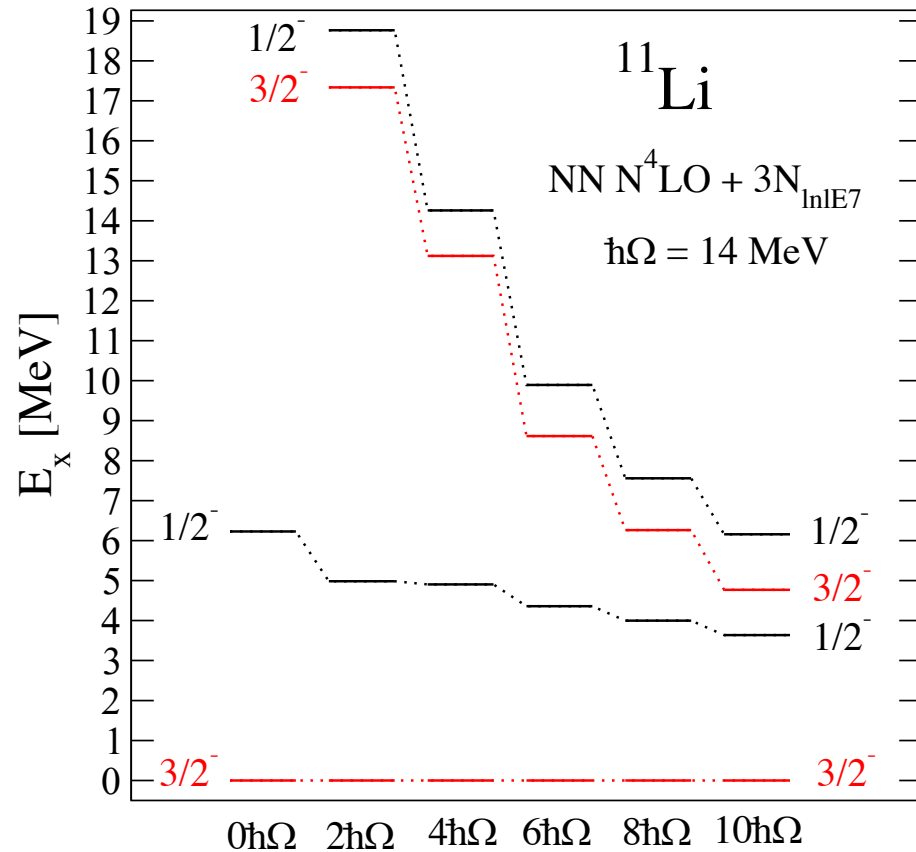


Large-scale NCSM calculations of ^{11}Li

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Large-scale NCSM calculations of ^{11}Li



Next step – including continuum via NCSMC – $^{11}\text{Li} \sim {}^9\text{Li} + n + n$

Conclusions

- *Ab initio* nuclear theory
 - Makes connections between the low-energy QCD and many-nucleon systems
- No-core shell model with continuum capable of describing halo nuclei from first principles
 - Successful reproduction of ^{11}Be parity inversion – three-nucleon interaction important
 - Ground-state halo of ^8B significantly impacts the proton capture on ^7Be
 - Three-body cluster NCSMC applied to the Borromean halo nucleus ^6He
- Next frontier – *ab initio* NCSMC description of ^{11}Li as a three-cluster $^9\text{Li} + n + n$ system