

Exchange correlation effects in nuclear novel phenomena

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Lanzhou University

International Symposium Commemorating the 40th Anniversary of the Halo Nuclei
October 12 - 18, 2025 | Capital Hotel, Beijing





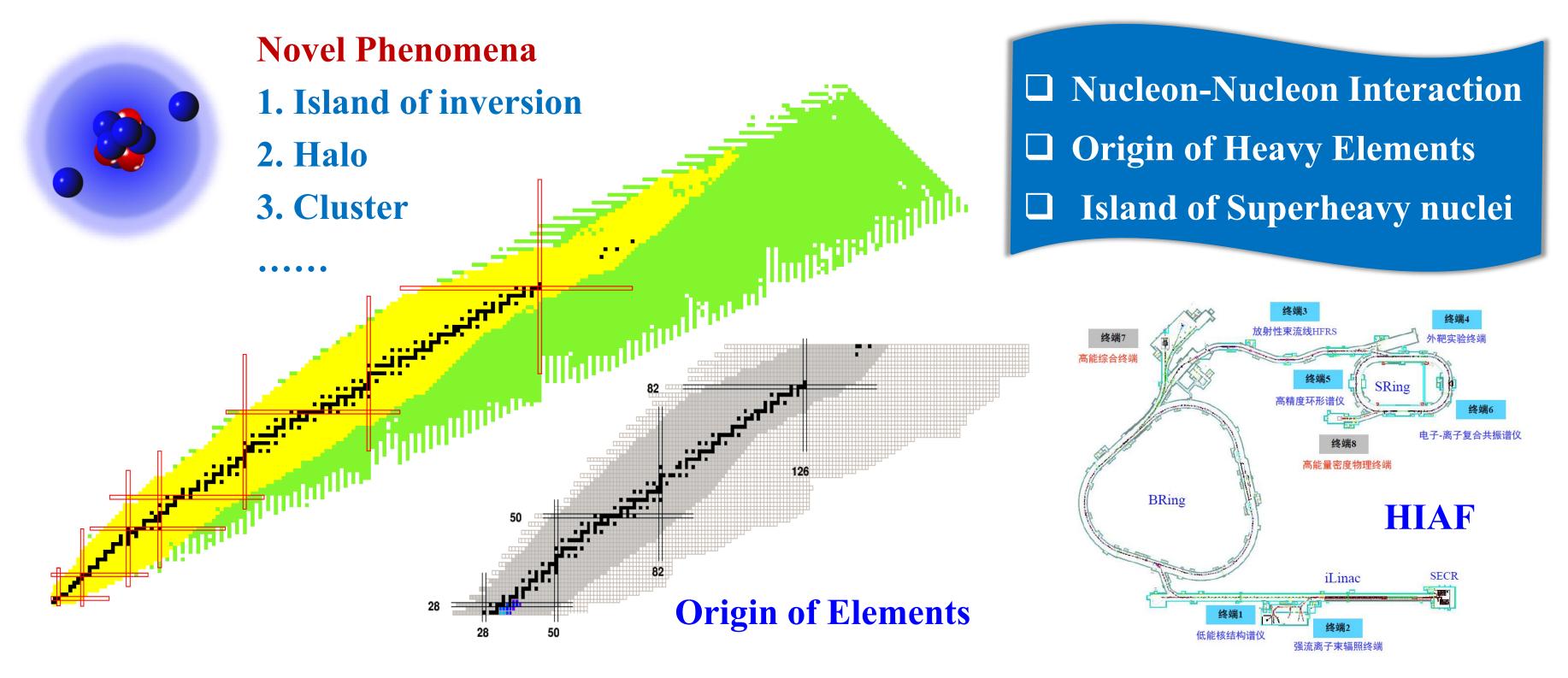
OUTLINE

- Introductions
 - Nuclear novel phenomena and the challenges
 - > Relativistic Hartree-Fock (RHF) theory and the extensions
- Recent developments of the RHF theory
 - > Deformed relativistic Hartree-Fock-Bogoliubov (D-RHFB) model
 - Configuration-Interaction relativistic Hartree-Fock (CI-RHF) model
- **■** Exchange correlation effects
 - > Spherical halo phenomena at the drip line of Ce isotopes
 - ➤ Deformed halo structure in ¹⁹C
 - Nature of the island of inversion exhibited by ³²Mg
- Conclusions



Unstable Nuclei





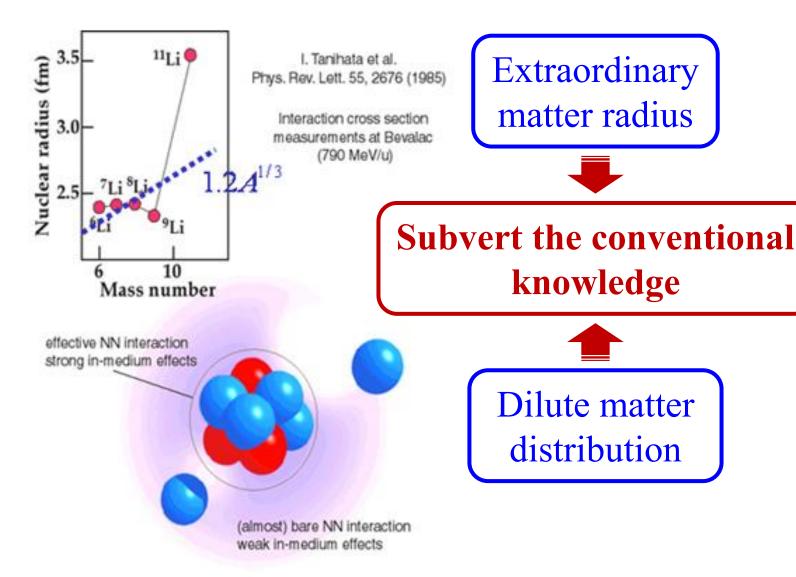
Call for comprehensive understanding on the nature of unstable nuclei



Halo and the challenges



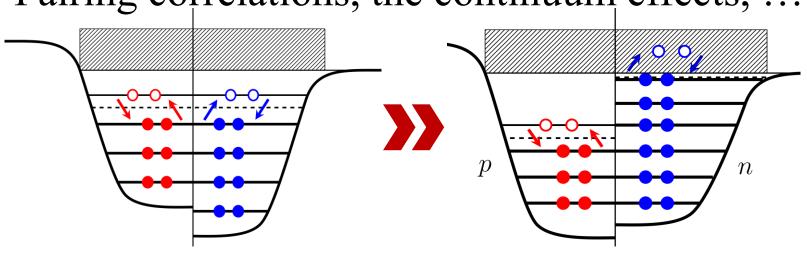
☐ The first candidate of halo nuclei: ¹¹Li



I. Tanihata et al., PRL **55** (1985) 2676

1. Weak binding mechanism

Pairing correlations, the continuum effects, ...



Meng et al., PPNP 57 (2006) 470

2. Nuclear deformation

Deformed Halo, cluster, island of inversion,...

3. Nuclear many-body state

Mean field approach, AMP, GCM, SM, ...

How nucleons spontaneously form an atomic nucleus

Nuclear force & Nuclear many-body method



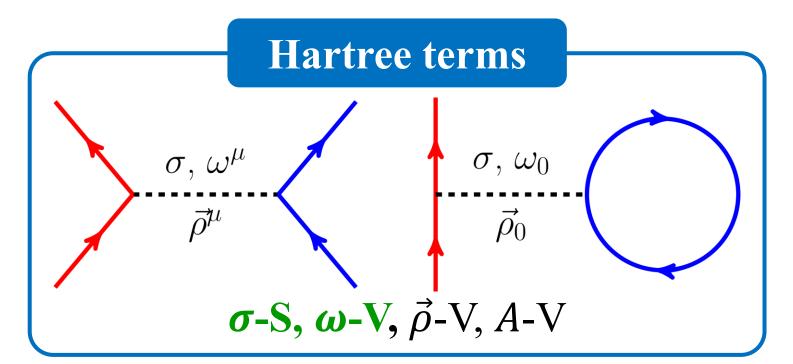
Relativistic Hartree-Fock (RHF) theory

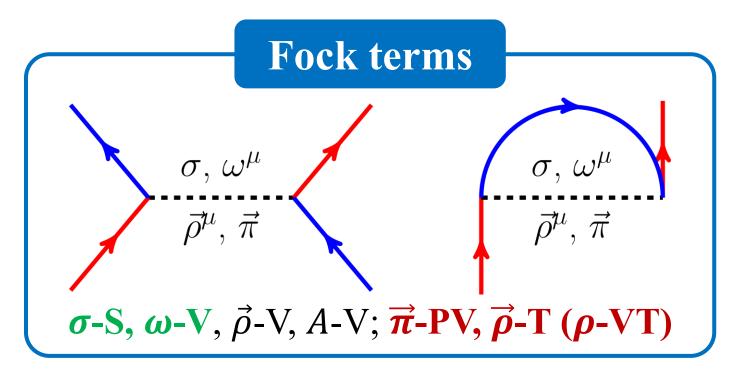


□ Relativity: Interplay between strong scalar attraction and vector repulsion

Natural spin-orbit coupling, the origin of pseudo-spin symmetry and spin symmetry, ...

Meng et al., PPNP 57, 470 (2006); Ginocchio, PRL 78, 436 (1997); Meng et al., PRC 58, R628 (1998); Zhou et al., PRL 91, 262501 (2003).





Fock terms: pure exchange degrees of freedom (DoF) such as π -PV and ρ -T (ρ -VT) Impact the modeling of nuclear binding, nuclear tensor force, ...

PRC 100, 051301(R) (2019)

PRC 91, 034326 (2015); CPC 42, 024101 (2018); PRC 98, 034313 (2018); PRC 101, 064306 (2020)

Exchange correlation effects in nuclear novel phenomena?

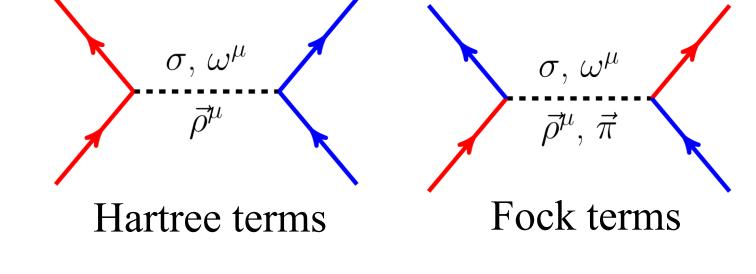


Effective RHF Lagrangians

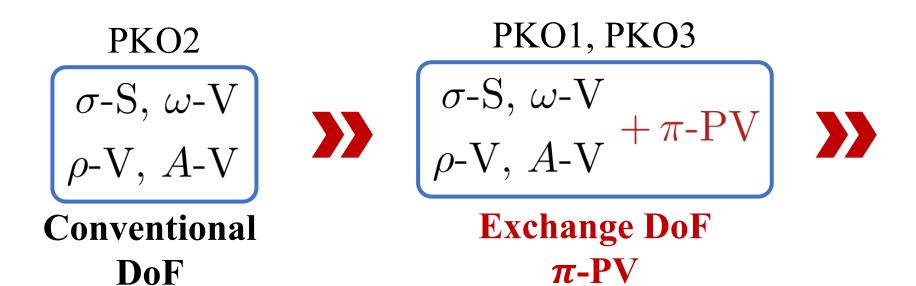


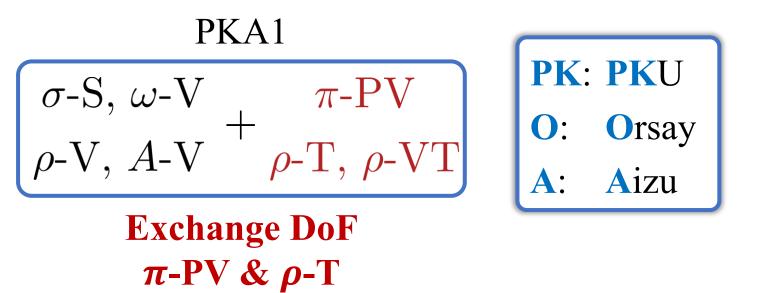
- Nuclear in-medium effects
 - 1. Non-linear self-couplings of meson fields
 - 2. Density-dependent coupling strengths

 Effects of many-body correlations in nuclei
- RHF Lagrangians PKO*i* and PKA1 similar quantitative accuracy as popular DFT



WHL, Giai, Meng, PLB **640**, 150 (2006); WHL, Sagawa, Giai, Meng, PRC **76**, 034314 (2007); EPL **82**, 12001 (2008).





Pave a way to reveal the underlying effects of the exchange correlations

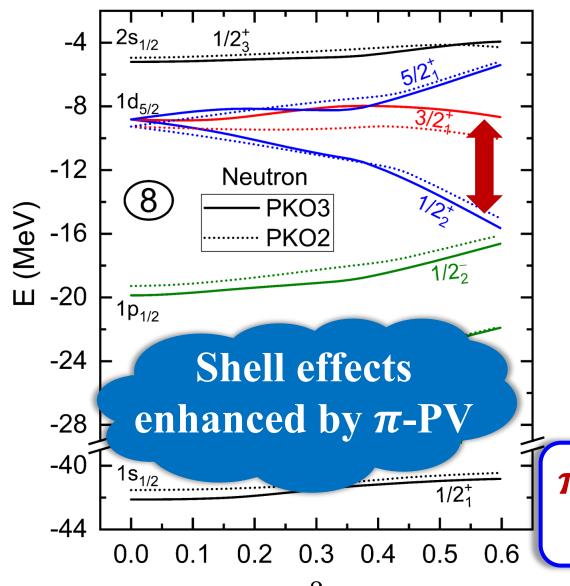


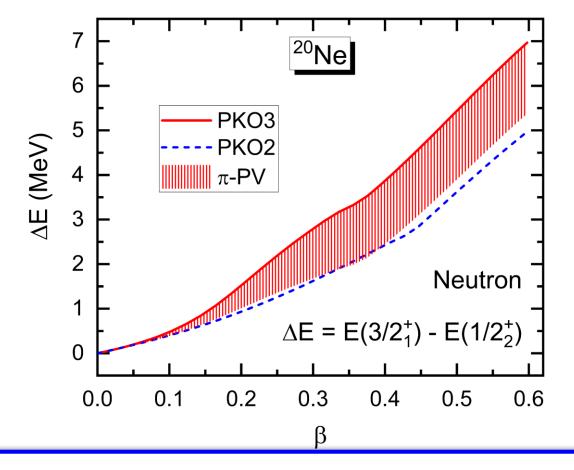
Axially deformed RHF (D-RHF) model: deformation



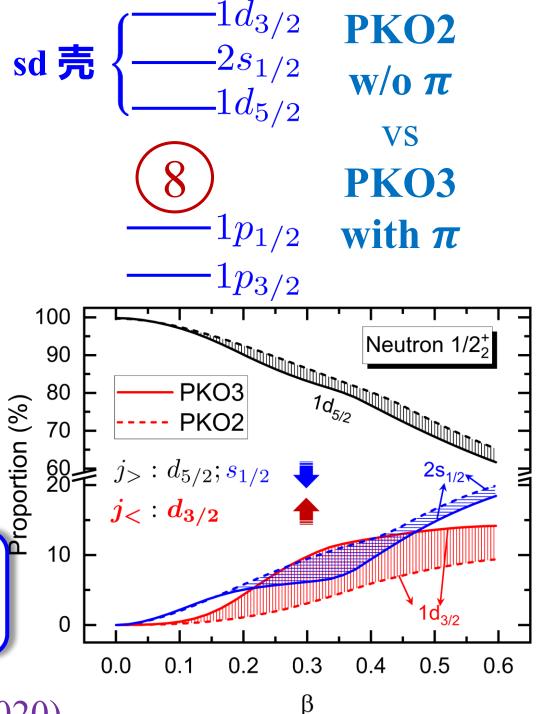
☐ Spherical Dirac Woods-Saxon (DWS) basis

- Zhou, Meng, Ring, PRC 68, 034323 (2003)
- > Appropriate asymptotic behaviors of wave functions
- > Microscopic insight into deformed nuclei: tensor force effects





 π -T: more $1d_{3/2}$ but less $1d_{5/2}$ waves mixing in $1/2_2^+$ orbit



Geng, Xiang, Sun, WHL, PRC 101, 064302 (2020)



Axially deformed RHFB (D-RHFB) model



Bogoliubov transformation: Bloch-Messiah decomposition

$$\overset{\text{HF}}{\text{s.p. space}} (c, c^{\dagger}) \xrightarrow{D} (a, a^{\dagger}) \xrightarrow{\overline{U}} \overset{\overline{U}}{\overline{V}} \xrightarrow{\overline{V}} (\alpha, \alpha^{\dagger}) \xrightarrow{C} (\beta, \beta^{\dagger}) \overset{\text{Bogoliubov}}{\text{q.p. space}}$$



$$\psi(x) = \sum_{k} \left[\psi_k^U(x) \beta_k + \psi_{\widetilde{k}}^V(x) \beta_k^{\dagger} \right] \qquad H = \widehat{T} + \sum_{\phi} \widehat{V}_{\phi}$$

$$H = \widehat{T} + \sum_{\phi} \widehat{V}_{\phi}$$

Continuum Pairing

$$p$$

HFB g.s.: β |HFB \rangle =0

$$\widehat{T} = \int d\boldsymbol{x} \sum_{kk'} \bar{\psi}_{k}^{V} (-i\boldsymbol{\gamma} \cdot \boldsymbol{\nabla} + M) \psi_{k'}^{V} \beta_{k} \beta_{k'}^{\dagger} \qquad \text{Geng, WHL, PRC 105, 034329 (2022)}$$

$$\widehat{V}_{\phi} = \frac{1}{2} \iint d\boldsymbol{x}_{1} d\boldsymbol{x}_{2} \sum_{k_{1}k'_{1}} \sum_{k_{2}k'_{2}} \bar{\psi}_{k_{1}}^{V} (x_{1}) \bar{\psi}_{k_{2}}^{V} (x_{2}) \Gamma_{\phi} D_{\phi} \psi_{k'_{2}}^{V} (x_{2}) \psi_{k'_{1}}^{V} (x_{1}) \beta_{k_{1}} \beta_{k_{2}} \beta_{k'_{2}}^{\dagger} \beta_{k'_{1}}^{\dagger} \qquad \text{Mean field}$$

$$+ \frac{1}{2} \iint d\boldsymbol{x}_{1} d\boldsymbol{x}_{2} \sum_{k_{1}k'_{1}} \sum_{k_{2}k'_{2}} \bar{\psi}_{k_{1}}^{V} (x_{1}) \bar{\psi}_{k_{2}}^{U} (x_{2}) \Gamma_{\phi} D_{\phi} \psi_{k'_{2}}^{U} (x_{2}) \psi_{k'_{1}}^{V} (x_{1}) \beta_{k_{1}} \beta_{k_{2}}^{\dagger} \beta_{k'_{2}}^{\dagger} \beta_{k'_{1}}^{\dagger} \qquad \text{Pairing correlations}$$

Unified treatment: spin-orbit and tensor forces, deformation, pairing correlations & the continuum



Configuration-Interaction RHF (CI-RHF) model



■ Effective Hamiltonian based on RHF model

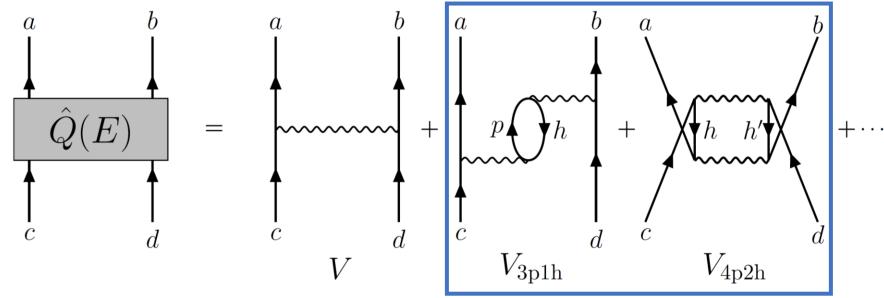
$$H = E_{\text{core}} + \sum_{ii'} \varepsilon_{ii'} \cdot c_i^{\dagger} c_{i'} \cdot + \frac{1}{4} \sum_{iji'j'} \overline{V}_{iji'j'} \cdot c_i^{\dagger} c_j^{\dagger} c_{j'} c_{i'} \cdot$$

EKK scheme

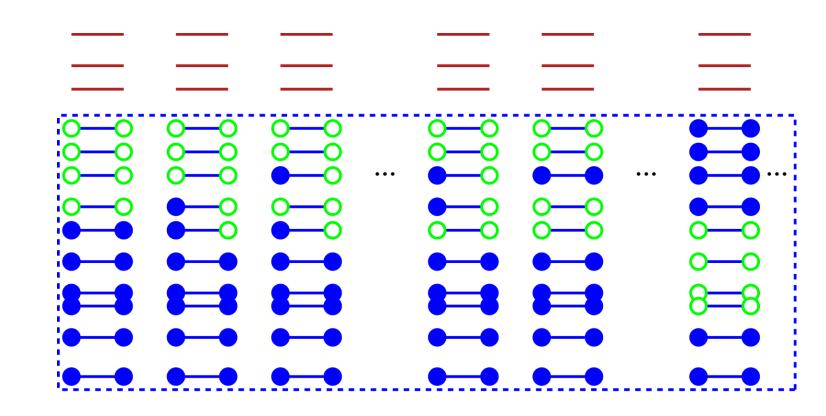
Introductions

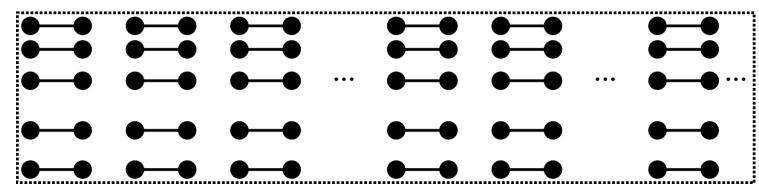
$$H_l^{\text{eff}} = H^{\text{BH}}(E_0) + \sum_{k=1}^{\infty} \frac{1}{k!} \frac{d^k \widehat{Q}(E_0)}{dE_0^k} \left\{ H_{l-1}^{\text{eff}} - E_0 \right\}^k$$
 Folded terms

$$\hat{Q}\text{-box} \atop \textbf{expansion} \ \hat{Q}(E) = \sum_{l=0}^{\infty} P \Big[V \frac{Q}{E-H_0} \Big]^l VP \\ \textbf{Core polarization}$$



J. Liu, Y.F. Niu, WHL, CPC 49, 064104 (2025)





Optimized 1p-1h HF basis excitation

2p-2h excitation

np-nh excitation

$$|\Psi\rangle = \sum_{n} C_n |\Phi_n\rangle , \quad |\Phi_n\rangle = \prod_{i=1}^{A-N_{\text{core}}} c_{n_i}^{\dagger} |\text{core}\rangle$$

Brown et al. PLB **695** (2011) 507...



CI-RHF model: ¹⁸O and Ne isotopes



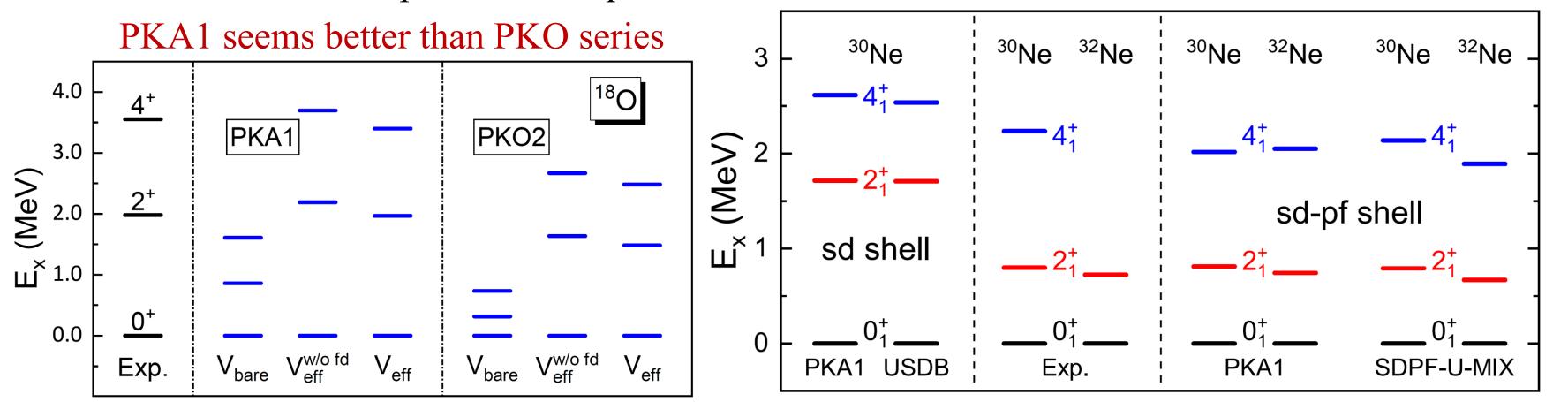
■ Without introducing any new parameters, the effective Hamiltonian constrained by the existing RHF Lagrangians works well for describing the low-lying excitations

Folded terms & core polarizations

are necessitated for precise description

Extending from sd shell to sd-pf shell

J. Liu, Y.F. Niu, WHL, CPC 49, 064104 (2025)



An efficient tool to describe the g.s. and low-lying excitations of unstable nuclei



Exchange correlation effects at the RHF level



 \square Exchange degrees of freedom (π -PV and ρ -T) change the modeling of nuclear binding

RMF (DD-ME2) \longrightarrow RHF (PKO3: plus π -PV) \longrightarrow RHF (PKA1: plus π -PV & ρ -T)

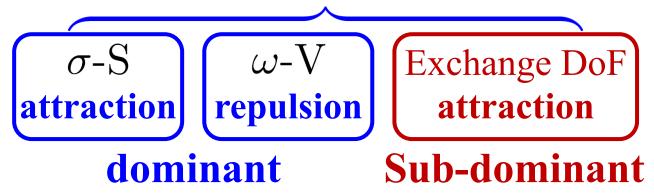
Modeling of nuclear binding for ²⁰⁸Pb

Models	$E_{\rm kin.+c.m.}$	$E_{\sigma+\omega}$	$E_{\rho+\pi}$	$E_{\rm cou.}$	Total
DD-ME2	2629.1	-5194.0	+100.3	827.2	-1637.4
PKO3	2823.2	-4609.7	-648.7	798.4	-1636.8
PKA1	2498.9	-3298.9	-1634.8	798.6	-1636.3

Fock terms, particularly the exchange degrees of freedom (DoF), impact the nuclear binding,

Geng, Li, WHL, Niu, Chang, PRC 100, 051301(R) (2019)

Picture of nuclear binding

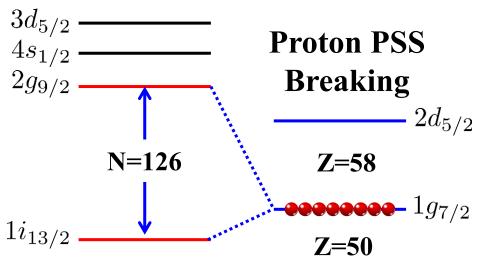


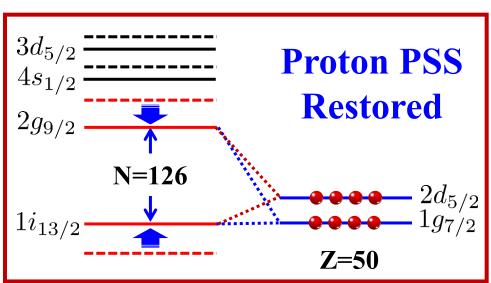


Giant halo in Ce isotopes

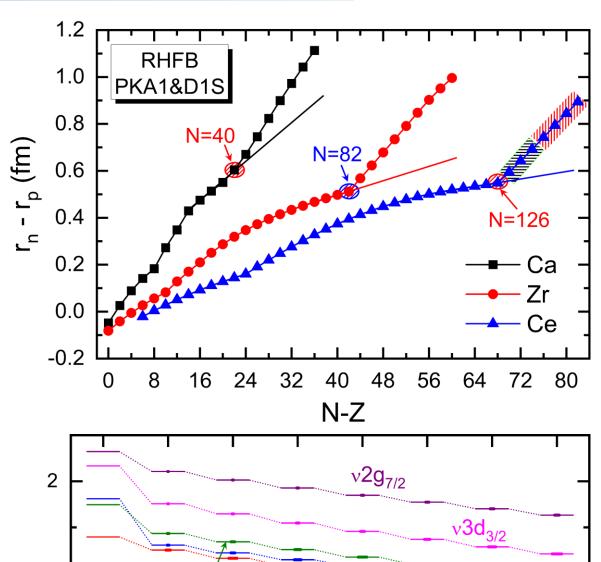


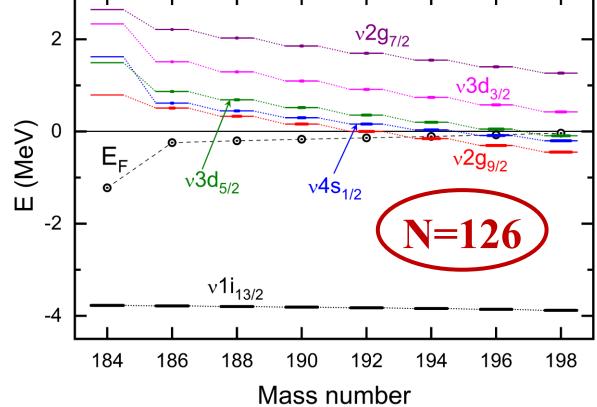
■ Exchange DoFs enhance the link between the restored proton PSS and the squeezed N=126 shell





	a	$= \pi 2d_{5/2}$	2	a	$a=\pi 1g_{7}$	/2	N=126
b	V_{ab}	ρ-Τ	π-PV	ρ-Τ	π -PV	V_{ab}	squeezed
$\nu 1 i_{13/2}$	-0.257	25.1%	2.1%	18.5%	9.8%	-0.386	
$-\nu 4s_{1/2}$	-0.097	17.1%	3.9%	20.2%	3.5%	-0.052	Nodal effects
$- u 2g_{9/2}$	-0.217	24.5%	3.0%	19.0%	8.8%	-0.157	CITCUS





WHL, Ring, Meng, Giai, Bertulani, PRC **81**, 031302(R) (2010)



Continuum effects in the halo structure of ¹⁹C

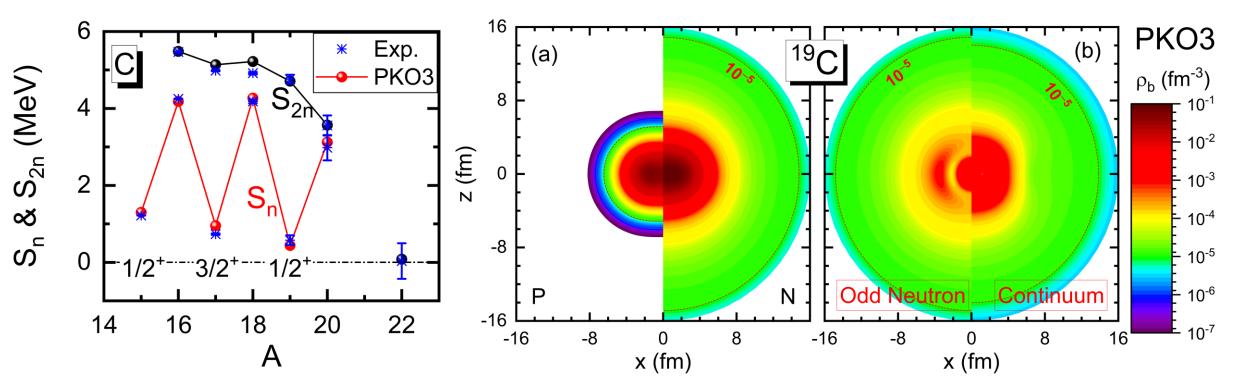


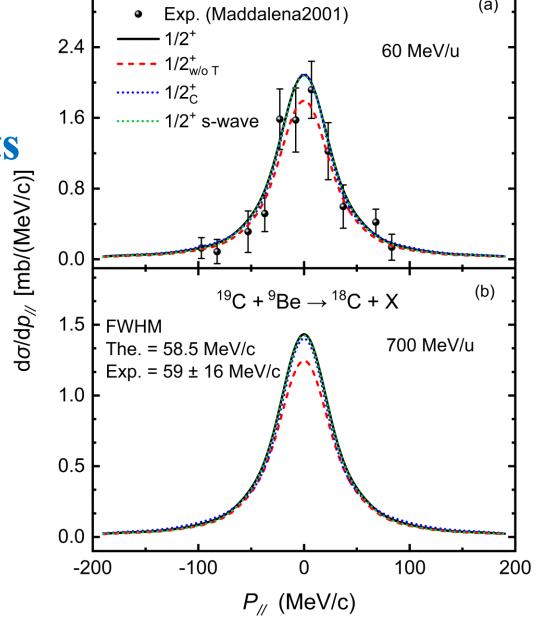
 \Box ¹⁹C: single-neutron halo, spin-parity $J^{\pi} = 1/2^+$

D-RHFB structure & Glauber reaction models

Structure & reaction observables: critical continuum effects

J. L. An, H. C. Li, J. Geng, WHL, S. S. Zhang, JPG 52, 055106 (2025)





Structure: spin-parity, neutron separation energy, density

Reaction: longitude momentum

☐ D-RHFB model can provide reliable structure input for the reactions of unstable nuclei



The island of inversion nuclide: ³²Mg



- □ Nature of the island of inversion exhibited by ³²Mg
 - \triangleright Neutron magic shell N = 20 vanishes, accompanying with notable g.s. deformation
 - > Rotational-like g.s. band and the co-existence of prolate and spherical shapes

Motobayashi, Ikedaa *et al.*, PLB **346**, 9 (1995); Pritychenko, Glasmacher *et al.*, PLB **461**, **322** (1999). Church, Campbell *et al.*, PRC **72**, 054320 (2005); Wimmer, Kroll *et al.*, PRL **105**, 252501 (2010).

- ☐ Shell model with the SDPF-M interaction
 - A remarkable cross-shell excitation from sd- to pf-shells is revealed for deformed g.s. 0_1^+

Caurier, Nowacki, Poves, PRC 90, 014302 (2014); Kitamura, Wimmer, Poves et al., PLB 822, 136682 (2021).

Angular momentum projection based on the mean-field models: intrusion of *pf*-shell **Reproduce the deformed g.s.** but fail in describing the low-lying excitation energies

Rodríguez-Guzmán, Egido, Robledo, NPA **709**, 201 (2002); Vaquero, Egido, Rodríguez, PRC **88**, 064311 (2013). Yao, Meng, Ring *et al.*, PRC **79**, 044312 (2009).

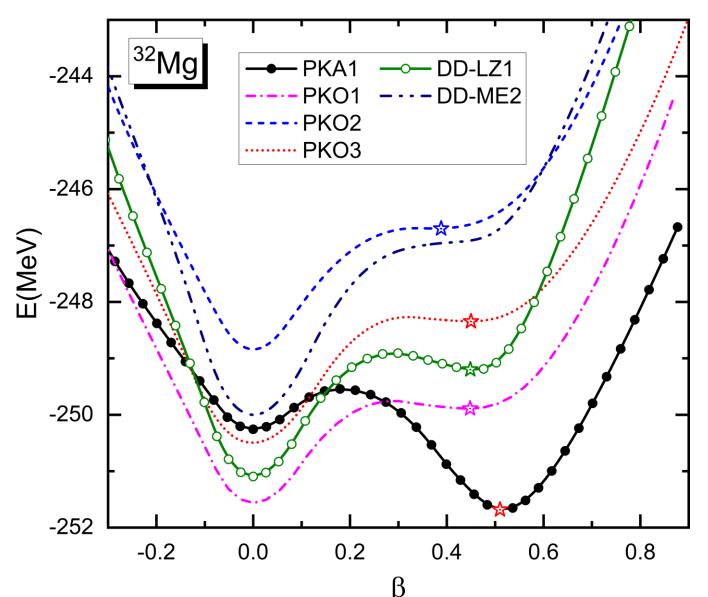
What can we learn from the nature of the island of inversion exhibited by ³²Mg



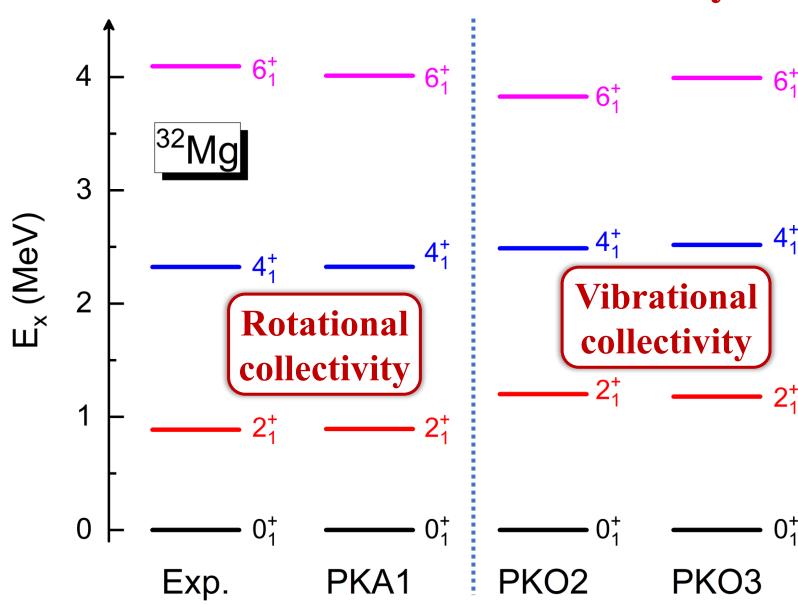
Nature of the island of inversion exhibited by ³²Mg



☐ D-RHFB calculations: **deformation**



☐ CI-RHF calculations: collectivity



Link between D-RHFB and CI-RHF

Part of many-body correlations in CI-RHF model are considered by the D-RHFB model

Peng, Liu, Geng, Niu, WHL, CPC 49, 064112 (2025)



Island of inversion: underlying mechanism



☐ Cross-shell excitation: deformation picture in ³²Mg

$R = E(4_1^+)/E(2_1^+)$	R = E	$(4_1^+)_{1}$	$/E(2_1^+)$
-------------------------	-------	---------------	-------------

	$1d_{5/2}$	$2s_{1/2}$	$1d_{3/2}$	$1f_{7/2}$	$2p_{3/2}$	R	
PKA1	5.75	1.91	2.27	1.50	0.56	2.61	
PKO2	5.73	1.92	2.38	1.81	0.16	2.07	
PKO3	5.73	1.92	2.56	1.57	0.23	2.14	

stable deformation

Vibrational collectivity

soft deformation

Rotational collectivity



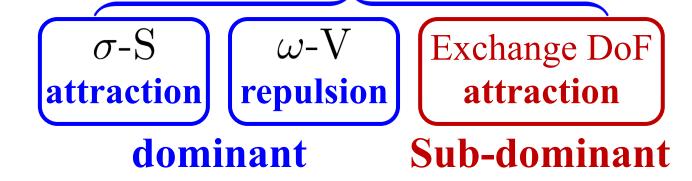
Peng, Liu, Geng, Niu, WHL, CPC 49, 064112 (2025)

Impact of exchange DoFs: binding of ³²Mg

	$E_{\rm kin.}$	$E_{\sigma+\omega}$	$E_{ ho}$	E_{π}	$E_{\rm oth.}$	E
PKA1	400.8	-417.0	-217.5	-31.5	14.9	-250.2
	409.6	-434.4	-214.4	-29.9	17.4	-251.7
PKO3	438.1	-606.1	-59.7	-37.2	14.4	-250.5
	448.7	-616.3	-59.2	-36.1	14.6	-248.3
DD-LZ1	412.6	-701.3	18.3		19.4	-251.1
	414.1	-702.7	17.6	_	21.8	-249.2

Exchange DoFs exhibit substantial contributions to nuclear binding

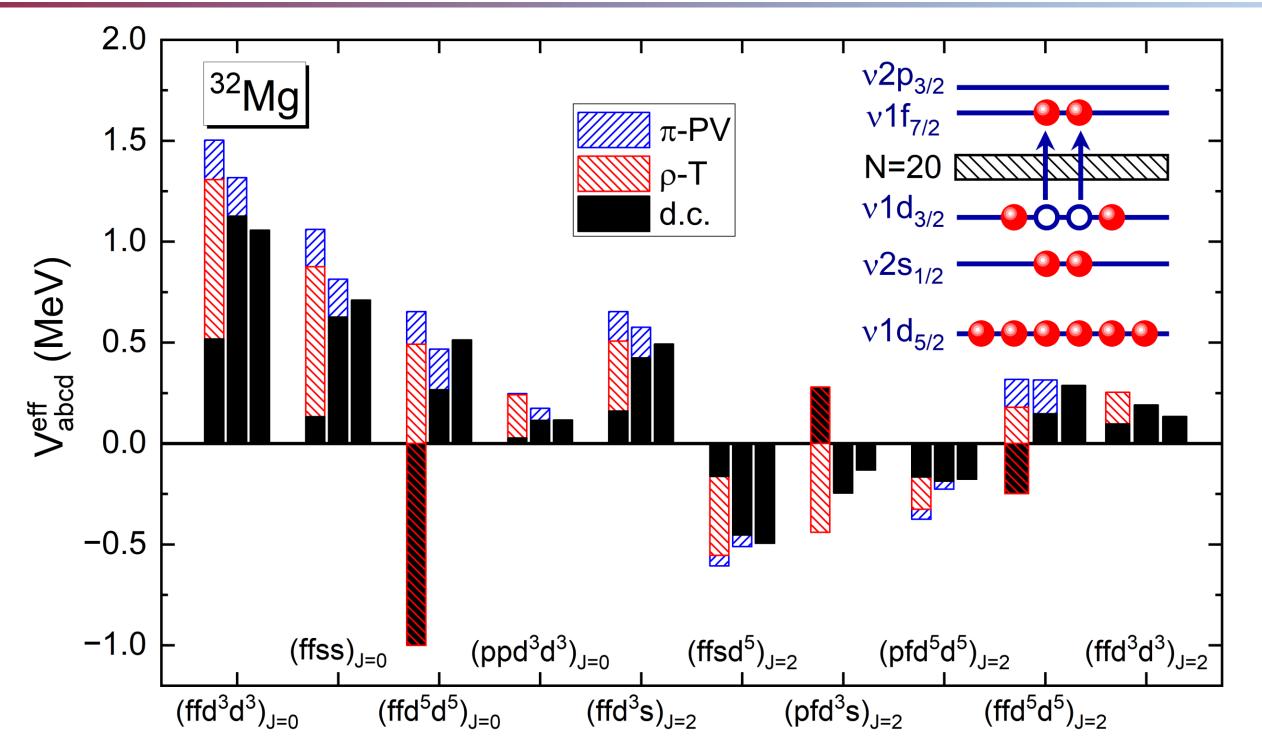
Picture of nuclear binding





Configuration Interactions: ρ -T and π -PV





PKA1 (left)
PKO3 (mid.)
PKO2 (right)



Discernible π -PV & compressive ρ -T



Be determined at the Hartree-Fock level

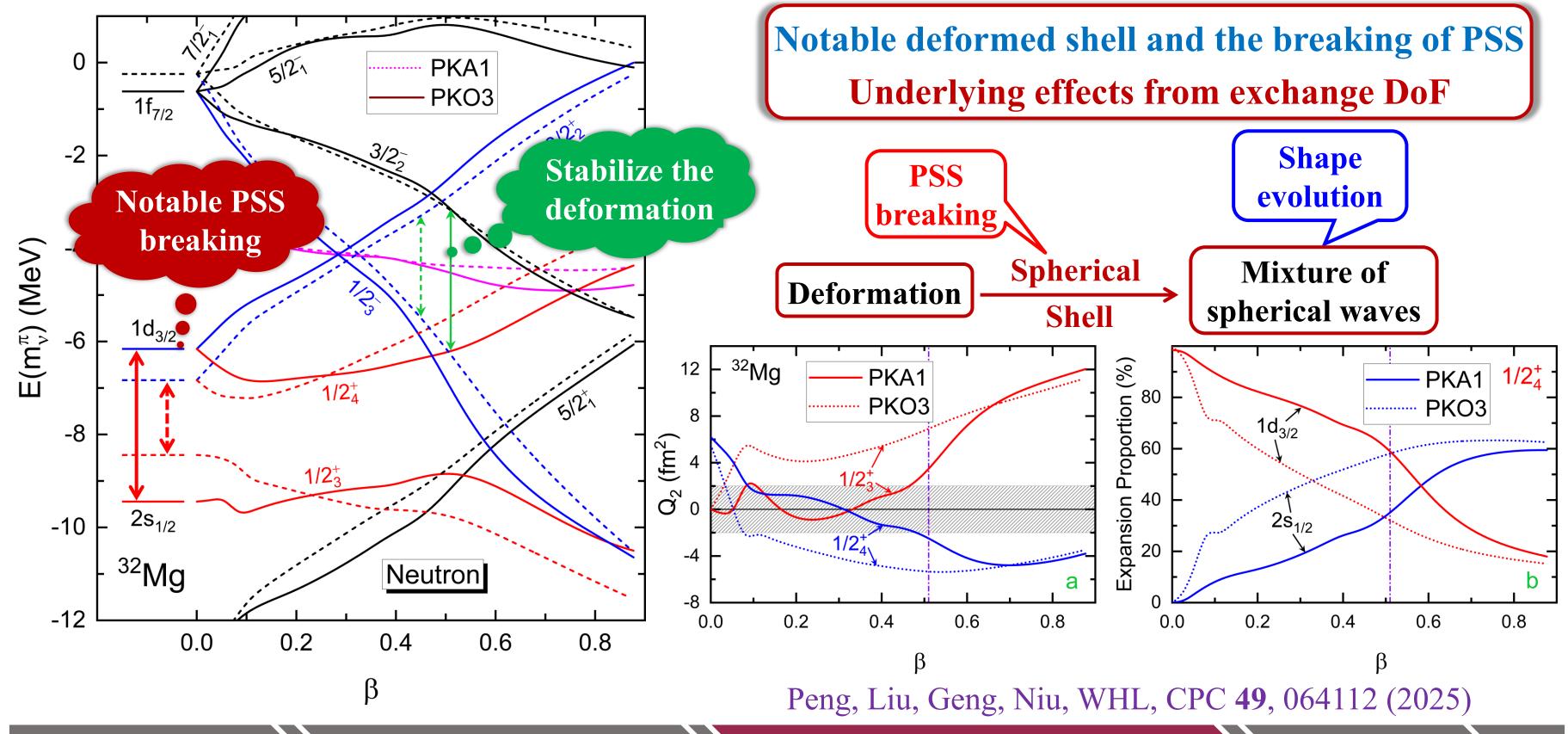
Configuration Interactions shaped by the exchange DoFs - deformation stability?

Peng, Liu, Geng, Niu, WHL, CPC 49, 064112 (2025)



PSS breaking: an implemented insight



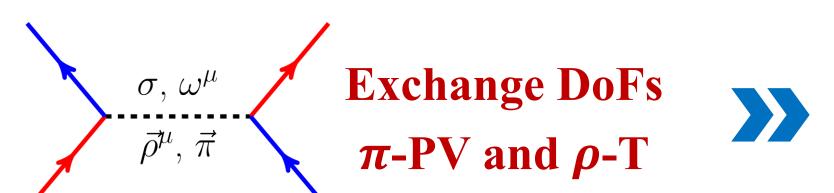




Conclusions: exchange correlation effects



☐ At the mean field level & beyond, exchange DoFs are more important than expected

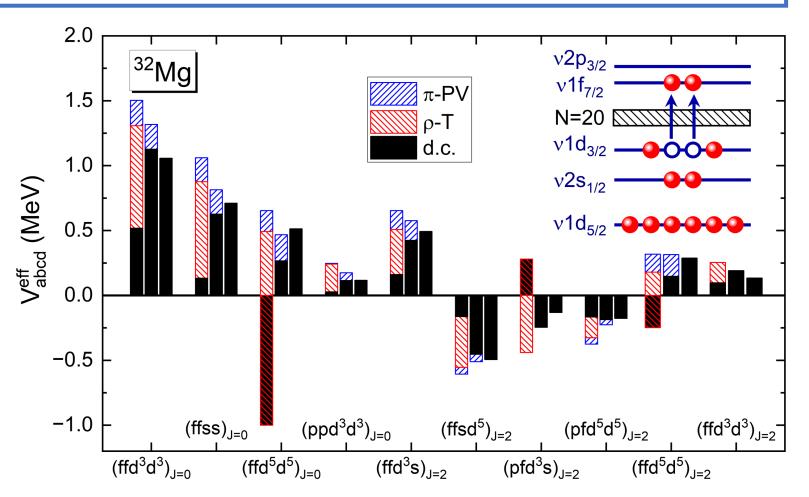


■ Exchange DoFs shall be predetermined at the mean field level

Picture of nuclear binding

How nucleons form an atomic nucleus

- Impact the modeling of nuclear binding
- > Enhance the neutron-proton interactions



Explicit Fock terms are complicated, which calls for an economical solution

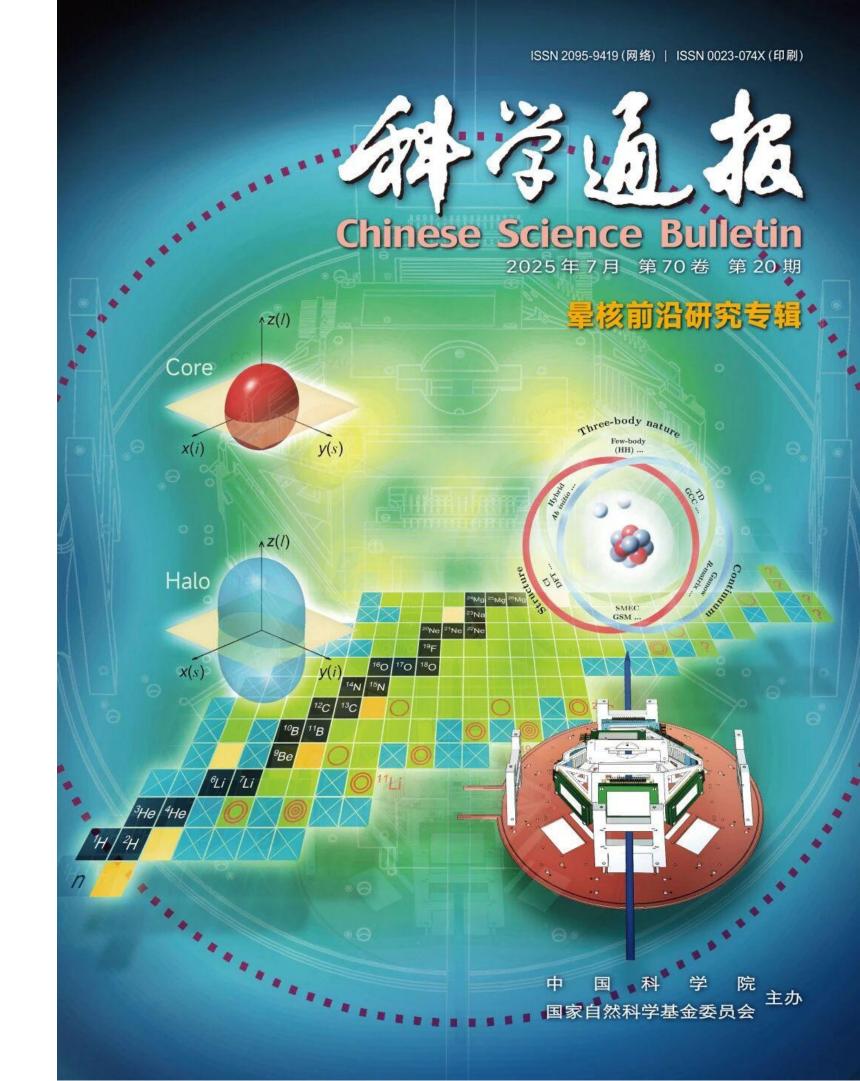
Acknowledgements

Collaborators:

- J. Geng, Y. Peng, Y.F. Niu @LZU
- J. Liu @USC
- S.S. Zhang, J.L. An, H.C. Li @BUAA
- J. Meng, P.W. Zhao @PKU
- J.J. Li, Z.P. Li @SWU
- P. Ring @TUM
- N. Van Giai @Paris-Saclay University
- C.A. Bertulani @ Texas A&M University-Commerce

Thank you for your attentions

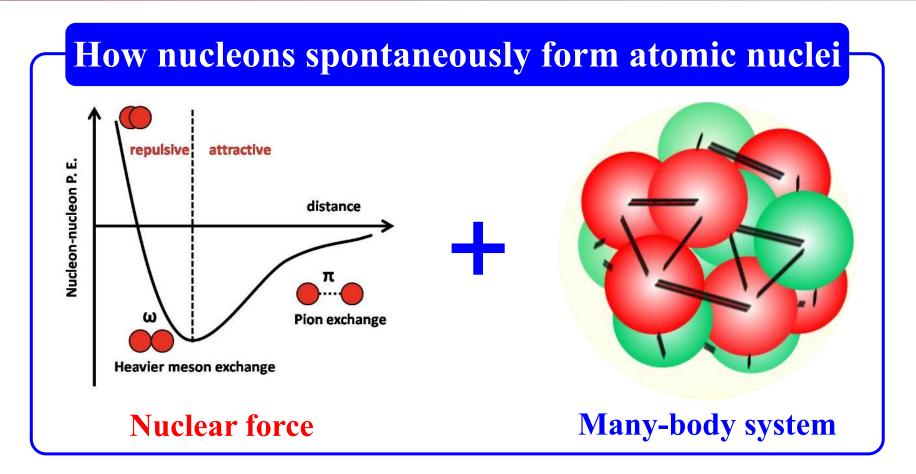
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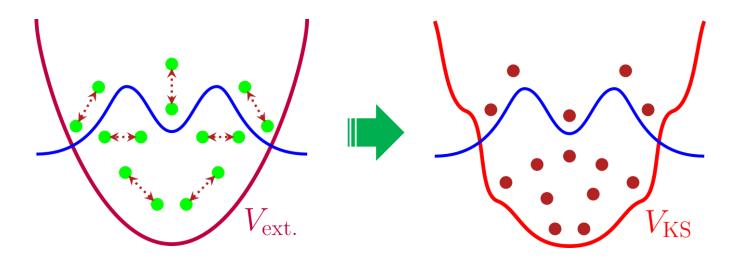
Fundamental issue in nuclear physics





Density functional theory (DFT)

many-body problem → one-body problem



Drut, Furnstahl, Platter, PPNP 64 (2010) 120

☐ A relativistic representation of nuclear binding given by relativistic DFT

Interplay between strong scalar attraction and vector repulsion

bound nucleus

Natural spin-orbit coupling, the origin of pseudo-spin symmetry and spin symmetry, ...

Meng et al., PPNP 57, 470 (2006); Ginocchio, PRL 78, 436 (1997); Meng et al., PRC 58, R628 (1998); Zhou et al., PRL 91, 262501 (2003).

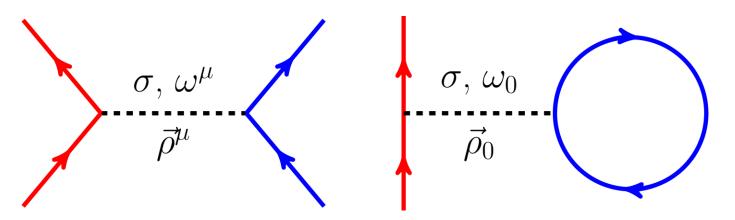
What can we learn from the novel phenomena exhibited by unstable nuclei?

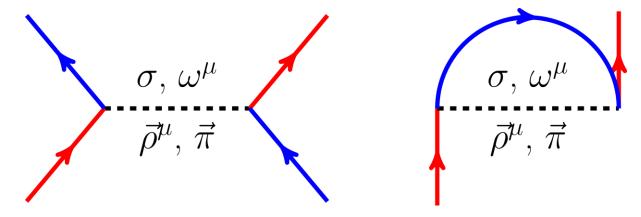


Relativistic Hartree-Fock (RHF) theory



☐ Fock terms give rise to pure exchange degrees of freedom





Hartree terms: σ -S, ω -V, $\overrightarrow{\rho}$ -V, A-V

Fock terms: σ -S, ω -V, $\vec{\rho}$ -V, A-V; π -PV, ρ -T (ρ -VT)

Bouyssy (1987), Bernardos (1993), Shi (1995), Marcos (2004), Long (2004-now), Geng (2020-now), ...

- Exchange degrees of freedom (DoF) impact the modelling of nuclear binding π -PV and ρ -T (ρ -VT) PRC 100, 051301(R) (2019)
- Nuclear tensor force can be naturally introduced: e.g., via the π -PV and ρ -T couplings PRC 91, 034326 (2015); CPC 42, 024101 (2018); PRC 98, 034313 (2018); PRC 101, 064306 (2020)
- □ Not exactly a DFT: explicit Fock terms (meson exchange) are not a functional of local density

$$E(\widehat{\rho}) = \langle \Psi | \widehat{H} | \Psi \rangle \approx \langle \Phi | H_{\text{eff.}} | \Phi \rangle = T + E_H(\rho) + E_{\text{xc}}(\rho)$$

Exchange correlations



Axially deformed RHF (D-RHF) model



☐ Spherical Dirac Woods-Saxon (DWS) basis

Zhou, Meng, Ring, PRC 68, 034323 (2003)

- > Appropriate asymptotic behaviors of wave functions
- > Microscopic insight into deformed nuclei: effects of tensor force
- ☐ Yukawa propagator: the decomposition in spherical coordinate

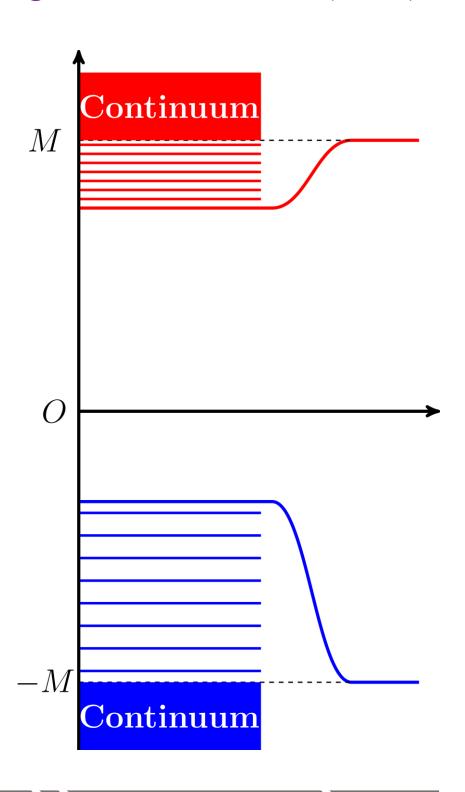
$$\frac{1}{4\pi} \frac{e^{-m|\mathbf{x}-\mathbf{x}'|}}{|\mathbf{x}-\mathbf{x}'|} = \frac{1}{rr'} \sum_{\lambda=0}^{\infty} I_{\lambda+\frac{1}{2}}(mr_{<}) K_{\lambda+\frac{1}{2}}(mr_{>}) \sum_{\mu=-\lambda}^{\lambda} Y_{\lambda\mu}(\mathbf{\Omega}) Y_{\lambda\mu}^{*}(\mathbf{\Omega})$$

$$\Rightarrow \text{Singularity appears when } \mathbf{x} = \mathbf{x}'$$

$$V_{\phi} = \frac{1}{2} \int d\boldsymbol{x} d\boldsymbol{x}' \sum_{\alpha\beta;\alpha'\beta'} \left[\bar{\psi}_{\alpha} \Gamma_{\phi} \psi_{\alpha'} \right]_{\boldsymbol{x}} \frac{1}{4\pi} \frac{e^{-m_{\phi}|\boldsymbol{x}-\boldsymbol{x}'|}}{|\boldsymbol{x}-\boldsymbol{x}'|} \left[\bar{\psi}_{\beta} \Gamma^{\phi} \psi_{\beta'} \right]_{\boldsymbol{x}'} c_{\alpha}^{\dagger} c_{\beta}^{\dagger} c_{\beta'} c_{\alpha'}$$

 $(\lambda \mu)$ -terms are truncated by the DWS basis to avoid singularity

Geng, Xiang, Sun, WHL, PRC 101, 064302 (2020)

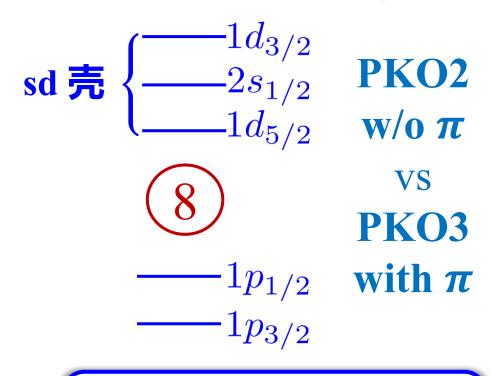




Tensor force effects in ²⁰Ne



☐ Tensor force changes the mixing of spherical waves



 π -T enhances the intrude of $1d_{3/2}$ waves

Shell effects enhanced by π -PV

PRC 101, 064302 (2020)

